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### Article

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# Supernova environments in J-PLUS

## Normalized cumulative-rank distributions and stellar-population synthesis combining narrow- and broad-band filters<sup>★</sup>

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### ABSTRACT

We investigated the local environmental properties of 418 supernovae (SNe) of all types using data from the Javalambre Photometric Local Universe Survey (J-PLUS), which includes five broad-band and seven narrow-band imaging filters. Our study involves two independent analyses: (1) the normalized cumulative-rank (NCR) method, which utilizes all 12 single bands along with five continuum-subtracted narrow-band emission and absorption bands, and (2) simple stellar population (SSP) synthesis, where we build spectral energy distributions (SED) of the surrounding 1 kpc<sup>2</sup> SN environment using the 12 broad- and narrow-band filters. Improvements on previous works include: (i) the extension of the NCR technique to other filters (broad and narrow) and the use a set of homogeneous data (same telescope and instruments); (ii) a correction for extinction to all bands based on the relation between the  $g - i$  color and the color excess  $E(B - V)$ ; and (iii) a correction for the contamination of the [N II]  $\lambda 6583$  line that falls within the  $H\alpha$  filter. All NCR distributions in the broad-band filters, tracing the overall light distribution in each galaxy, are similar to each other. The main difference is that type Ia, II, and IIb SNe are preferably located in redder environments than the other SN types. The radial distribution of the SNe shows that type IIb SNe seem to have a preference for occurring in the inner regions of galaxies, whereas other types of SNe occur throughout the galaxies without a distinct preference for a specific location. For the  $H\alpha$  filter we recover the sequence from SNe Ic, which has the highest NCR, to SNe Ia, which has the lowest; this is interpreted as a sequence in progenitor mass and age. All core-collapse SN types are strongly correlated to the [O II] emission, which also traces star formation rate (SFR), following the same sequence as in  $H\alpha$ . The NCR distributions of the Ca II triplet show a clear division between II-IIb-Ia and Ib-Ic-IIn subtypes, which is interpreted as a difference in the environmental metallicity. Regarding the SSP synthesis, we found that including the seven J-PLUS narrow filters in the fitting process has a more significant effect on the core-collapse SN environmental parameters than for SNe Ia, shifting their values toward more extincted, younger, and more star-forming environments, due to the presence of strong emission lines and stellar absorptions in those narrow bands.

**Key words.** methods: observational – methods: statistical – techniques: photometric – supernovae: general – galaxies: general – galaxies: photometry

<sup>★</sup> Full Tables A.1–A.5 are available at the CDS via anonymous ftp to [cdsarc.cds.unistra.fr](https://cdsarc.cds.unistra.fr) (130.79.128.5) or via <https://cdsarc.cds.unistra.fr/viz-bin/cat/J/A+A/684/A104>

## 1. Introduction

Supernovae (SNe) are one of the final stages in stellar evolution and are key in driving the chemical evolution of galaxies. Classical SNe (excluding superluminous ones) are basically divided in two main types: thermonuclear (those of type Ia) and core-collapse (CC). Type Ia SNe are those triggered by the thermonuclear explosion of a carbon and oxygen white dwarf (WD; Hoyle & Fowler 1960). Different progenitor scenarios, such as WDs reaching the Chandrasekhar limit ( $M_{Ch} \sim 1.44 M_{\odot}$ ) by mass accretion in a binary system or mergers of WDs, and explosion mechanisms, such as internal detonations of  $M_{Ch}$ , WD, or surface explosions on WDs with masses lower than  $M_{Ch}$ , have been proposed to explain how the WDs explode (see reviews by, e.g., Maoz et al. 2014). No direct progenitor detection has been reported to date (however, see McCully et al. 2014), but in all cases, the main spectral features are the lack of H lines and the presence of strong Si II lines up to a couple of weeks after maximum brightness (Filippenko 1997). SNe Ia occur in galaxies of all types, including elliptical galaxies which only contain old stellar populations, providing strong evidence that SNe Ia have long-lived, low-mass progenitors (Han et al. 2010). SNe Ia are more luminous than CC SNe, comprising about 30% of the observed SNe (Graur et al. 2017). Their distinctive feature lies in the fact that generally normal SNe Ia exhibit similar spectra and light curves, making them valuable for tracing cosmological distances (e.g., Phillips et al. 1999; Riess et al. 1996; Perlmutter et al. 1997, 1999; Astier et al. 2006).

Core collapse SNe are those resulting from the gravitational collapse of the iron core of a massive star ( $>8 M_{\odot}$ ; Arnett et al. 1989). CC SNe are divided into three main subtypes depending on their spectral features, which reflect the state of the outer layers of the progenitor star at the moment of explosion. Type II SNe show H lines because their progenitors have kept the H-rich outer envelope intact, type Ib SNe show He but no H since the progenitor has lost the H envelope, and type Ic SNe lack both H and He lines due to the loss of both H- and He-rich layers (Filippenko 1997). Additionally, SNe IIn show narrow H emission lines most probably due to interaction with circumstellar material (Smith et al. 2011), and SNe Iib are those intermediate between SNe II and Ib, that show H only for a few days after explosion indicating that there was still a thin H layer before explosion (Filippenko 1988; Nomoto et al. 1993). There is a debate on the process responsible for the mass-loss leading to stripped envelope (SE) SN types (Ib, Ic, Iib). Mass-loss through radiation-driven wind from single hot massive stars or the removal of the outer envelope via tidal stripping by a less massive companion in a binary are the two most viable possibilities (Filippenko 1997; Gal-Yam et al. 2017; Prentice & Mazzali 2017; Shivvers et al. 2017; Taddia et al. 2018).

One outstanding question in the field of massive stellar evolution concerns the link between the nature of the progenitor and the resultant SN. Direct constraints on the nature of SN progenitors are limited. The bulk of these come from the small number of nearby events where deep, high-resolution (usually from the *Hubble* Space Telescope) pre-explosion imaging exists to directly image the progenitor star (e.g., Maund et al. 2013; Smartt 2015; Van Dyk 2017). For almost all discovered SNe, this is unfeasible. As such, methods investigating statistical properties of SNe have been developed to exploit the large numbers of observed SN environments (pre-SN or once the SN has faded), including historical ones. One of those statistical methods is the normalized cumulative-rank (NCR; Fruchter et al. 2006; James & Anderson 2006) method. The NCR method,

when applied to H $\alpha$  narrow imaging, measures the correlation between SN locations and star-formation intensity (i.e., H $\alpha$  emission) in their host galaxies. The strength of this correlation is an indicator of the progenitor lifetime, a proxy for the initial mass of the progenitor. The method has been employed on samples of SNe to deduce a sequence of ascending mass for the common SN types Ia, II, Ib, and Ic (Anderson & James 2008; Anderson et al. 2012; Kangas et al. 2013). Subsequently, Kangas et al. (2017) established a connection between the NCRs of these SNe and those of massive stars in the LMC and M33. This link indicated initial masses of approximately  $>20 M_{\odot}$  for Ic SNe progenitors and  $>9 M_{\odot}$  for II and Ib SNe. For the stripped-envelope SN types, Ic and Ib, these findings would imply that single Wolf-Rayet stars and/or massive binaries are the most probable progenitors for the former, while for the latter interacting relatively low-mass binary systems, such as the directly detected 10–12  $M_{\odot}$  progenitor of the SN Ib iPTF13bvn (Eldridge & Maund 2016), would be dominant.

Another method that has proved useful in putting constraints on SN progenitors is the study of the stellar populations at SN locations. Galbany et al. (2014, 2016) presented the first statistical study of nearby SN host galaxies using integral field spectroscopy. These were provided by the CALIFA survey, which consists of 132 SNe of all types in 115 galaxies, and it was later extended to 272 SNe in 232 galaxies in Galbany et al. (2018). Stellar parameters were inferred by fitting a set of single stellar population (SSP) models to the spectra of the locations where SNe occurred. They found that CC SNe tend to explode at positions with younger stellar populations than the galaxy average, while at Ia SN locations local properties were, on average, the same as global ones. They also found a sequence from higher to lower metallicity, from SN Ia to SN Ic-BL, and a significant increasing relative number of SNe Ic at higher metallicities compared to other CC SNe types, which supports a large fraction of this SN subtype occurring under the single-star scenario (see also Kuncarayakti et al. 2013, 2018).

In this paper, we present a comprehensive analysis of 418 SNe environments using the 12 filter system of the Javalambre Photometric Local Universe Survey (J-PLUS<sup>1</sup>; Cenarro et al. 2019), whose spatial resolution ( $\lesssim 382$  pc, corresponding to an angular resolution of 1.14 arcsec at  $z = 0.016$ ) and overall data quality of this dataset allow us to study the SN environments. Our investigation involves two main methodologies: the normalized cumulative-rank (NCR) analysis and the simple stellar population (SSP) synthesis. Compared to previous works in the literature, our study offers several advancements. Firstly, we extend the NCR technique to a number of broad- and narrow-band filters other than H $\alpha$ , which enables us to explore correlations with new spectral lines. This extension opens up possibilities to probe additional environmental factors, such as stellar population metallicity or age. Previous studies, such as Kangas et al. (2013, with UV imaging) and Anderson et al. (2015, with NIR imaging), have explored these factors to some extent. However, our research operates on a much larger scale and utilizes a set of homogeneous data acquired from the same telescope and instrument. Another improvement over past works is the correction for the contamination from the [N II]  $\lambda 6583$  line that falls within the H $\alpha$  filter. Vilella-Rojo et al. (2015) presented a method to correct for such extra emission by combining narrow- and broad-band data and providing pure H $\alpha$  emission from J-PLUS observations. This method can also be extended to other narrow-band filters.

<sup>1</sup> [www.j-plus.es](http://www.j-plus.es)

## 2. Data

### 2.1. J-PLUS

The Javalambre Photometric Local Universe Survey is designed to observe 8500 deg<sup>2</sup> of the northern sky from the Observatorio Astrofísico de Javalambre (OAJ, Teruel, Spain; [Cenarro et al. 2014](#)) with the 83 cm Javalambre Auxiliary Survey Telescope (JAST80) and T80Cam, a panoramic camera of 9.2k × 9.2k pixels that provides a 2 deg<sup>2</sup> field of view (FoV) with a pixel scale of 0.55 arcsec pix<sup>-1</sup> ([Marín-Franch et al. 2015](#)). The J-PLUS filter system comprises the 12 broad, intermediate and narrow-band optical filters. J-PLUS is particularly designed to carry out the photometric calibration of the Javalambre Physics of the Accelerating Universe Astrophysical Survey (J-PAS<sup>2</sup>; [Benitez et al. 2014](#)), which has observed the northern sky from Javalambre using 59 optical narrow-band filters. For this reason, some J-PLUS filters are located at key stellar spectral features that allow us to retrieve very accurate spectral energy distributions (SED) for more than five millions of stars in our galaxy. The 12 filters of J-PLUS are listed in Table 1 and their transmission is shown in Fig. 1, together with two template spectra of a star-forming and a passive galaxies to highlight the regions of interest covered by the intermediate- and narrow-band filters. The third J-PLUS Data Release (DR3<sup>3</sup>) was made public in December 2022, comprising 1642 pointings observed in the 12 optical bands amounting to ~3200 sq. deg with around 30 million sources detected at magnitudes  $r < 21$ . We cross-matched all SN coordinates from the Open Supernova Catalogue (OSC; [Guillochon et al. 2017](#)) to the central coordinates of all J-PLUS DR3 tiles, and found 2168 SNe positions within the 2 sq. deg of the J-PLUS tile. The width of the narrow H $\alpha$  filter (J0660 filter; ~150 Å) puts an upper limit on the redshift of galaxies that are useful for an NCR study at  $z \sim 0.0163$  (or about 60 Mpc), since the J0660 transmission falls down dramatically at 6672 Å. Thus, in order to keep the H $\alpha$  emission of the SN host galaxy within the coverage of the J-PLUS J0660 filter, we excluded all objects with a redshift larger than 0.0163. Moreover, NCR becomes less useful with distance, and samples at significantly different median distances may not be directly comparable ([Kangas et al. 2017](#)). This redshift cut is passed by 282 SNe of the following types: 88 SNe Ia, 126 SNe II, 7 SNe IIn, 22 SNe Ib, 17 SNe Ic, 17 SNe Iib, and another five type Ibc SNe.

### 2.2. A dedicated SN host galaxy program at JAST80

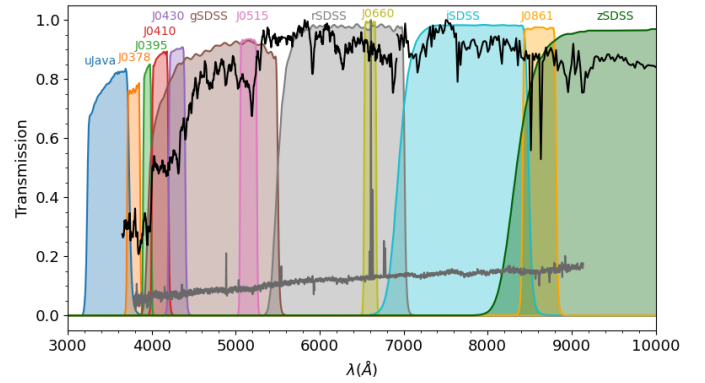
To reach a significantly large number of SNe that cover all types compared to previous studies, and within the redshift limit imposed by the J-PLUS H $\alpha$  filter width, we had to include objects outside the J-PLUS footprint, taking special care to increase the number of less represented subtypes (Ic, Ib, Iib, IIn). With this objective in mind, we initiated a dedicated program to acquire 73 additional fields using the same instrumental configuration and exposure times as in J-PLUS, which involves T80 and 12 filters. These fields encompassed the positions and host galaxies of 136 SNe. For each of the 12 filters, we obtained three individual frames, all of which were subsequently processed using the same J-PLUS pipeline ([Cenarro et al. 2019](#)).

