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Limits on Radioactive-Powered Emission Associated With a Short-Hard GRB 070724A in a Star-Forming Galaxy

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ABSTRACT
We present results of an extensive observing campaign of the short duration, hard spectrum gamma-ray burst (GRB) 070724A, aimed at detecting the radioactively-powered emission that might follow from a binary merger or collapse involving compact objects. Our multi-band observations span the range in time over which this so-called Li-Paczynski mini-supernova could be active, beginning within 3 hours of the GRB trigger, and represent some of the deepest and most comprehensive searches for such emission. We find no evidence for such activity and place limits on the abundances and the lifetimes of the possible radioactive nuclides that could form in the rapid decompression of nuclear-density matter. Furthermore, our limits are significantly fainter than the peak magnitude of any previously detected broad-lined Type Ic supernova (SN) associated with other GRBs, effectively ruling out a long GRB-like SN for with this event. Given the unambiguous redshift of the host galaxy ($z = 0.456$), GRB 070724A represents one of a small, but growing, number of short-hard GRBs for which firm physical/restframe quantities currently exist. The host of GRB 070724A is a moderately star-forming galaxy with an older stellar population component and a relatively high metallicity of $12+\log(O/H)_{K,D02} = 9.1$. We find no significant evidence for large amounts of extinction along the line of sight that could mask the presence of a SN explosion and estimate a small probability for chance alignment with the putative host. We discuss how our derived constraints fit into the evolving picture of short-hard GRBs, their potential progenitors, and the host environments in which they are thought to be produced.

Key words: Gamma-rays: Bursts: Individual: GRB 070724A, short GRBs, GRB host galaxies

1 INTRODUCTION
Long duration gamma-ray bursts (GRBs), lasting more than $\sim$ 2 seconds (Kouveliotou et al. 1993) are thought to originate from the collapse of massive stars (see Woosley & Bloom 2006). Short-duration, hard spectrum, gamma-ray bursts (SHBs), with a duration of less than $\sim$ 2 seconds have long been assumed to have a different astrophysical origin than long-duration events (Nakar 2007; Lee & Ramirez-Ruiz 2007), namely the coalescence of compact binaries, the most widely discussed being neutron star binaries (NS-NS) or a neutron star and a black hole binary (NS-BH) (Paczynski 1986; Eichler et al. 1989; Paczynski 1991; Narayan et al. 1991). The discovery of afterglows associated with SHBs (Gehrels et al. 2005; Hjorth et al. 2005b) led to the inference that they are associated with an older population of stars...
(e.g. Bloom et al. 2006). Subsequent follow-up observations supported the idea that these events would not be accompanied by supernovae (SNe) (Hjorth et al. 2005b; Bloom et al. 2006).

Nominally, without explosive nucleosynthesis of $^{56}$Ni to form a supernova, there should be no late-time optical emission after the afterglow has faded. However, during a NS-NS or NS-BH merger, dense material stripped from the star has been predicted to form large tidal tails (Rosswog 2007). Depending on the details of the encounter and the neutron star equation of state, a fraction of this can be dynamically ejected from the system. The subsequent decompres-

2. OBSERVATIONS & ANALYSIS

2.1 Swift BAT and XRT

The detection of GRB 070724A prompted an automated slew of the Swift spacecraft followed by XRT observations of the field beginning at $T + 72.1$ sec and which continued for roughly $\sim 10^6$ sec, at which point the source faded below the detector’s sensitivity threshold. We obtained the Swift BAT and XRT data on GRB 070724A from the Swift Archive\footnote{ftp://legacy.gsfc.nasa.gov/swift/data}. The data were processed with version 0.11.4 of the xrtpipeline reduction script from the HEAsoft 6.3.1\footnote{http://heasarc.gsfc.nasa.gov/docs/software/lheasoft/} software release, where we have employed the latest (2007-07-09) calibration files at time of writing. The reduction of XRT data from cleaned event lists output by xrtpipeline to science ready light curves and spectra is described in detail in Butler & Kocevski (2007a).

Our best estimate of the position of the XRT detection is $\alpha = 01^h51^m13.9^s, \delta = -18°35'39.1''$ with an error of $\sim 2$ arcsec. The fluence in the 15-150 keV BAT energy band is $3.0 \pm 0.7 \times 10^{-8}$ erg cm$^{-2}$. This is a extremely low value when compared to other Swift detected GRBs, being among the faintest 2% of Swift bursts and the faintest 20% of Swift detected short bursts. Given the redshift $z = 0.457$, reported by Cucchiara et al. (2007), the isotropic-equivalent energy released in gamma-rays is estimated at $E_{iso} = 1.55 \times 10^{49}$ ergs.

The XRT light curve (Figure 1) is consistent with an unbroken power law with a decay index $(F_\nu \propto t^{-\alpha})$ of $\alpha_X = 1.37 \pm 0.03$ over the entire span of the Swift observations. The winded timing mode spectrum from $t = 74$ sec to 4.1 ksec after the burst is well fit by a power law with photon index $(dN/dE \propto E^{-\Gamma})$ of $\Gamma = 1.5$ or correspondingly an energy index $(F_\nu \propto \nu^{-\beta_X})$, of $\beta_X = 0.5$ and photoelectric absorption yielding $N_H = 2.4 \pm 0.9 \times 10^{22}$ cm$^{-2}$, significantly greater than Galactic ($N_{H,\text{Galactic}} = 1.2 \times 10^{20}$ cm$^{-2}$) (Dickey & Lockman 1990).

We note that there are significant and rapid variations at early times in the X-ray hardness, which suggest that the excess column is a spurious result of the simple model assumption and that this emission is not due to the external shock afterglow (cf. Butler & Kocevski 2007a). Indeed, there are no hardness variations after $t \approx 300$ sec, and the column density for the photon counting (PC) mode spectrum at these times is consistent with the Galactic value. The hardness is consistent with the $\Gamma = 2.0$ commonly observed in other XRT detected GRBs, both long and short. (Butler & Kocevski 2007b).

Assuming $N_H = 1.2 \times 10^{20}$ cm$^{-2}$ and $\Gamma = 2.0$ ($\beta_X = -1.0$), the conversion from the 0.3 – 10.0 keV count rate to $\frac{dN}{dt}$.

The paper is organized as follows. In §2 we present a summary of our observations. The results of the mini-SN modeling are presented in §3, while the properties of the suggested host galaxy and surrounding galaxy population is discussed in §4. Finally, the implications of our results are presented in §5. Throughout the paper we assume $H_0 = 71$ km s$^{-1}$ Mpc$^{-1}$ and a $\Lambda$CDM cosmology with $\Omega_m = 0.27$, $\Omega_\Lambda = 0.73$. 
flux in \( \mu \text{Jy} \) at 1 keV is 0.044 \( \mu \text{Jy} \) cps\(^{-1}\). The unabsorbed flux after \( t \approx 300 \text{ sec} \) until \( t \approx 550 \text{ ksec} \) is \( F_E = (2.6\pm0.4) \times 10^{-5} \left( t/[10^3 \text{ s}] \right)^{-1.37\pm0.03} \mu \text{Jy} \). We find that the afterglow flux at \( t = 10 \text{ hrs} \) is \( F_{X,10} \approx 6.0 \times 10^{-14} \text{ ergs s}^{-1} \), which at a redshift of \( z = 0.457 \) gives a luminosity of \( L_{X,10} = 4.6 \times 10^{43} \text{ ergs s}^{-1} \). Assuming that this flux comes from an adiabatically expanding external shock in the slow-cooling regime, the expected (unabsorbed) \( R \)-band optical flux is a factor \( \approx 20 - 520 \) times higher, depending on the location of the cooling break below the \( X \)-ray band (e.g. Sari et al. 1998).