Summing up the observations from the main survey and our dedicated program, 418 SNe are included in our sample, and they are listed in Table A.1. Taking into account the main SN types, we have 156 SNe II, 20 SN IIn, 108 SN Ia, 51 SN Ib, 49 SN Ic,

**Table 1.** J-PLUS photometric system properties.

Filter name	Central $\lambda$ (nm)	FWHM (nm)	Comments
uJAVA	348.5	50.8	ultraviolet continuum
J0378	378.5	16.8	[O II] emission
J0395	395.0	10.0	Ca H+K absorption
J0410	410.0	20.0	H $\delta$ emission
J0430	430.0	20.0	G-band absorption
gSDSS	480.3	140.9	SDSS green continuum
J0515	515.0	20.0	Mgb Triplet absorption
rSDSS	625.4	138.8	SDSS red continuum
J0660	660.0	14.5	H $\alpha$ emission
iSDSS	766.8	153.5	SDSS near infrared continuum
J0861	861.0	40.0	Ca II triplet absorption
zSDSS	911.4	140.9	SDSS near infrared continuum

**Notes.** Transmission curves can be found at the J-PLUS website: <http://www.j-plus.es/survey/instrumentation>.



**Fig. 1.** Transmission of 12 J-PLUS filters, on top of a typical star-forming (in gray) and passive galaxy (in black) spectra. Labels, which represent the filter names, highlight the main spectral features covered by the intermediate and narrow band filters.

27 SN Iib, and seven other type Ibc SNe. The analysis is carried out with the sample divided into these seven main SN types, where Ibc SNe include Ib, Ic, Iib, and other Ibc.

## 3. Generation of individual 3D data cubes

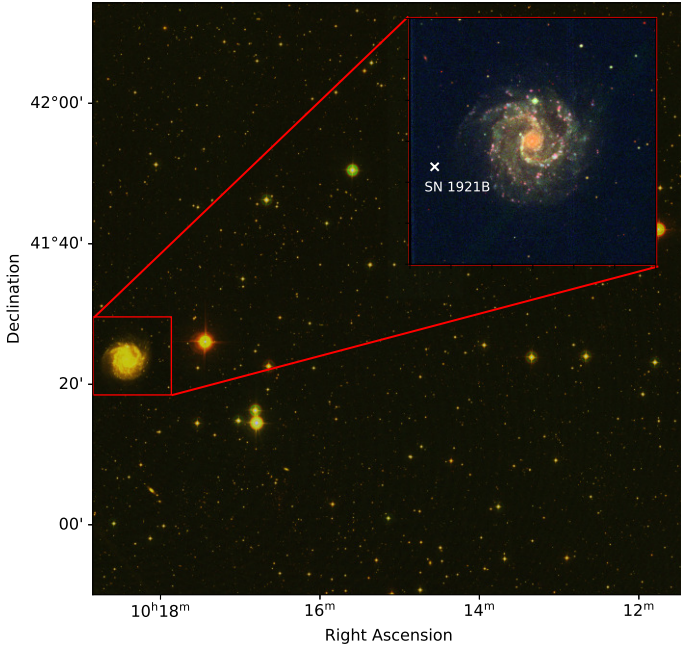
Starting from the initial 12 images per field, one for each filter, we performed a number of steps to produce one 3D data cube per galaxy, namely galaxy cutout, flux calibration, correction for dust extinction, and continuum subtraction. All these steps are described in the following subsections.

### 3.1. Galaxy cutouts

J-PLUS images cover ~2 sq. deg of the sky with a pixel scale of 0.55 arcsec, with a median seeing of 1.14 arcsec, corresponding to  $\lesssim 382$  pc at  $z = 0.016$ . After registering all 12 images to the same reference frame, we took a squared cutout of a size that includes the full extent of the SN host galaxy centered at its core, defining the size by visual inspection (see Fig. 2). The 12 cutouts are then stored in a 3D cube, mimicking integral field spectroscopy observations, where two spatial dimensions correspond to RA and Dec, and the third dimension provides a spectral energy distribution (SED) at any position within the cutout.

<sup>2</sup> <http://www.j-pas.org/survey>

<sup>3</sup> [http://www.j-plus.es/datareleases/data\\_release\\_dr3](http://www.j-plus.es/datareleases/data_release_dr3)



**Fig. 2.** Example of J-PLUS field frame. The zoomed-in image corresponds to the cutout of the galaxy NGC 3184, which contains the SN 1921B; marked in white. The images were constructed combining the uJAVA (blue), gSDSS (green), and rSDSS (red) images as false colors.

### 3.2. Flux calibration

To convert the electronic counts stored in each pixel to radiative flux, we use the expression from [Logroño-García et al. \(2019\)](#):

$$F_{\lambda} = C \cdot 10^{-0.4(ZP+48.6)} \frac{c}{\lambda_{\text{pivot}}^2}, \quad (1)$$

where  $C$  is the number count in the pixel,  $c$  is the speed of light,  $ZP$  is the zero point of the band used for the calibration to the standard AB magnitude system, and  $\lambda_{\text{pivot}}$  is the pivot wavelength of the filter, which is a source-independent measurement of the characteristic wavelength of a given pass band by

$$\lambda_{\text{pivot}}^2 = \frac{\int T(\lambda) d\lambda}{\int T(\lambda) \lambda^{-2} d\lambda}, \quad (2)$$

where  $T(\lambda)$  represents the transmission curve of the filter. The values of  $ZP$  and pivot wavelength are given in Table 2.

### 3.3. Dust extinction correction

The Milky Way interstellar dust reddening was corrected using the python code DUSTMAPS ([Green 2018](#)) and the *Gaia* Total Galactic Extinction map ([Delchambre et al. 2023](#)) as the source to obtain the  $E(B - V)$  color excess for every galaxy. In addition, we applied a correction to the reddening due to dust in the SN host galaxy. Following [Calzetti et al. \(2000\)](#), the difference between the intrinsic and observed flux due to extinction is given by

$$F_i(\lambda) = F_0(\lambda) 10^{0.4E(B-V)k'(\lambda)}, \quad (3)$$

where  $F_i(\lambda)$  and  $F_0(\lambda)$  are the intrinsic and observed fluxes at a wavelength  $\lambda$ ,  $E(B - V)$  is the color excess, and  $k'(\lambda)$  is the extinction law: an empirical relationship between the amount

of extinction and wavelength. We employed the [Calzetti et al. \(2000\)](#) extinction law, defined as

$$k'(\lambda) = \begin{cases} 2.659(-1.857 + 1.040/\lambda) + R'_V & \text{for } 0.63 \mu\text{m} \leq \lambda \leq 2.20 \mu\text{m} \\ 2.659(-2.156 + 1.509/\lambda - 0.198/\lambda^2 + 0.011/\lambda^3) + R'_V & \text{for } 0.12 \mu\text{m} \leq \lambda \leq 0.63 \mu\text{m}, \end{cases} \quad (4)$$

where  $R'_V$  is the obscuration in the  $V$  band. Here, we used the same value found by [Calzetti et al. \(2000\)](#),  $R'_V = 4.05 \pm 0.80$ , for stellar extinction<sup>4</sup>. To estimate the color excess  $E(B - V)$ , we followed [Vilella-Rojo et al. \(2015\)](#), who found a relation between the observed  $g' - i'$  color and the spectroscopically measured  $E(B - V)$  obtained after convolving the SDSS spectra with the J-PLUS photometric system. These authors obtained the following expression by fitting a power-law function:

$$E(B - V) = 0.206(g' - i')^{1.68} - 0.0457, \quad (5)$$

and assuming  $E(B - V) = 0$  for  $g' - i' < 0.4$ .

This is an empirical expression for the gas extinction. [Vilella-Rojo et al. \(2015\)](#) mixed in the same process of correction to the stellar and gas components for all the J-PLUS data; so, we used the same method to be coherent. The gas extinction is, at minimum, the same as the stellar extinction. However, on average, the extinction in HII regions can be twice as high ([Cid et al. 2005](#)). Therefore, it is essential to note that the extinction correction applied by this method represents a conservative lower bound.

With all these ingredients, we corrected the fluxes of the 12 images for all SNe in our sample, pixel by pixel. This improvement results in a change in our NCR values compared to studies that do not apply this correction to the fluxes. Hence, we caution against making direct comparisons of NCR values to previous works that employed different methodologies.

### 3.4. J0660 continuum subtraction

J-PLUS narrow-band filters were designed to cover wavelength regions of a number of emission lines of interest. However, they also include the flux of the underlying stellar continuum that needs to measure the gas-phase component exclusively. In particular, both the rSDSS and J0660 filters are collecting the flux of the  $H\alpha$  emission line, so to study the  $H\alpha$  emission, which is a tracer of the star formation, we have to subtract the stellar continuum from the J0660 flux.

[Vilella-Rojo et al. \(2015\)](#) presented two methods to remove the underlying continuum from a narrow-band filter. The simplest is by combining the narrow band that contains the emission with an adjacent or overlapping broad filter that traces the continuum. For  $H\alpha$ , we can use the J0660 and the rSDSS filters, but taking into account that J0660 also includes the contribution of the forbidden transition of [N II] doublet (at 654.8 nm and 658.53 nm). Assuming a flat continuum, the flux of the three emission lines can be recovered using the following expression:

$$F_{H\alpha+[N II]} = \Delta_{J0660} \frac{\bar{F}_{J0660} - \bar{F}_{rSDSS}}{1 - \frac{\Delta_{J0660}}{\Delta_{rSDSS}}}, \quad (6)$$

where  $\bar{F}_{J0660}$  and  $\bar{F}_{rSDSS}$  correspond to the average flux inside the J0660 and rSDSS filters, respectively, and  $\Delta_x$  is defined for

<sup>4</sup> The stellar extinction mainly affects those filters without emission lines, and gas extinction mainly affects those with emission lines ([Thomas et al. 2013](#)).

**Table 2.** Pivot wavelengths and zero points of the J-PLUS bands.

Filter	uJAVA	J0378	J0395	J0410	J0430	gSDSS
$\lambda_{\text{pivot}}$ (nm)	352.29	378.64	395.06	410.07	430.04	474.46
ZP (AB mag)	21.078 (049)	20.416 (036)	20.332 (037)	21.283 (029)	21.380 (028)	23.553 (035)
Filter	J0515	rSDSS	J0660	iSDSS	J0861	zSDSS
$\lambda_{\text{pivot}}$ (nm)	514.98	622.98	659.98	767.65	860.25	892.20
ZP (AB mag)	21.500 (040)	23.494 (023)	20.953 (023)	23.136 (023)	21.439 (007)	22.498 (007)

any pass band  $x$  at any wavelength of interest  $\lambda_s$  as

$$\Delta_x \equiv \frac{\int P_x(\lambda)\lambda d\lambda}{P_x(\lambda = \lambda_s)\lambda_s}, \quad (7)$$

where  $P_x$  is the transmission of the pass band  $x$  as a function of wavelength. In our case, we took  $\lambda_s = \lambda_{H\alpha}$ .

This method assumes a linear continuum, so the  $H\alpha$  absorption is not taken into account, resulting in an overall bias of approximately 9% in all the results (Vilella-Rojo et al. 2015). When including more filters, such as those available in J-PAS, this method of continuum subtraction improves significantly (Martínez-Solaache et al. 2021).

### 3.5. [N II] removal

The continuum-subtracted  $J0660$  flux includes both the  $H\alpha$  and the [N II] doublet emission. To remove the [N II] contribution, we used the bimodal empirical relation between the spectroscopic dust-corrected  $H\alpha$  flux and the total  $H\alpha + [\text{N II}]$  flux found by Vilella-Rojo et al. (2015). This bimodality can be disentangled by using the same  $g' - i'$  color used for the dust correction. Vilella-Rojo et al. (2015) obtained the following expressions by fitting a line to each branch:

$$\log(F_{H\alpha}) = \begin{cases} 0.989 \log(F_{H\alpha + [\text{N II}]}) - 0.193, & \text{if } g' - i' \leq 0.5 \\ 0.954 \log(F_{H\alpha + [\text{N II}]}) - 0.753, & \text{if } g' - i' > 0.5 \end{cases}. \quad (8)$$

In addition, we can easily obtain the [N II] flux by simply taking the difference between the  $H\alpha$  flux and  $F_{H\alpha + [\text{N II}]}$ ,

$$F_{[\text{N II}]} = F_{H\alpha + [\text{N II}]} - F_{H\alpha}. \quad (9)$$

We did not apply this procedure for SNe at  $z > 0.014$ , since the [N II]  $\lambda 6583$  line is outside the  $J0660$  filter range at a higher redshift and a subtraction of this line cannot be correctly performed. Therefore, we assume all emission in the  $J0660$  filter comes from  $H\alpha$  flux for  $z > 0.014$ .

### 3.6. Other continuum subtractions

We repeated the subtraction procedure for other narrow filters by applying a different method that was presented in Pascual et al. (2007) and Vilella-Rojo et al. (2015), in which three instead of two filters were used, one narrow filter that contain the feature of interest and two filters (one at each side if the main feature) to trace the continuum,

$$F_{cs} = \frac{(\bar{F}_{B1} - \bar{F}_{B2}) - \left(\frac{\alpha_{B1} - \alpha_{B2}}{\alpha_N - \alpha_{B2}}\right)(\bar{F}_N - \bar{F}_{B2})}{\beta_{B1} - \beta_N \left(\frac{\alpha_{B1} - \alpha_{B2}}{\alpha_N - \alpha_{B2}}\right)}, \quad (10)$$

where  $F_{cs}$  is the flux of the narrow filter with the continuum subtracted,  $\bar{F}_x$  is the average flux, and  $\beta$  and  $\alpha^5$  are defined as

$$\beta_x \equiv \frac{P_x(\lambda = \lambda_s)\lambda_s}{\int P_x(\lambda)\lambda d\lambda}; \quad \alpha_x \equiv \frac{\int P_x(\lambda)\lambda^2 d\lambda}{\int P_x(\lambda)\lambda d\lambda}. \quad (11)$$

The subindices  $B1$ ,  $B2$ , and  $N$  refer to the first broad filter, the second one, and the narrow filter, respectively. We note that in this case the continuum is not flat, but it has a slope defined by the two reference filters. Similarly to the method of subtracting the continuum described above, this method also improves significantly when additional filters are incorporated (Martínez-Solaache et al. 2021).

We used this method to subtract the continuum of the  $J0378$  and  $J0861$  filter using the uJAVA and gSDSS and the iSDSS and zSDSS filters, respectively. In the latter case, this is an estimate of the calcium triplet absorption. More negative NCR values for absorption filters indicate deeper absorption. Therefore, we inverted the sign of the resulting fluxes when constructing the NCR with this continuum-subtraction method to maintain coherence during analysis.

### 3.7. Error calculation

Throughout the preceding calculations, it is essential to consider the respective error estimation through robust propagation. The total uncertainty in our measurements of the flux (Eq. (1)) for a source and a given filter (Molino et al. 2014, Logroño-García et al. 2017) is given by

$$\sigma_{F_\lambda} = \sqrt{\left(\frac{\partial F_\lambda}{\partial ZP} \sigma_{ZP}\right)^2 + \left(\frac{\partial F_\lambda}{\partial C} \sigma_{bn}\right)^2 + \left(\frac{\partial F_\lambda}{\partial C} \sigma_{ec}\right)^2}, \quad (12)$$

$\sigma_{ZP}$  being the error of the zero point reflected in Table 2,  $\sigma_{ec}$  the uncertainty in the electron counting of the CCD, and  $\sigma_{bn}$  the large-scale background noise variation; the last two are given by

$$\sigma_{ec} = \sqrt{\frac{C}{G}}, \quad \sigma_{bn} = S_{\text{fit}} \sqrt{N_{\text{pix}}} (a_{\text{fit}} + b_{\text{fit}} \sqrt{N_{\text{pix}}}), \quad (13)$$

where  $G$  is the gain of the detector,  $N_{\text{pix}}$  the number of pixels; and  $S_{\text{fit}}$ ,  $a_{\text{fit}}$ , and  $b_{\text{fit}}$  are the resulting coefficients from the fitting. This information is found within headers of each image. We calculated the error of the other magnitudes by a quadratic propagation of errors. The error of the flux with the dust correction (Eq. (3)) is given by

$$\sigma_{F_i(\lambda)} = \sqrt{(D_{F_0} \cdot \sigma_{F_0(\lambda)})^2 + (D_k \cdot \sigma_{k(\lambda)})^2 + (D_E \cdot \sigma_{E(B-V)})^2}, \quad (14)$$

<sup>5</sup> We need to make a calculation that involves an integration of the transmission curve of the pass band; the data of the transmission curve are the same and can be found on the J-PLUS webpage.

where  $D_{F_0} = \frac{\partial F_i(\lambda)}{\partial F_0(\lambda)}$ ,  $D_{k'} = \frac{\partial F_i(\lambda)}{\partial k'(\lambda)}$ , and  $D_E = \frac{\partial F_i(\lambda)}{\partial E(B-V)}$ .  $\sigma_{E(B-V)}$  and  $\sigma_{k'(\lambda)}$  are the errors of the color excess (Eq. (5)) and the error given for the extinction law (Eq. (4)), respectively, that are given by

$$\sigma_{E(B-V)} = \sqrt{\left(\frac{\partial E(B-V)}{\partial i'} \sigma_{i'}\right)^2 + \left(\frac{\partial E(B-V)}{\partial g'} \sigma_{g'}\right)^2}, \quad (15)$$

where  $\sigma_{g'}$  and  $\sigma_{i'}$  are the error of the flux in the  $g_{\text{SDSS}}$  and  $i_{\text{SDSS}}$  bands given by Eq. (12), and

$$\sigma_{k'(\lambda)} = \sqrt{\left(\frac{\partial k'(\lambda)}{\partial \lambda} \sigma_{\lambda}\right)^2 + \left(\frac{\partial k'(\lambda)}{\partial R'_V} \sigma_{R'_V}\right)^2}. \quad (16)$$

Knowing the error of the fluxes with the dust correction, we can now calculate the error of the continuum subtracted fluxes. In the case of the subtraction of the red continuum of the  $J0660$  filter (Eq. (6)), we have

$$\sigma_{F_{\text{H}\alpha} + [\text{N II}]} = \sqrt{\left(\frac{\partial F_{\text{H}\alpha} + [\text{N II}]}{\partial \bar{F}_{F600}} \sigma_{\bar{F}_{F600}}\right)^2 + \left(\frac{\partial F_{\text{H}\alpha} + [\text{N II}]}{\partial \bar{F}_{r'}} \sigma_{\bar{F}_{r'}}\right)^2}, \quad (17)$$

with  $\sigma_{\bar{F}_{r'}}$  and  $\sigma_{\bar{F}_{F600}}$  being the error of the extinction-corrected flux of the  $J0660$  and  $r_{\text{SDSS}}$  filters obtained by Eq. (14).

We can then calculate the error of the  $\text{H}\alpha$  flux with the  $[\text{N II}]$  line subtracted (Eq. (8)) and the error of the  $[\text{N II}]$  line flux (Eq. (9)) as

$$\sigma_{F_{\text{H}\alpha}} = \frac{\partial F_{\text{H}\alpha}}{\partial F_{\text{H}\alpha} + [\text{N II}]} \sigma_{F_{\text{H}\alpha} + [\text{N II}]} \quad (18)$$

and

$$\sigma_{F_{[\text{N II}]}} = \sqrt{\left(\frac{\partial F_{[\text{N II}]}}{\partial F_{\text{H}\alpha}} \sigma_{F_{\text{H}\alpha}}\right)^2 + \left(\frac{\partial F_{[\text{N II}]}}{\partial F_{\text{H}\alpha} + [\text{N II}]} \sigma_{F_{\text{H}\alpha} + [\text{N II}]}\right)^2}. \quad (19)$$

Finally, we can calculate the error of the  $J0378$  and  $J0861$  flux with the continuum subtracted (Eq. (10)) as

$$\sigma_{F_{\text{cs}}} = \sqrt{\left(\frac{\partial F_{\text{cs}}}{\partial \bar{F}_N} \sigma_{\bar{F}_N}\right)^2 + \left(\frac{\partial F_{\text{cs}}}{\partial \bar{F}_{B1}} \sigma_{\bar{F}_{B1}}\right)^2 + \left(\frac{\partial F_{\text{cs}}}{\partial \bar{F}_{B2}} \sigma_{\bar{F}_{B2}}\right)^2}, \quad (20)$$

with the subindices  $B1$ ,  $B2$ , and  $N$  being referred to the first broad filter, the second one, and the narrow filter, respectively. Errors on  $\alpha_x$ ,  $\beta_x$ , and  $\Delta_x$  are negligible and are therefore not included in the calculation in order to simplify the procedure.

### 3.8. Final 3D data cubes

Once all these corrections are applied, we store these newly generated images as new slices in the 3D data cube, where fluxes are in the primary extension and flux errors are in the first extension. The final cube has 17 slices: 12 corresponding to the extinction-corrected images in the 12 J-PLUS filters, the  $\text{H}\alpha + [\text{N II}]$  continuum subtracted flux, the  $J0378$  and  $J0861$  continuum-subtracted flux, the flux corresponding only to the  $\text{H}\alpha$  line, and the  $[\text{N II}]$  emission line flux. These resulting data cubes are available at Zenodo<sup>6</sup> for all 418 SNe host galaxies presented in this work. Figure 3 shows the  $\text{H}\alpha$  image of the same galaxy from Fig. 2 at four different steps of the analysis: (a) the initial cutout from the observed image; (b) after dust reddening correction; (c) after  $r$ -band continuum subtraction; and (d) after  $[\text{N II}]$  emission removal.

<sup>6</sup> <https://zenodo.org/records/10514632>

## 4. Analysis

### 4.1. Normalized cumulative-rank distributions

The NCR is obtained by sorting the flux values in increasing order (see Fig. 4), constructing the cumulative distribution, and normalizing it to the total emission of the galaxy. This procedure associates each pixel with an NCR value between 0 and 1, where the brightest pixel in the galaxy has an NCR value of one, and all pixels corresponding to sky have an NCR value of zero. Then, the ranked value in this distribution corresponding to the pixel where the SN occurred<sup>7</sup> is the NCR associated with the SN. By compiling significant numbers of NCR values for different SN types, we can build NCR distributions and study the differences among types and use this information to infer properties of their progenitors.

The NCR method has previously been applied to  $\text{H}\alpha$  imaging of SN hosts. Assuming that the  $\text{H}\alpha$  emission scales by the number of stars that are formed (Kennicutt 1998), a diagonal cumulative NCR distribution with a mean value of 0.5 indicates that the population traces the observed light, and the SN type follows the number of stars formed and mapped by that particular SF tracer. Therefore, if SNe explode predominantly in locations with high SFR, they will favor higher NCRs. On the other hand, SNe that explode in random locations over the galaxy favor low NCRs.

Following James & Anderson (2006), we included all negative flux values in the construction of the NCR function; such values are the result of a background subtraction during the processing with the J-PLUS pipeline (Cenarro et al. 2019). This NCR calculation is applied independently to the 17 images of the 418 SN host galaxies.

### 4.2. SN environment SED fitting

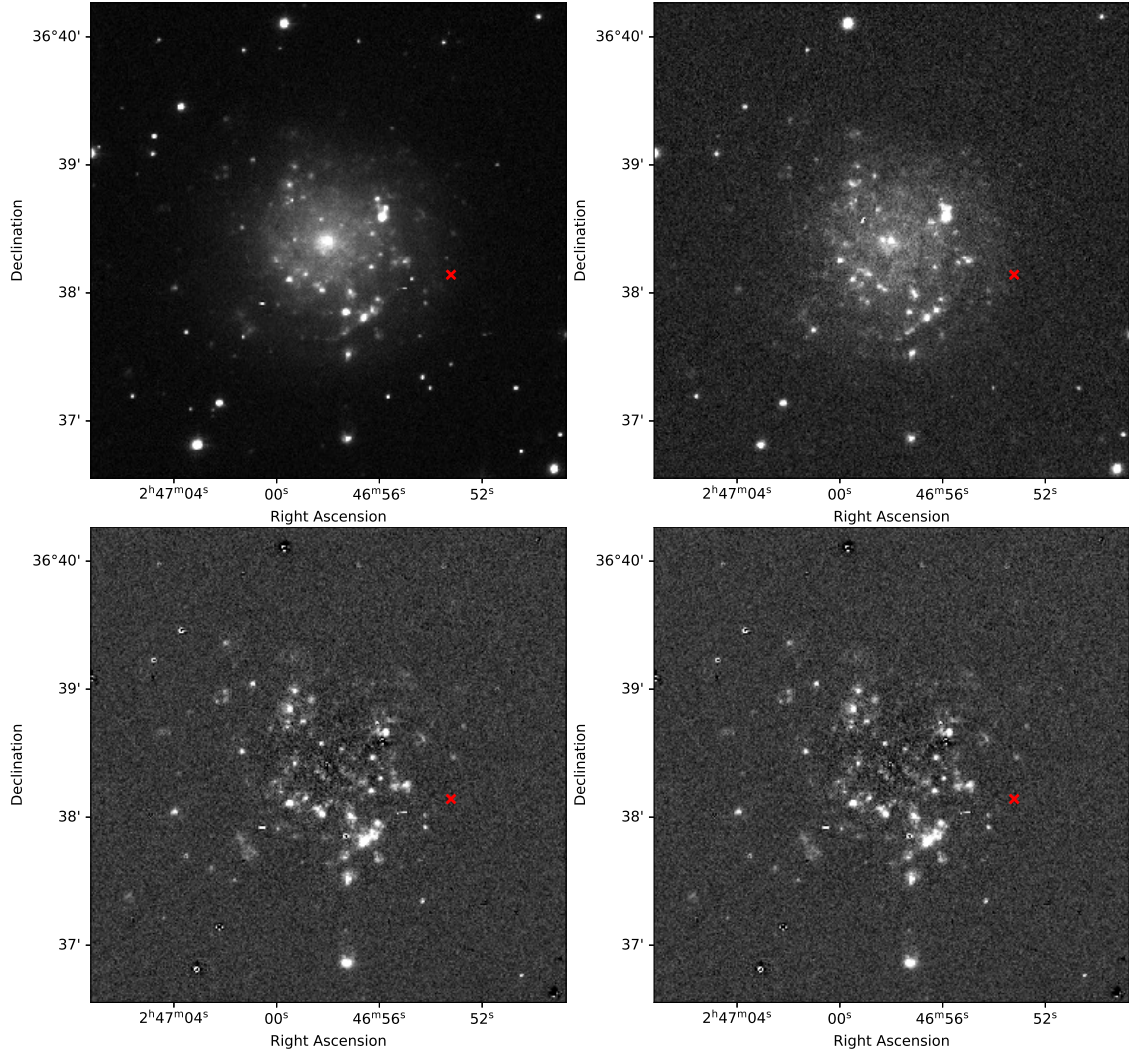
To extract the main properties of SN environments, we construct SEDs by measuring aperture photometry in all 12 J-PLUS bands in circular apertures of 1 kpc<sup>2</sup> centered at SN locations. To ensure the analysis yields meaningful results, the fluxes need to be only corrected for Milky Way extinction, so here we use the constructed 3D cubes with only this correction applied (at the stage outlined in Sect. 3.3). The size of the aperture is selected so that the flux contained in the aperture has a high enough signal-to-noise ratio to reliably obtain SSP parameters. We made calculations using the host galaxy redshift obtained from the OSC and assuming a flat  $\Lambda$ CDM cosmology with  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ . Within the redshift range of our sample this translates into apertures of a 26 to 2 arcsec radius. Figure 5 shows a couple of examples of the  $r_{\text{SDSS}}$  filter galaxy cutout and the SED from the 12 J-PLUS bands. We note how the SED shape of an SN environment with ongoing star formation (SN 1997dq on top) shows a bluer continuum with a clear  $\text{H}\alpha$  emission, while the SED of a SN environment in a passive galaxy (SN 1957B on the bottom) looks redder and smoother.

Once the 418 SEDs were built, we used FAST++ (Schreiber et al. 2018)<sup>8</sup>, a C++ implementation of FAST (Kriek et al. 2009), to fit them and obtain a number of galaxy physical properties. Given galaxy photometry, FAST++ determines the best-fit SED from a library of simple stellar population (SSP) models. These models are initially synthesized in composite stellar populations (CSPs) on a 5D grid, with each grid point

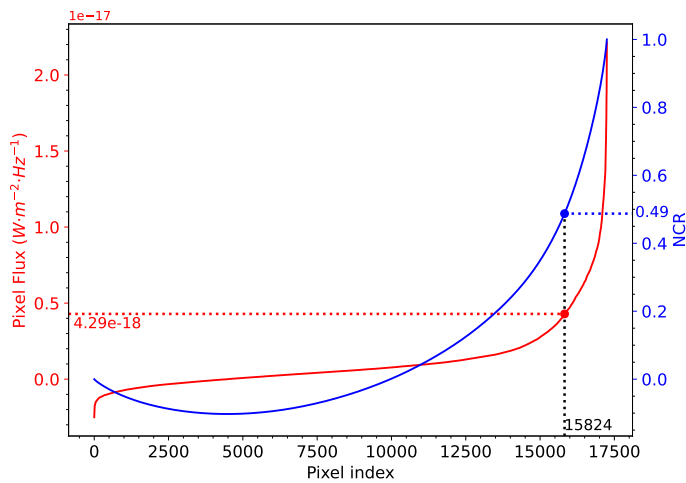
<sup>7</sup> The precision of SN locations is lower than the image resolution (see Sect. 1), so what we are measuring is the immediate SNe environment.

<sup>8</sup> Available at <https://github.com/cschreib/fastpp>





**Fig. 3.** J-PLUS NGC1058 J0660 image with the type II SN1961V marked in red. Top left: before dust extinction correction. Top right: after dust extinction correction. Bottom left: after continuum subtraction. Bottom right: after [N II] removal.



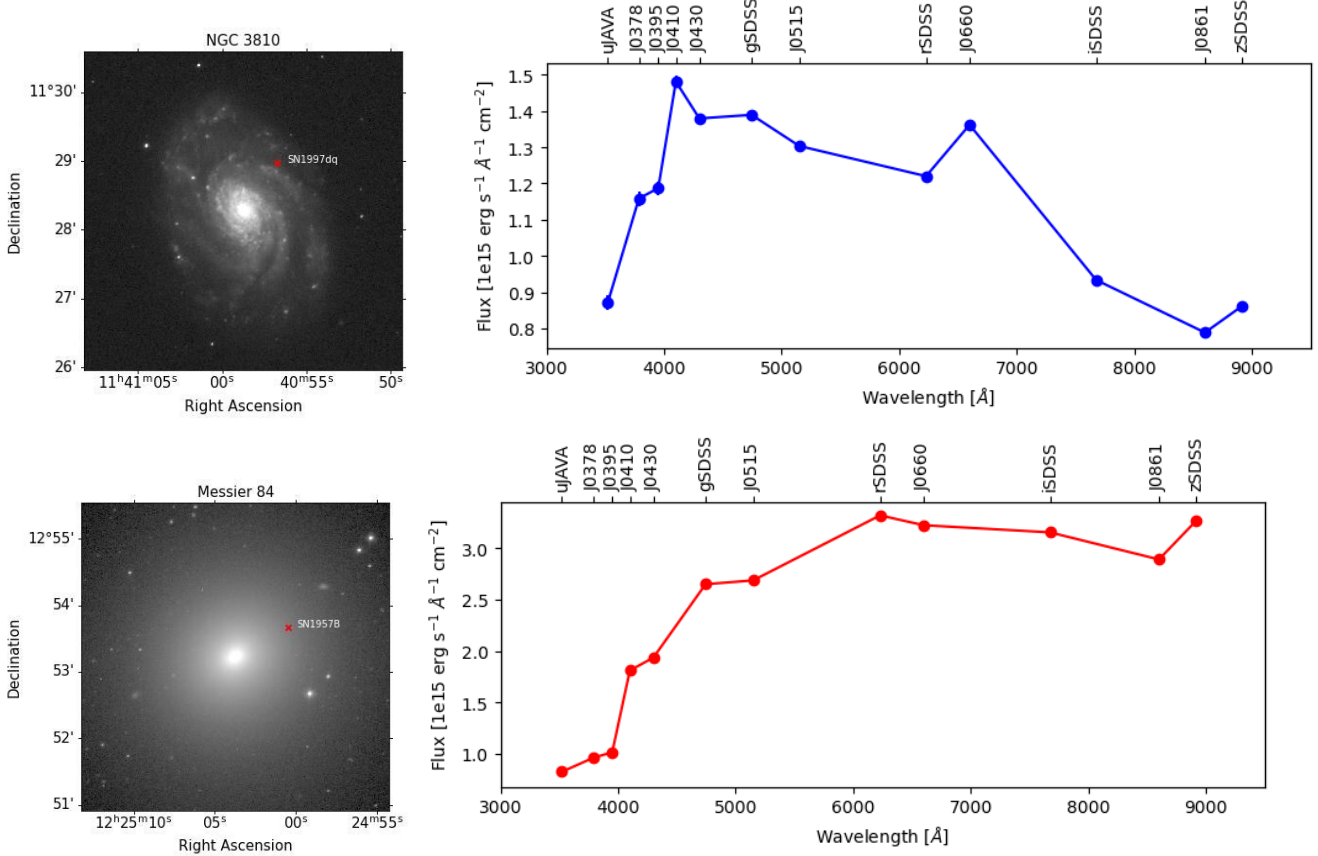
**Fig. 4.** Example of calculation of NCR value for SN pixel. The red curve is the sorted pixel value (left flux scale), while the blue curve is the cumulative distribution of the flux values, i.e., the normalized cumulative-rank pixel-value function (NCRPVF or NCR for simplicity). The blue dot represents the position and corresponding NCR value of the SN 2008az pixel, and the red dot represents the corresponding flux value.

corresponding to a CSP SED with age  $t_*$ , stellar metallicity  $Z_*$ ,  $V$  band extinction  $A_V$ , timescale  $\tau$ , and star formation history (SFH) at a given redshift  $z$ . At each point on the grid, a  $\chi^2$  value is calculated as

$$\chi_F^2 = \sum_i \frac{[F_{\lambda,i} - F_{\lambda,\text{mod}}(t_*, \tau, A_V, Z_*, z)]^2}{\sigma_i^2}, \quad (21)$$

where  $N$  is the number of photometric points and  $\sigma_i$  is the error for point  $i$ . One-sigma confidence intervals for parameters and derived quantities (such as stellar mass and SFR) are determined using Monte Carlo sampling of the grid around the lowest  $\chi^2$ . In this work, we used the BC03 (Bruzual & Charlot 2003) SSP library calculated using a delayed exponential (delayed- $\tau$ ) SFH parametrization and a Chabrier (2003) initial mass function. Dust extinction is modeled with a Calzetti et al. (2000) dust extinction law with a uniform and constant foreground dust screen. In all fits, the redshift was fixed to the reported spectroscopic redshift of the SN host galaxy.

The output of FAST++ includes the distribution of stellar populations of different ages and metallicities present at that location. Although several degeneracies are at play, with large enough samples one can disentangle the underlying differences



**Fig. 5.** Example of two SEDs of SN local environments of one squared kiloparsec extracted from J-PLUS observations. On top, the star-forming environment of SN 1997dq in its host galaxy NGC 3810 clearly shows a blue continuum based on the five broad-band filters and increased flux for the narrow-band filters where the main emission lines fall. On the bottom, for SN 1957B in the Messier 84 galaxy, a smooth red continuum is seen with no bumps at the narrow-band filter wavelengths.

between the average properties of these parent populations for different SN types. In particular, we mainly focused on stellar population age, dust extinction, stellar mass, and star formation rate.

## 5. Results

### 5.1. NCR distributions

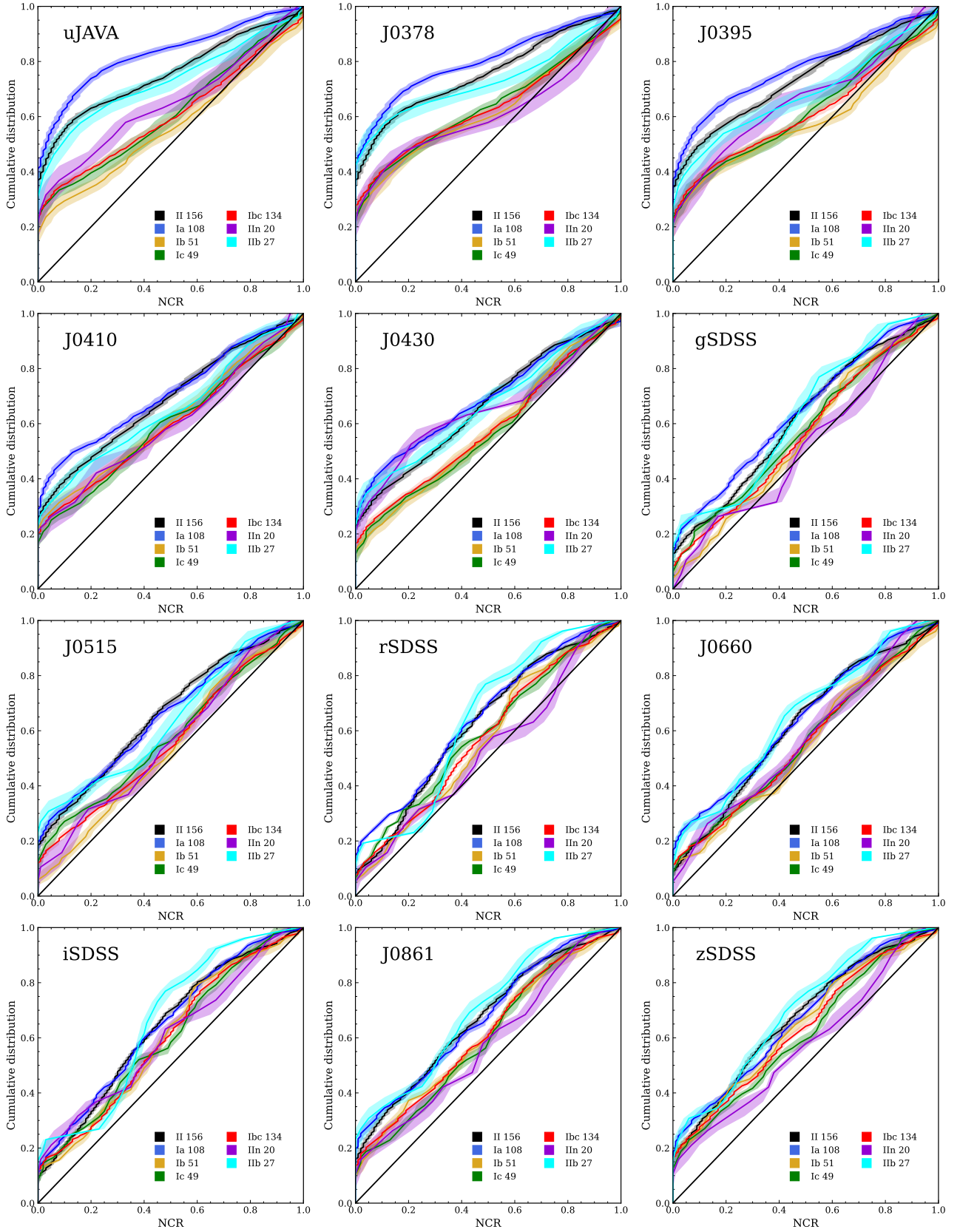
The resulting NCRs for all individual SNe are provided in Tables A.2–A.4. With the NCR and error for all SNe in our sample, we constructed cumulative distributions for all SN types and in all 17 frames of the data cubes as follows. We built a probability distribution for every NCR using 30 000 NCR trials of the form  $NCR + (\sigma_{NCR} \cdot \alpha)$ ,  $\alpha$  being a random variable between  $-1$  and  $1$ . We sorted all the NCRs in increasing order, and the median and one sigma of those distributions of 30 000 NCRs are shown in the plots. In Fig. 6, we show the 12 panels, one for each broad and narrow band, with the corresponding NCR distributions for the seven SN types.

In general, we observe that as we move to redder wavelengths (from gSDSS to zSDSS), all SN type distribution tends to be equal, following the black diagonal distribution, with an average  $NCR \gtrsim 0.4$  for every distribution. Type II, IIb, and Ia SN distributions begin to separate from the diagonal, and type Ibc, Ib, Ic, and IIn distributions begin to separate at bluer wavelengths. Also, type Ib SNe distribution gets closer to the diagonal as we move to bluer bands. The order Ia-II-IIb-IIn-Ibc remains

mostly equal in all filters; however, IIn and IIb types present higher and lower NCRs in red bands, respectively.

We conducted Kolmogorov–Smirnov (KS) tests to compare two samples and assess the probability of both being drawn from the same probability distribution. By setting a significance level of  $2\sigma$ , if the  $p$  value computed by the KS test falls below the threshold of 0.05, it means the rejection of the null hypothesis. In other words, it indicates they originated from different populations. Table 3 shows the  $p$  value for every pair of NCR distributions and the diagonal hypothetical distribution that accurately traces the respective flux ([1:1]), and also for every filter. The table summarizes what is shown in Fig. 6. For the 12 J-PLUS filters, all distributions for filters redder than J0410 show  $p$  values  $> 0.05$  (with the exception of the combination of II vs. Ibc types for J0515 and Ia vs. Ibc for J0410). For uJAVA, J0378, and J0395, Ia and II-IIb-IIn shows  $p$ -value  $> 0.05$ . Only combinations of CC-types with Ia, and combinations of II with Ib, Ic and Ibc presents  $p$ -values  $< 0.05$ , indicating that types Ia, II, IIb, and IIn come from different possible progenitors than type Ib, Ic, and Ibc SNe that are located in redder environments in comparison with the rest of the types.

The distributions of the five continuum-subtracted narrow-band filters are presented in Fig. 7. We note again that SNe with  $z > 0.014$  are not included in the NII distributions as explained in Sect. 3.5. The most significant characteristic is the high number of zero NCR values in all cases. As we discuss later (Sect. 6), we associated part of the zero values found with the small aperture of the JAST80 telescope and the exposure times



**Fig. 6.** Cumulative NCR distributions of 12 broad- and narrow-band filters. The straight black diagonal line represents a hypothetical distribution, infinite in size, which accurately traces the respective observed flux. Each SN distribution corresponds to the median and one sigma of the 30 000 NCR trial distributions as detailed in Sect. 5.1 of Lyman et al. (2021). It can be observed that as we move to redder filters (from gSDSS to zSDSS), the distributions for each type tend to become more similar along the diagonal. In contrast, there is a clearer distinction between them in the blue filters with more zero values.

**Table 3.**  $p$ -values of the Kolmogorov–Smirnov test for every SN type combination and filter.

Combination	uJAVA	J0378	J0395	J0410	J0430	gSDSS	J0515	rSDSS	J0660	iSDSS	J0861	zSDSS	H $\alpha$ + [N II]	[O II]	Ca II triplet	H $\alpha$	[N II]
II vs. Ia	0.213	0.740	0.709	0.757	0.646	0.896	0.999	0.442	0.408	0.997	0.974	0.973	0.076	0.013	0.990	0.399	0.527
II vs. Ib	0.002	0.096	0.023	0.638	0.392	0.386	0.083	0.279	0.169	0.496	0.632	0.925	0.095	0.347	0.158	0.037	0.017
II vs. Ic	0.021	0.164	0.184	0.438	0.251	0.810	0.318	0.777	0.298	0.421	0.367	0.391	0.017	0.461	0.001	0.001	0.002
II vs. Ibc	0.001	0.037	0.019	0.194	0.073	0.289	0.023	0.229	0.117	0.399	0.380	0.495	0.000	0.134	0.003	0.000	0.000
II vs. IIa	0.410	0.321	0.773	0.695	0.819	0.343	0.487	0.211	0.750	0.814	0.446	0.434	0.193	0.364	0.109	0.069	0.026
II vs. IIb	0.998	0.918	0.817	0.899	0.995	0.985	0.751	0.695	0.889	0.805	0.971	0.987	0.421	0.848	0.903	0.859	0.543
Ia vs. Ib	0.000	0.017	0.021	0.334	0.117	0.270	0.135	0.224	0.173	0.614	0.777	0.977	0.000	0.009	0.256	0.001	0.002
Ia vs. Ic	0.000	0.027	0.046	0.104	0.165	0.770	0.525	0.669	0.439	0.704	0.588	0.740	0.000	0.003	0.014	0.000	0.000
Ia vs. Ibc	0.000	0.003	0.006	0.046	0.074	0.217	0.092	0.342	0.217	0.506	0.697	0.959	0.000	0.000	0.026	0.000	0.000
Ia vs. IIa	0.086	0.184	0.362	0.449	0.926	0.155	0.595	0.271	0.774	0.832	0.557	0.557	0.006	0.076	0.263	0.014	0.010
Ia vs. IIb	0.627	0.837	0.700	0.938	0.999	0.837	0.893	0.700	0.999	0.771	0.996	0.999	0.018	0.043	0.965	0.310	0.121
Ib vs. Ic	0.988	0.999	0.926	0.999	1.000	0.879	0.701	0.890	0.999	0.994	0.991	0.984	0.990	1.000	0.409	0.380	0.895
Ib vs. Ibc	0.967	1.000	0.975	1.000	0.999	0.985	0.924	0.998	0.999	1.000	1.000	0.999	0.989	0.997	0.998	0.978	0.995
Ib vs. IIa	0.806	0.999	0.889	0.999	0.405	0.851	0.987	0.494	0.999	0.947	0.967	0.687	0.903	0.707	0.947	0.754	0.754
Ib vs. IIb	0.131	0.352	0.862	0.987	0.553	0.528	0.221	0.295	0.271	0.624	0.760	0.972	0.964	0.964	0.770	0.630	0.630
Ic vs. Ibc	1.000	1.000	0.999	1.000	1.000	0.999	0.999	0.991	1.000	0.999	0.999	0.999	0.997	0.999	0.744	0.827	0.995
Ic vs. IIa	0.934	0.997	0.966	0.999	0.469	0.731	0.975	0.789	0.999	0.995	0.951	0.951	0.857	0.707	0.984	0.681	0.883
Ic vs. IIb	0.286	0.349	0.897	0.835	0.620	0.793	0.865	0.512	0.438	0.280	0.488	0.512	0.798	0.964	0.118	0.179	0.335
Ibc vs. IIa	0.965	0.995	0.927	1.000	0.466	0.853	0.999	0.579	0.999	0.973	0.903	0.772	0.569	0.652	0.999	0.840	0.813
Ibc vs. IIb	0.264	0.393	0.927	0.957	0.732	0.673	0.613	0.393	0.346	0.443	0.673	0.788	0.810	0.880	0.372	0.314	0.487
IIa vs. IIb	0.890	0.688	0.996	0.996	0.996	0.551	0.749	0.284	0.565	0.714	0.376	0.578	0.549	0.873	0.498	0.553	0.553
II vs. [1:1]	0	0	0	0	0	0.00038	0	0	0	0	0	0	0	0	0	0	0
Ia vs. [1:1]	0	0	0	0	0	0.00125	0	0.00027	0.00028	0.00055	0	0	0	0	0	0	0
Ib vs. [1:1]	0.014	0.001	0.001	0.003	0.128	0.357	0.667	0.107	0.311	0.035	0.060	0.019	0.047	0.056	0.003	0.100	0.004
Ic vs. [1:1]	0.002	0.000	0.001	0.022	0.071	0.257	0.069	0.213	0.339	0.216	0.144	0.071	0.295	0.018	0.278	0.908	0.193
Ibc vs. [1:1]	0.000	0.000	0.000	0.000	0.000	0.146	0.016	0.031	0.036	0.004	0.002	0.000	0.011	0.001	0.000	0.053	0.000
IIa vs. [1:1]	0.028	0.021	0.036	0.146	0.016	0.855	0.643	0.952	0.568	0.366	0.496	0.496	0.396	0.473	0.317	0.711	0.554
IIb vs. [1:1]	0.000	0.000	0.001	0.016	0.005	0.107	0.017	0.020	0.068	0.015	0.023	0.023	0.037	0.017	0.000	0.007	0.001

**Notes.** We also included the  $p$ -values for the SN types and the diagonal hypothetical distribution that accurately traces the respective flux ([1:1]). Black font boxed values represent those distributions without an underlying causative relationship; i.e., we reject the null hypothesis that the two samples were drawn from the same probability distribution ( $p$ -value  $< 0.05$ ), and both NCR distributions come from different populations. The distribution used for the five continuum-subtracted filters have been rebuilt by proportionally removing the number of zero NCR values in each distribution that correspond to the lowest observed fraction of zeros in an individual distribution (see Sect. 6).

of the images used, which is not enough to collect enough photons to provide a high signal-to-noise ratio in these continuum-subtracted images for all galaxies. Moreover, H $\alpha$  is associated with stars of  $M > 15$ – $20 M_{\odot}$ , so many SN progenitors are not expected to be associated with this emission, contributing to more zero NCRs in the distributions. The removal of the [N II] line has no appreciable effect on the NCR distributions, with the plots of H $\alpha$ + [N II], H $\alpha$  and [N II] being similar. In any case, we observe a similar scenario between the broad- and narrow-band filters, where the type Ia distribution is clearly the one with the lowest NCRs in all cases, followed by type II. The [O II] emission also traces the star formation rate (Kennicutt 1998), similarly to H $\alpha$ . With the change of signs in Ca II triplet fluxes, a higher NCR after subtracting the continuum suggests deeper absorption, which is analogous to high NCR values in emission-line filters and indicates strong emission (see Sect. 3.6). The order of the distributions appears to remain consistent, with Ia, II, and IIb displaying lower NCR values, distinguishing them from the other SN types. However, the large amount of zero values makes the analysis difficult in these cases.

## 5.2. SN environment parameter distributions

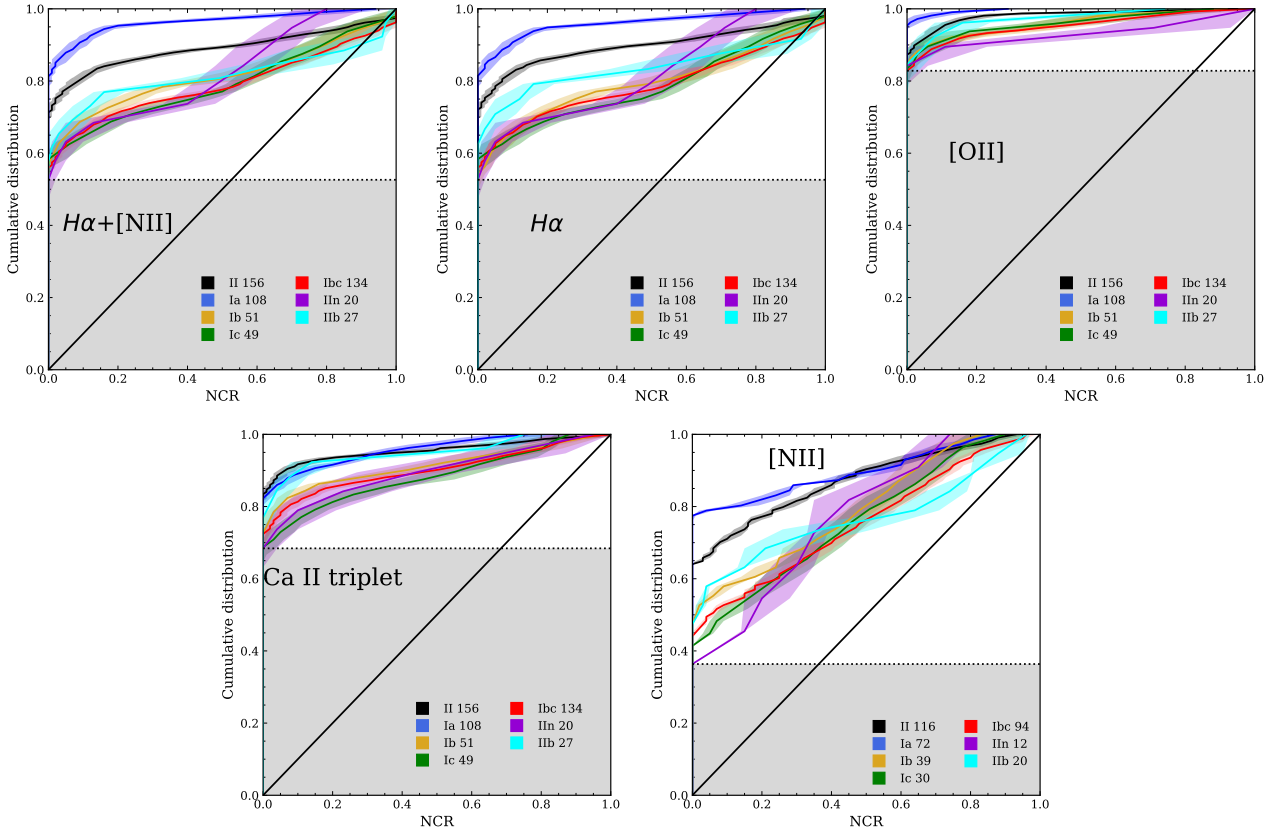
The local stellar population parameters  $A_V$ ,  $t_*$ , SFR, and sSFR for all 418 SN environments are reported in Table A.5. Additionally, their distributions are illustrated in Fig. 8 and segregated according to the same seven SN groups. These have been constructed similarly to in the previous section, that is, by performing 30 000 different distributions where in each realization the best output values from the best FAST++ fits are randomly

varied within their  $1\sigma$  uncertainty, and the median and the  $1\sigma$  variation of all realizations is taken as the final distribution.

We recover previous trends already found in the literature. The average local  $A_V$  is larger for SNe Ibc ( $1.13 \pm 0.08$  mag) compared to SNe Ia ( $0.63 \pm 0.09$  mag), with the SNII distribution being in the middle of the two ( $0.85 \pm 0.07$  mag), as previously reported by Galbany et al. (2017). This trend is driven by the SNII and SNIIc distributions, which are clearly shifted to higher values compared to all others. Similarly, all CC SNe distributions show younger average ages ( $38 \pm 10$  Myr for SNe II and  $43 \pm 13$  Myr for SNe Ibc) compared to those of SNe Ia environments ( $222 \pm 73$  Myr), and CC SNe environments have both higher SFR ( $0.08 \pm 0.06$  and  $0.16 \pm 0.04 M_{\odot} \text{ yr}^{-1}$  for SNe Ibc and SNe II, respectively) and sSFR ( $1.88 \pm 1.40$  and  $4.35 \pm 1.29 \times 10^{-8} \text{ yr}^{-1}$  for SNe Ibc and SNe II, respectively) than SNe Ia ( $0.01 \pm 0.01 M_{\odot} \text{ yr}^{-1}$  and  $0.19 \pm 0.16 \times 10^{-8} \text{ yr}^{-1}$ ) that occur in more passive local environments (Galbany et al. 2018).

## 5.3. Radial distributions of the SNe

To investigate potential correlations between the emission of a specific band and the distribution of SNe in galaxies, we calculated two parameters: the normalized galactocentric distance (NGCD) and the fraction of the  $r$ -band emission flux (Fr; Anderson & James 2009; Haberman et al. 2010). Our approach involved fitting an elliptical 2D Gaussian model to the rSDSS images of the galaxies to determine the shape and size of each galaxy. If the SN was located at the center of the ellipse, NGCD was assigned a value of 0; if the SN was positioned at the outer ellipse's edge, the NGCD equaled 1. We then fit a second ellipse,



**Fig. 7.** Cumulative NCR distributions of five constructed filters.  $H\alpha+[NII]$ ,  $[OII]$  and Ca II triplet distributions correspond to the  $J0660$ ,  $J0378$ , and  $J0861$  continuum-subtracted filters, respectively. The  $H\alpha$  filter is the result of removing the  $[NII]$  contribution, which is only applied to SNe with  $z < 0.014$ . All distributions correspond to the median and one sigma of the 30 000 NCR trial distributions as outlined in Sect. 5.1 of Lyman et al. (2021). The straight black diagonal line represents a hypothetical distribution, infinite in size, that accurately traces the respective observed flux. The gray shading marks the lowest values of zeros in a distribution in each filter.

with the same parameters as the first one, but scaled to place the SN at the edge of this second ellipse.  $Fr$  was subsequently computed as the ratio between the integrated fluxes in the rSDSS band of the second and first ellipses.

We find in Fig. 9 that all SN types exhibit nearly identical distributions with respect to each other, except for type IIb SNe, which appear to be concentrated in the inner regions of galaxies (within the range of 0.2 to 0.6 in NGCD). For the other SN types, there does not seem to be a preference for any specific location within the galaxies.

The  $Fr$  are also similar among all SN types, and close to the diagonal, tracing the rSDSS emission as the NCR distributions of Fig. 6 do. Since all J-PLUS bands present the NCR distributions close to the diagonal, and there is no preference in the location of the SN, we only traced the general shape and size of the galaxies.

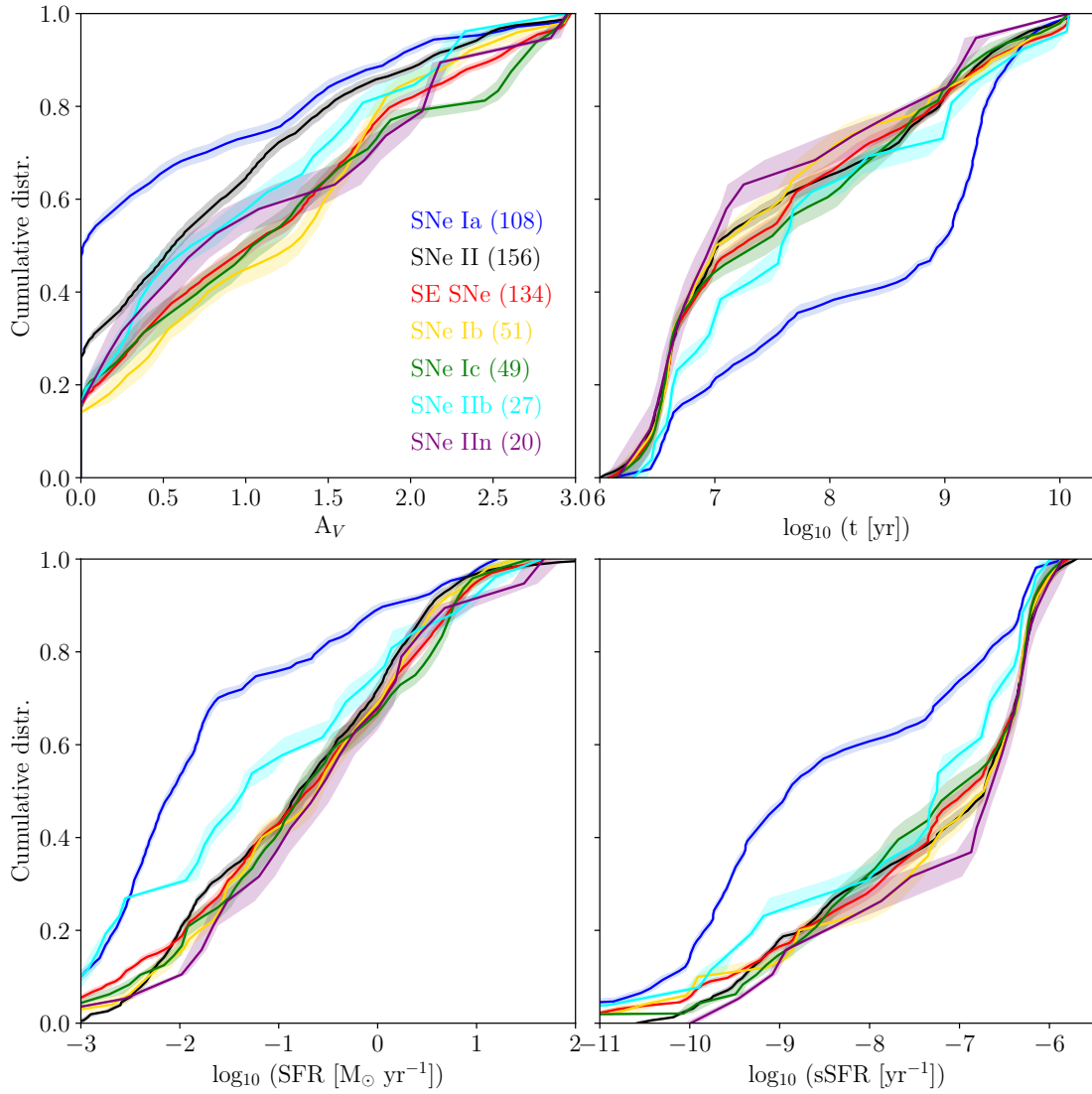
## 6. Discussion

### 6.1. NCR results

Concerning the NCR distributions of the broad- and narrow-band filters and the average NCR values, it is worth noting that the broad filters exhibit slightly lower errors than the narrow-band filters, particularly for those where the continuum has been subtracted. The errors for the uJAVA filter are higher than the rest of the broad-band filters. This is what we expected since the broad filters have higher S/Ns, and the uJAVA filter has the

lower FWHM of all J-PLUS broad-band filters. We also observe approximately the same sequence between the main types in all panels in Fig. 6; this is type Ibc between II and Ia for bluer wavelengths, moving upward in the cumulative plot and downward for NCRs for redder wavelengths. This shift of the Ibc distribution is due to the Ib contribution, which correlates more with bluer bands, similarly to what was reported by Kangas et al. (2013). Kelly & Kirshner (2012) also found that the environments of SN Ib exists preferentially in environments with bluer surface brightness, especially in the ultraviolet band. In our case, we also found that the filter that best distinguishes the type Ib distribution is uJAVA.

By looking at continuum-subtracted filters, the main characteristic is the large amount of zeros present in all distributions. For instance, in the  $H\alpha$  plot we have a noticeable amount of zero NCR values:  $\sim 35\%$  for Ibc,  $\sim 53\%$  for II, and  $\sim 65\%$  for Ia. This is partly because many SNe do not follow the  $H\alpha$  emission. However, this fraction of zeros is larger than in previous studies (e.g., Anderson et al. 2012 found  $\sim 25\%$  for Ibc,  $\sim 39\%$  for II, and  $\sim 58\%$  for Ia). Since these filters represent the residual flux once a pseudo-continuum (from broad-band filters) has been subtracted from a narrow-band filter, and since the diagonal line assumes a population that specifically traces the observed light, the distribution will include a fraction of zeros that depends on the image depth. Therefore, we have two possible explanations for this large amount of zeros; either most of the flux collected by the narrow-band filters is due to the underlying continuum, or the combination of a small telescope aperture (0.8 m) and low



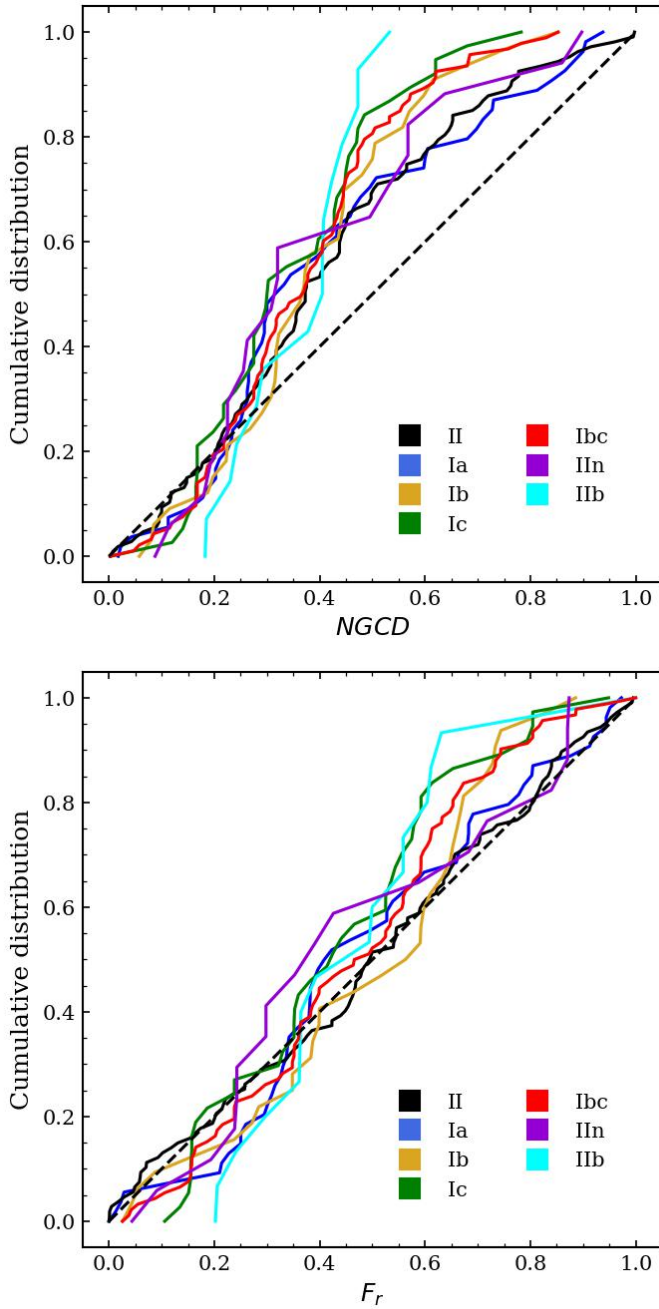
**Fig. 8.** Cumulative distributions of each of three main SN types for main environmental parameters obtained from best FAST++ fits to SED: extinction  $A_V$ , average stellar age  $\langle t \rangle$ , star-formation rate  $\log_{10}$  SFR, and specific star-formation rate  $\log_{10}$  SFR/Mass.

exposure time was not enough to get significant contrast in these continuum-subtracted filters.

To further analyze this, we rebuilt the NCR distributions, this time proportionally removing the number of zero NCR values in each distribution that correspond to the lowest observed fraction of zeros in an individual distribution, following a procedure similar to Ransome et al. (2022). This can be done as all our images are of a similar depth, since all the images were obtained using the same telescope with the same conditions, exposure time, and instrument configuration. This is equivalent to excluding the shaded region in all panels of Fig. 7, because we assume we have reached the background level at that point, and tracing the diagonal from the intersection of the dotted horizontal line with the zero NCR to the upper right corner. This assumption is grounded in the fact that we successfully retrieved a certain level of signal, and anything below that threshold corresponds to the background we were not able to recover due to the combination of the telescope and exposure time used, as previously mentioned. In Fig. 10, we show the resulting distributions.  $H\alpha + [N II]$ ,  $H\alpha$ , and  $[N II]$  plots still show similar distributions; the main type that correlates the emission is Ic, with the high-

est NCR values on average (Fig. 11). For the  $J0660$  (Fig. 6) and  $H\alpha$  (Fig. 10) panels for type Ia, II, and all Ibcs combined, we recovered the tendency of the distributions seen in previous works (e.g., Anderson & James 2008), where type Ic/Ibc SNe is the closest distribution to the diagonal line, followed by type II and type Ia SNe. We interpreted this as a sequence in progenitor age and then inferred a sequence of progenitor mass, with the distributions closer to the diagonal being the SN type with higher mass progenitors (Kuncarayakti et al. 2018).

Type IIc SN distribution also correlates with the  $H\alpha$  emission, with lower NCR values than Ic ( $\sim 0.05$  on average), which is also consistent with the literature (Ransome et al. 2022), with the same methodology), however, our nonzero IIc sample is reduced to only seven SNe. All CC SNe types are strongly correlated to the  $[O II]$  emission, that also traces SFR (Kennicutt 1998) following the sequence of Ia-II-Ib-Ic-Ibc-Ib-IIc. If we exclude the bias in the IIc type, which is limited to only four SNe and not numerous enough to perform reliable statistics, we can observe that type Ib exhibits the most pronounced correlation with this emission, primarily because this type of SN is more closely associated with bluer bands, as previously mentioned. The same



**Fig. 9.** Cumulative distributions of normalized galactocentric distances (top) and  $F_r$  (bottom) for every SN type.

correlations between SN types are found for the  $H\alpha + [N II]$ ,  $H\alpha$ , and  $[N II]$  filters. Besides, the Ia, IIb, and IIIn types are correlated when the  $[N II]$  is subtracted, the  $p$ -value being  $> 0.05$  after the removal.

The Ca II triplet distribution makes a clear division between types II, IIb, and Ia and types Ib, Ic, and IIIn (closer to the diagonal because of the change in the sign). These differences between populations are also evident when computing the  $p$  value of the KS test (Table 3): II versus Ic, II versus Ibc, Ia versus Ic, and Ia versus Ibc present  $p$  values  $< 0.05$  for the Ca II triplet NCRs; II versus Ia, II versus IIb, and Ia versus IIb present  $p$  values  $> 0.05$ , as well as the rest of CC SNe. The interpretation of the intensity of the Ca II triplet absorption is not trivial. Some works (e.g., Díaz et al. 1989; García-Vargas et al. 1998) have shown

that the Ca II triplet equivalent width (EW) presents a degeneracy between stellar population age and metallicity. At high metallicity, the Ca II EW reaches its maximum value for stellar populations of  $\sim 10$  Myr, primarily attributed to the presence of red supergiant (RSG) stars. A secondary, lower peak emerges around 100 Myr, driven by stars in the asymptotic giant branch (AGB) phase. Furthermore, for ages over 1 Gyr, the Ca II EW becomes a reliable indicator of increasing metallicity. Based on our results, and compared to SNIi that explode in low Ca II NCR environments, one possible interpretation would be to attribute the higher values for Ib/Ic to young progenitors exploding at locations with metal-rich older populations. For SNIIn, although with low numbers, our results are in line with these progenitors being a combination of very young ( $\sim 10$  Myr peak) and older ( $\sim 100$  Myr peak) stars (Galbany et al. 2018).

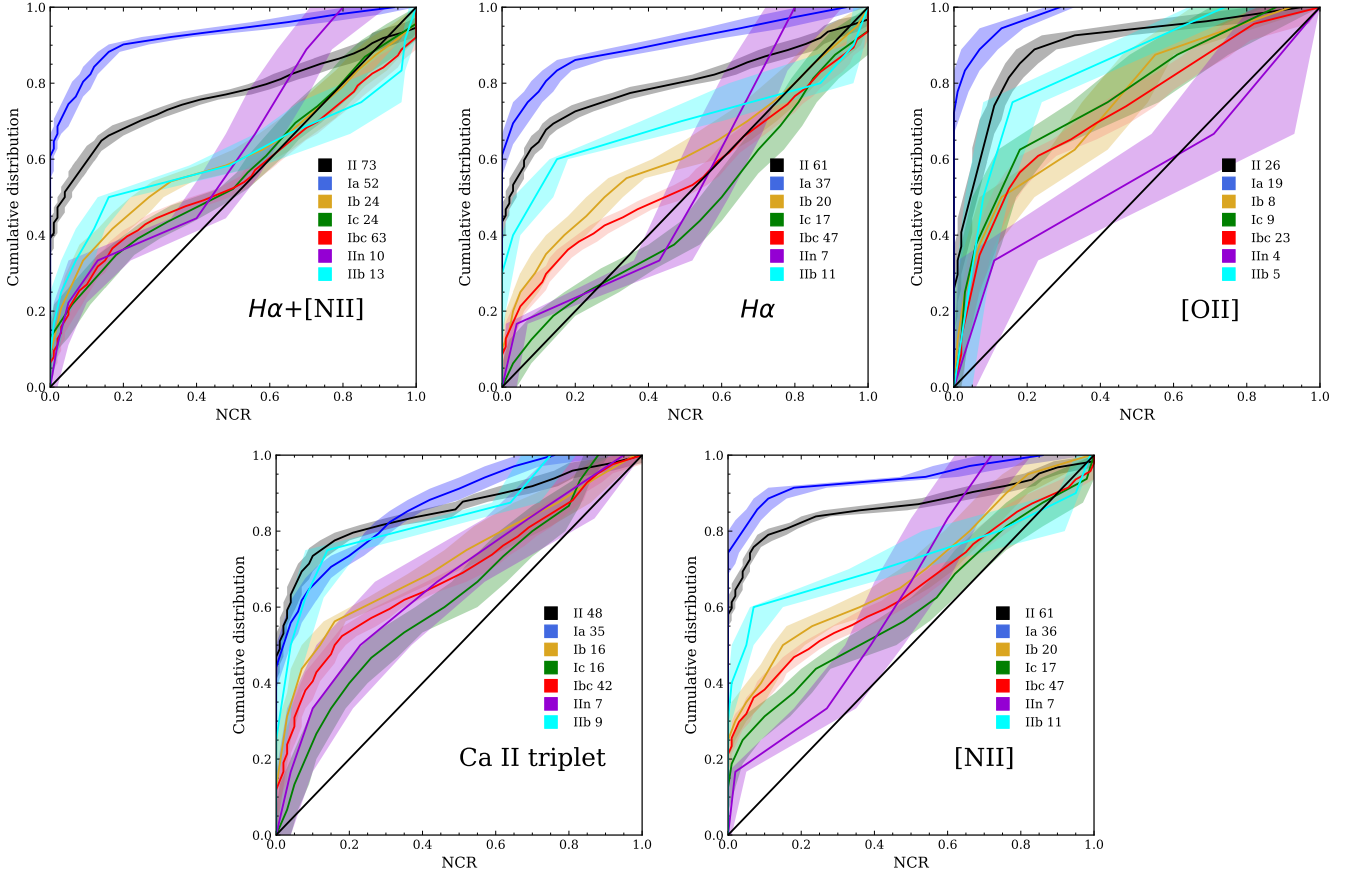
Finally, in Fig. 11 we show the average NCR values in all filters for the 7 SN types. On the left panel, we see that on average, through the 12 J-PLUS filters, SNe Ia present the lower NCR values (between  $0.10 \pm 0.02$  and  $0.33 \pm 0.02$ ), followed by  $II < IIb < IIIn < Ibc$ , which are the highest NCR values for Ibc-Ib-Ic types, ranging between  $0.36 \pm 0.01$  and  $0.48 \pm 0.01$ . The IIIn oscillates among the Ia-II-IIb types and Ibc, obtaining values from  $0.35 \pm 0.02$  up to  $0.49 \pm 0.01$  in the gSDSS and rSDSS bands. Moreover, all NCRs of the Ia, II, and IIb types tend to increase as we move to redder wavelengths (up to 0.2 in NCR), while Ib, Ic, and Ibc tend to remain constant.

For the five continuum-subtracted filters, shown in the right panel of Fig. 11, Ic SNe tend to have higher NCR values on average in comparison with the corresponding J-PLUS bands, the average NCR being, for instance,  $0.39 \pm 0.02$  in the  $J0660$  filter and  $0.54 \pm 0.05$  in  $H\alpha$ , and  $\sim 0.5$  for the rest, except the  $[O II]$ , whose average NCR falls to  $0.26 \pm 0.02$ . Moreover, the figure illustrates the low NCR values on average for Ia SNe in comparison with the rest of the SN types; these are especially low for  $[O II]$ . The insignificant effect of the  $[N II]$  removal from the  $J0660$  filter for the rest of the SN types is also evident since the average NCR values for  $H\alpha$  and  $H\alpha + [N II]$  are similar in all distributions (except in type Ic, where the difference is  $\sim 0.1$ ). The IIIn type is the next in terms of the highest NCRs.  $H\alpha$  presents a value of  $0.45 \pm 0.11$ , which is consistent with the  $0.521 \pm 0.038$  that Ransome et al. (2022) found; however, we need to take into account that we only have seven type IIIn SNe in our nonzero sample. The ordination of SN types between  $J0378$  and  $[O II]$  and between  $J0861$  and Ca II triplet remains the same, but with systematically lower NCRs after subtracting the continuum.

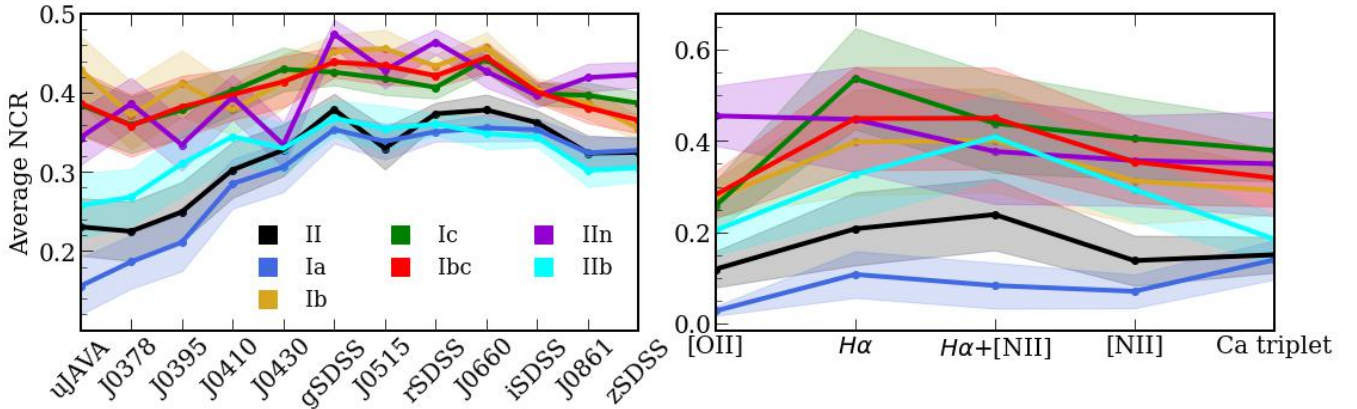
## 6.2. SED fitting

We recovered the same relations between the main SN types previously found in other works. Although the main focus of this work was to construct NCR distributions of the narrow-band filters, the fact that we recover these relations proves that the data is good enough to perform this kind of study. However, to our knowledge, this is the first time that SED fitting is performed adding these seven narrow-band filters on top of the five main SDSS broad bands.

To see how much information these narrow-band filters carry that is not included in the broad bands, we repeated the FAST++ fits this time only using the five broad-band filters. In Fig. 12, we show the differences between the twelve- and five-filter SED fits for the four main environmental parameters. Individual dots show the difference for each individual SN environment, the



**Fig. 10.** Same as Fig. 7, but rebuilding the NCR distributions proportionally removing the number of zero NCR values in each distribution that correspond to the lowest observed fraction of zeros in an individual distribution.



**Fig. 11.** Average NCR distributions as a function of the wavelength, for every SN type. The distributions are built averaging the NCR values for every J-PLUS photometric filter from the bluer to the redder filter (left panel). The right plot is built as the left one, but using the continuum-subtracted filters and proportionally removing the number of zero NCR values in each distribution that correspond to the lowest fraction of zeros in a distribution in each filter, as described in Sect. 6.1.

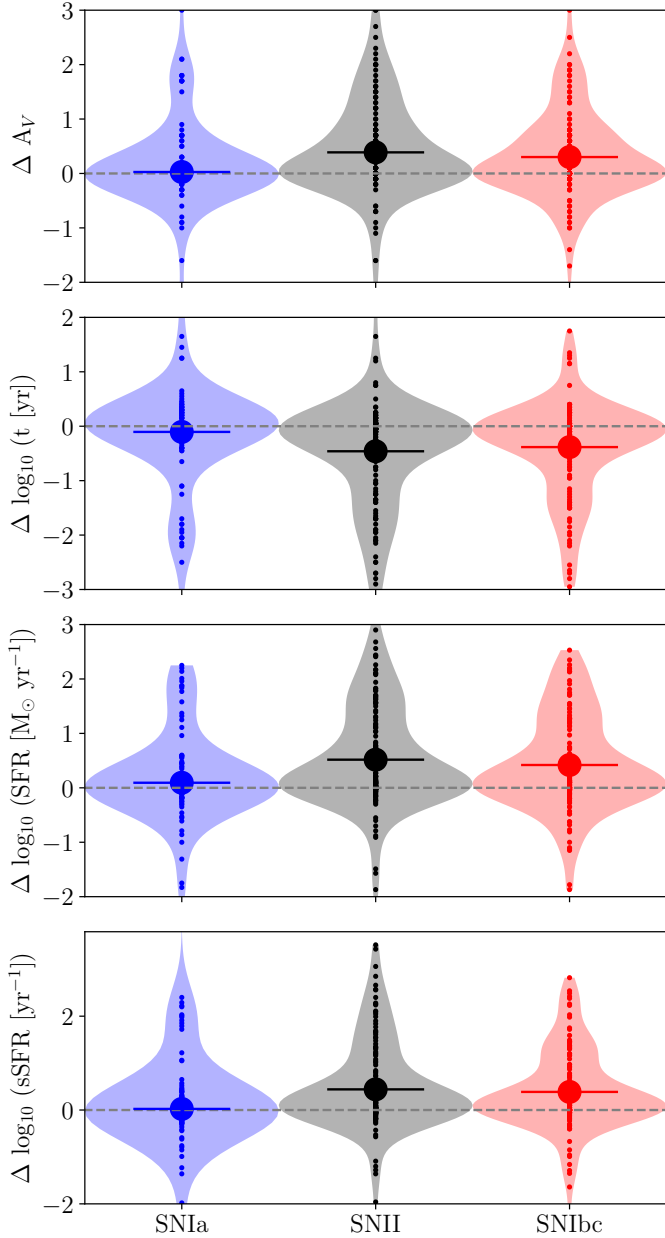
big dot represents the average of all differences, and the violin plots represent the shape of the distributions. While all violins show the most populated value, similar to the median, compatible with zero difference, the average values show small shifts, which are more significant for the two CC SNe types with respect to Ia SNe. Given that the narrow filters are chosen so that they cover some of the main emission lines in galaxy spectra, these shifts for CC SNe are expected because they occur next to HII regions that have intense emission lines. On the other hand, SNe Ia occur in all environments, most of them being

non-star-forming without emission lines, so the narrow-band filters should not change the SED and the results of the fitting significantly.

## 7. Conclusions

We present a study of 418 environments of different supernova types using broad- and narrow-band images obtained by the J-PLUS collaboration with the JAST80 telescope. We measured the NCR value for each SN and constructed NCR distributions





**Fig. 12.** Difference between FAST++ SED fits once all 12 J-PLUS filters are used and those only using the five broad-band SDSS filters, for the four main environmental parameters. Individual dots represent the difference of each parameter for each individual SN environment, while the big dot represents the average of all differences. Violin plots represent the shape of the distributions.

by SN type. Improvements over previous works include the extension of the original  $H\alpha$  NCR analysis to other optical bands, the homogeneity of the data of those 418 environments using the same telescope, instrument configuration, spatial resolution, and image depth, and a better handling of the extinction correction and the stellar-continuum subtraction for the narrow-band filters. In addition, we constructed SEDs of all these nearby environments and performed SSP synthesis to estimate parameters such as star formation rate, stellar age, and extinction.

Regarding the NCR analysis on the J-PLUS filters, all broad-band distributions are quite similar, all of them being close to the diagonal. The fact that almost all broad-band distributions are close to the diagonal, along with the absence of a clear loca-

tion preference for the SNe in the galaxies with respect to both NGCD and the fraction of rSDSS flux (with the exception of the IIb type, which occur preferably in the inner part of the galaxies, between 0.2 and 0.5), suggests that we are primarily tracing the overall shape of the galaxy. The bluer bands (uJAVA, J0378, and J0395) show the only difference with respect to the others; Ia, II, and IIb types present lower NCRs on average. When performing KS tests, the p values indicate that the underlying populations of Ia-II-IIb-IIc and Ibc-Ib-Ic are different, indicating that Ia-II-IIb progenitors trace the redder wavelengths better, also shown in the average NCR distributions.

The average NCR distributions for the five continuum-subtracted filters show a large fraction of zero values. For instance, in the  $H\alpha$  emission we have  $\sim 35\%$  for Ibc,  $\sim 53\%$  for II, and  $\sim 65\%$  for Ia. This is partly because many SNe do not follow the  $H\alpha$  emission. However, it is worth noting that there are more instances of zero values compared to previous studies (Anderson et al. 2012). We attribute those extra zero values to the low signal-to-noise ratio in the narrow-band filters due to the combination of a small telescope aperture (0.8 m) and low exposure time not being sufficient to obtain significant contrast in these continuum-subtracted filters. For this reason, the fraction of zero values is not directly comparable with other studies that use images with different depths.

After considering the lowest percent of zeros among the seven SN type distributions as our background, and proportionally removing the number of zero NCR values in each distribution that correspond to the lowest fraction of zeros in a distribution in each filter, we constructed new background-free NCR distributions, as performed by Ransome et al. (2022). The SN Ic distribution correlates most strongly with  $H\alpha$  emission. It is followed by IIc, Ibc, Ib, IIb, II, and Ia, showing that the progenitors of the Ic types are younger and more massive than type IIs, while Ia SNe progenitors are even older, recovering results shown by previous works. However, it is important to note that the distribution for type IIc is based on only seven SNe in this panel. All core-collapse types strongly correlate with the [O II] emission, which also traces SFR following the sequence of Ia-II-IIb-Ic-Ibc-Ib (excluding IIc since we have only 4 SNe). Besides this, the Ca II triplet makes a clear division between II-IIb-Ia and Ib-Ic-IIc progenitors, which we interpret as a difference in environmental stellar metallicity. The KS test performed for the five continuum-subtracted filters also shows that II versus Ic, II versus Ibc, Ia versus Ic, and Ia versus Ibc come from different populations. The opposite occurs with II versus Ia, II versus IIb, and Ia versus IIb, as well as the rest of CC SNe. In addition, after the [N II] removal, the distinction between the underlying populations of the Ia, IIb, and IIc types, as revealed in the KS test, disappears, but the NCRs remain unchanged. We should note that the removal of [N II] is oversimplified and may potentially introduce noise into this aspect of the analysis.

Regarding the SED analysis, we recovered the main relations previously reported among SN types for all parameters. Moreover, we found that including the J-PLUS narrow filters in the SED fitting has a more significant effect for the CC SNe environmental parameters, shifting all parameters due to the presence of strong emission lines in their typical star-forming regions.

In addition to the results that were known from the literature and that this paper confirms, the new findings presented can be summarized as follows:

- (i) NCR analysis for broad-band filters traces the overall shape of the galaxy.
- (ii) IIb type SNe occur preferably in the inner part of the galaxies.

- (iii) Ia, II, and IIb types occur preferably in redder environments.
- (iv) CC SNe strongly correlate the [OII] emission, tracing the SFR.
- (v) High Ca II NCR environments can be interpreted as young progenitors exploding at locations with metal-rich older populations.
- (vi) IIcn progenitors are a combination of older and very young stars.
- (vii) Including narrow filters in the SED analysis causes a shift in the four main environmental parameters for the CC SNe. The JAST80 telescope with J-PLUS filters has allowed us to study SN host galaxies in the northern sky up to  $z = 0.0163$ . The H $\alpha$  emission from galaxies at higher redshifts falls outside the width covered by the J0660 filter. In the near future, using the next generation of narrow-band wide-field surveys, such as J-PAS at the JST250 telescope (Benítez et al. 2014) and the PAU survey at the *William Herschel* Telescope (Padilla et al. 2019), with filters covering a contiguous range at redder wavelengths, we will be able to extend the redshift range by using the filter covering the observed H $\alpha$  emission at  $6563 \times (1 + z)$  Å.

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## References

- Anderson, J. P., & James, P. A. 2008, *MNRAS*, **390**, 1527  
 Anderson, J. P., & James, P. A. 2009, *MNRAS*, **399**, 559  
 Anderson, J. P., Haberman, S. M., James, P. A., & Hamuy, M. 2012, *MNRAS*, **424**, 1372  
 Anderson, J. P., James, P. A., Haberman, S. M., Galbany, L., & Kuncarayakti, H. 2015, *PASA*, **32**, e019  
 Arnett, W. D., Schramm, D. N., & Truran, J. W. 1989, *ApJ*, **339**, L25  
 Astier, P., Guy, J., Regnault, N., et al. 2006, *A&A*, **447**, 31  
 Benítez, N., Dupke, R., Moles, M., et al. 2014, arXiv e-prints [arXiv:1403.5237]  
 Bruzual, G., & Charlot, S. 2003, *MNRAS*, **344**, 1000  
 Calzetti, D., Armus, L., Bohlin, R. C., et al. 2000, *ApJ*, **533**, 682  
 Cenarro, A. J., Moles, M., Marín-Franch, A., et al. 2014, *Proc. SPIE*, **9149**, 914911  
 Cenarro, A. J., Moles, M., Cristóbal-Hornillos, D., et al. 2019, *A&A*, **622**, A176  
 Chabrier, G. 2003, *PASP*, **115**, 763  
 Cid, Fernandes R., Mateus, A., Sodr e, L., Stasińska, G., & Gomes, J. M. 2005, *MNRAS*, **358**, 363  
 Delchambre, L., Bailer-Jones, C. A. L., Bellas-Velidis, I., et al. 2023, *A&A*, **674**, A31  
 Diaz, A. I., Terlevich, E., & Terlevich, R. 1989, *MNRAS*, **239**, 325  
 Eldridge, J. J., & Maund, J. R. 2016, *MNRAS*, **461**, L117  
 Filippenko, A. V. 1988, *AJ*, **96**, 1941  
 Filippenko, A. V. 1997, *ARA&A*, **35**, 309  
 Fruchter, A. S., Levan, A. J., Strolger, L., et al. 2006, *Nature*, **441**, 463  
 Galbany, L., Stanishev, V., Mourão, A. M., et al. 2014, *A&A*, **572**, A38  
 Galbany, L., Stanishev, V., Mourão, A. M., et al. 2016, *A&A*, **591**, A48  
 Galbany, L., Mora, L., González-Gaitán, S., et al. 2017, *MNRAS*, **468**, 628  
 Galbany, L., Anderson, J. P., Sánchez, S. F., et al. 2018, *ApJ*, **855**, 107  
 Gal-Yam, A. 2017, in *Handbook of Supernovae*, eds. A. W. Alsabti, & P. Murdin, 195  
 Garcia-Vargas, M. L., Molla, M., & Bressan, A. 1998, *A&AS*, **130**, 513  
 Graur, O., Bianco, F. B., Modjaz, M., et al. 2017, *ApJ*, **837**, 121  
 Green, G. 2018, *J. Open Source Softw.*, **3**, 695  
 Guillochon, J., Parrent, J., Kelley, L. Z., & Margutti, R. 2017, *ApJ*, **835**, 64  
 Haberman, S. M., Anderson, J. P., & James, P. A. 2010, *ApJ*, **717**, 342  
 Han, D.-H., Park, C., Choi, Y.-Y., & Park, M.-G. 2010, *ApJ*, **724**, 502  
 Hoyle, F., & Fowler, W. A. 1960, *ApJ*, **132**, 565  
 James, P. A., & Anderson, J. P. 2006, *A&A*, **453**, 57  
 Kangas, T., Mattila, S., Kankare, E., et al. 2013, *MNRAS*, **436**, 3464  
 Kangas, T., Portinari, L., Mattila, S., et al. 2017, *A&A*, **597**, A92  
 Kelly, P. L., & Kirshner, R. P. 2012, *ApJ*, **759**, 107  
 Kennicutt, R. C. Jr. 1998, *ApJ*, **498**, 541  
 Kriek, M., van Dokkum, P. G., Labbé, I., et al. 2009, *ApJ*, **700**, 221  
 Kuncarayakti, H., Doi, M., Aldering, G., et al. 2013, *AJ*, **146**, 30  
 Kuncarayakti, H., Anderson, J. P., Galbany, L., et al. 2018, *A&A*, **613**, A35  
 Logroño-García, R., Vilella-Rojo, G., López-Sanjuan, C., et al. 2017, in *Highlights on Spanish Astrophysics IX*, eds. S. Arribas, A. Alonso-Herrero, F. Figueras, et al., 283  
 Logroño-García, R., Vilella-Rojo, G., López-Sanjuan, C., et al. 2019, *A&A*, **622**, A180  
 Lyman, J. D., Galbany, L., Sánchez, S. F., et al. 2021, *MNRAS*, **495**, 992  
 Maoz, D., Mannucci, F., & Nelemans, G. 2014, *ARA&A*, **52**, 107  
 Marín-Franch, A., Taylor, K., Cenarro, J., Cristóbal-Hornillos, D., & Moles, M. 2015, in *IAU General Assembly*, 2257381  
 Martínez-Solauche, G., González Delgado, R. M., García-Benito, R., et al. 2021, *A&A*, **647**, A158  
 Maund, J. R., Fraser, M., Smartt, S. J., et al. 2013, *MNRAS*, **431**, L102  
 McCully, C., Jha, S. W., Foley, R. J., et al. 2014, *Nature*, **512**, 54  
 Molino, A., Benítez, N., Moles, M., et al. 2014, *MNRAS*, **441**, 2891  
 Nomoto, K., Suzuki, T., Shigeyama, T., et al. 1993, *Nature*, **364**, 507  
 Padilla, C., Castander, F. J., Alarcón, A., et al. 2019, *AJ*, **157**, 246  
 Pascual, S., Gallego, J., & Zamorano, J. 2007, *PASP*, **119**, 30  
 Perlmutter, S., Gabi, S., Goldhaber, G., et al. 1997, *ApJ*, **483**, 565  
 Perlmutter, S., Aldering, G., Goldhaber, G., et al. 1999, *ApJ*, **517**, 565  
 Phillips, M. M., Lira, P., Suntzeff, N. B., et al. 1999, *AJ*, **118**, 1766  
 Prentice, S. J., & Mazzali, P. A. 2017, *MNRAS*, **469**, 2672  
 Ransome, C. L., Haberman, Mawson, S. M., Darnley, M. J., James, P. A., & Percival, S. M. 2022, *MNRAS*, **513**, 3564  
 Riess, A. G., Press, W. H., & Kirshner, R. P. 1996, *ApJ*, **473**, 88  
 Schreiber, C., Glazebrook, K., Nanayakkara, T., et al. 2018, *A&A*, **618**, A85  
 Shivvers, I., Modjaz, M., Zheng, W., et al. 2017, *PASP*, **129**, 054201  
 Smartt, S. J. 2015, *PASP*, **32**, e016  
 Smith, N., Li, W., Filippenko, A. V., & Chornock, R. 2011, *MNRAS*, **412**, 1522  
 Taddia, F., Stritzinger, M. D., Bersten, M., et al. 2018, *A&A*, **609**, A136  
 Thomas, D., Steele, O., Maraston, C., et al. 2013, in *The Intriguing Life of Massive Galaxies*, eds. D. Thomas, A. Pasquali, & I. Ferreras, *IAU Symp.*, **295**, 129  
 Van Dyk, S. D. 2017, *Philos. Trans. R. Soc. London Ser. A*, **375**, 20160277  
 Vilella-Rojo, G., Viironen, K., López-Sanjuan, C., et al. 2015, *A&A*, **580**, A47

## Appendix A: Sample list and NCRs

**Table A.1.** Supernova sample. Name of the SN, type, right ascension, declination, and redshift.

NAME	TYPE	RA	DEC	z
SN2014ge	Ib	181.2146	26.9963	0.0019
SN2017igf	Ia	175.7069	77.3702	0.0060
SN2019clq	Ib	114.1563	74.4465	0.0120
SN2019tdf	Ia	266.6207	30.7069	0.0155
SN2019upq	II	217.2838	27.4502	0.0143
SN2019xdf	II	141.9661	68.4119	0.0130
SN2020aapr	II	112.0179	66.3748	0.0150
SN2020acac	II	217.7526	25.4889	0.0151
...				

**Note:** The full table is available at the CDS.

**Table A.2.** NCR values for the broad-band filters.

Name(Type)	uJAVA	gSDSS	rSDSS	iSDSS	zSDSS
ASASSN-14kg(II)	0.484	0.775	0.745	0.749	0.780
ASASSN-15lf(II)	0.271	0.130	0.096	0.089	0.050
ASASSN-15tw(II)	0.571	0.324	0.305	0.269	0.232
ASASSN-16eu(II)	0.441	0.269	0.200	0.203	0.207
ASASSN14az(IIb)	0.603	0.185	0.122	0.087	0.077
ASASSN15bd(IIb)	0.819	0.661	0.576	0.578	0.569
AT2020drl(II)	0.062	0.000	0.034	0.016	0.160
...					

**Note:** The full table is available at the CDS.

**Table A.3.** NCR values for the narrow-band filters.

Name(Type)	J0378	J0395	J0410	J0430	J0515	J0660	J0861
ASASSN-14kg(II)	0.691	0.666	0.589	0.689	0.749	0.740	0.738
ASASSN-15lf(II)	0.120	0.262	0.044	0.122	0.110	0.097	0.081
ASASSN-15tw(II)	0.000	0.565	0.428	0.534	0.355	0.289	0.200
ASASSN-16eu(II)	0.664	0.405	0.280	0.326	0.346	0.289	0.271
ASASSN14az(IIb)	0.000	0.184	0.316	0.157	0.131	0.162	0.120
ASASSN15bd(IIb)	0.821	0.773	0.732	0.736	0.707	0.623	0.551
AT2020drl(II)	0.154	0.421	0.000	0.000	0.024	0.082	0.167
...							

**Note:** The full table is available at the CDS.

**Table A.4.** NCR values for the narrow-band filters after subtracting the corresponding underlying continuum.

Name(Type)	H $\alpha$ + [NII]	[OII]	Ca II triplet	H $\alpha$	[NII]
ASASSN-14kg(II)	0.952	0.000	0.000	0.953	0.892
ASASSN-15lf(II)	0.000	0.000	0.000	0.000	0.000
ASASSN-15tw(II)	0.171	0.000	0.000	0.179	0.030
ASASSN-16eu(II)	0.672	0.000	0.000	0.681	0.272
ASASSN14az(IIb)	0.035	0.000	0.000	0.040	0.000
ASASSN15bd(IIb)	0.753	0.584	0.000	0.759	0.659
AT2020drl(II)	0.000	0.000	0.000	0.000	0.000
...					

**Note:** The full table is available at the CDS.

**Table A.5.** Results of the SED fitting with all 12 filters.

Name	Type	$\log < t_s (yr) >$	$A_V$ (mag)	$\log$ SFR	$\log$ M
SN2021yok	II	$6.60^{0.08}_{-0.13}$	$2.10^{0.03}_{-0.15}$	$0.55^{0.19}_{-0.23}$	$6.85^{0.07}_{-0.12}$
SN2022prv	II	$6.80^{1.31}_{-0.17}$	$0.80^{0.41}_{-0.80}$	$-0.44^{0.18}_{-1.29}$	$6.06^{0.05}_{-0.16}$
SN1998ar	II	$6.60^{0.10}_{-0.11}$	$1.40^{0.03}_{-0.17}$	$0.12^{0.17}_{-0.27}$	$6.42^{0.07}_{-0.19}$
SN1990L	Ia	$7.65^{0.05}_{-0.10}$	$0.00^{0.00}_{-0.00}$	$-3.31^{0.07}_{-0.03}$	$3.99^{0.07}_{-0.04}$
SN2019ltw	IIb	$6.55^{0.01}_{-0.30}$	$1.40^{0.05}_{-0.09}$	$0.93^{0.52}_{-0.01}$	$7.18^{0.24}_{-0.02}$
SN2018cow	Ic	$9.50^{0.02}_{-2.47}$	$0.00^{1.43}_{-0.00}$	$-2.16^{1.79}_{-0.02}$	$6.87^{0.02}_{-0.28}$
SN2020acac	II	$8.40^{0.02}_{-1.78}$	$0.00^{1.50}_{-0.00}$	$-1.35^{1.74}_{-0.03}$	$6.72^{0.05}_{-0.05}$
...					

**Note:** The full table is available at the CDS.