The resulting BAT and XRT detections are plotted in blue and red respectively in Figure 1. The expected optical flux from the forward shock falls within the grey region shown in the plot. Our optical observations discussed below rule out a bright optical afterglow for this event, although a cooling break located just below the XRT bandpass would allow for optical emission to have gone undetected and still fit within the framework of the standard external shock afterglow model. The lack of a bright optical afterglow for GRB 070724A is consistent with the trend for low-fluence bursts to have faint afterglows noted by Gehrels et al. (2008) and Nysewander et al. (2008). This makes GRB 070724A particularly interesting because it allows for the search of SN related optical emission without the need to contend with external shock powered afterglow emission.

### 2.2 Optical/NIR

Our earliest observations of the location of GRB 070724A were taken with the United Kingdom Infra-Red Telescope (UKIRT) at Mauna Kea, beginning with a 20 sec exposure in \( K \)-band at 13:11 UT on July 7th, roughly 2.5 hours after the trigger. Several epochs of observations followed in the \( J, H \), and \( K \) bands over the next 24 hours, ending approximately 1.1 days after the burst.

Our imaging campaign in the optical passbands began with the use of StanCam at the Nordic Optical Telescope (NOT) on La Palma/Spain starting at 04:18 UT, roughly 17.5 hours after trigger. These observations consisted of \( 3 \times 300 \text{ sec} \) in \( B \) and \( I \)-band and \( 3 \times 600 \text{ sec} \) in \( R \)-band. Comparison images of \( 6 \times 600 \text{ sec} \) in \( R \) were taken at 03:35 UT on August 16th 2007, 22.7 days after trigger.

We obtained an additional 600 sec \( J \)-band image with the ESO-VLT equipped with the FORS2 instrument starting at 09:08 UT (22.1 hours after trigger), followed by spectroscopic observations, which will be described below.

Our most extensive observations of the error circle of GRB 070724A were performed using the ANDICAM instrument operated by the SMARTS Consortium on the 1.3m
Figure 2. (Left) Stacked Keck R-band image (taken 18 days after the event) of the region around the putative host galaxy. The 3 sigma XRT error circle is plotting in red. (Right) Spatially resolved limiting magnitude contour map, derivied after removing an elliptical isophotal model for the host.

telescope at CTIO. Optical/IR imaging in J and I-bands began at 08:14 UT (21.3 hours after trigger) and continued for an additional 8 epochs with the last observations occurring at 07:31 UT on August 21st (27.8 days post trigger). Several dithered images were obtained in each filter, with total summed exposure times of 30 minutes in J and 36 minutes in I.

Finally, we obtained deep optical imaging with the 10m Keck I Telescope equipped with the LRIS instrument (Oke et al. 1995) at Mauna Kea, starting at 13:40 UT on August 11th. The co-addition of successive observations yields a total effective exposure time of 2720 seconds in R and 2120 seconds in g′-bands, followed by spectroscopic observations, which will be described below. All observations were conducted under photometric conditions and zeropointing was performed using a standard star field observed at an similar airmass.

2.3 Photometry

All of our observations were reduced using standard CCD packages in IRAF. Once reduced, the co-addition of successive observations was performed using the Swarp and Shiftadd software packages\(^3\) to produce weighted sum images of the host field. Astrometry was performed relative to USNO-B1 using at least 15 sources in common between our summed image and the catalog.

A custom pipeline was then used to perform photometry on individual and coadded frames, using aperture photometry via the Sextractor software package (Bertin & Arnouts 1996), to estimate instrumental magnitudes. An aperture size of 6/5× seeing (FWHM) was used for the analysis across all our science data. Our instrumental magnitudes were then compared to a standard star field typically taken on the night of each observations for zeropoint determination of each reduced frame. The aperture size used to photometrize the standard star frames was typically 2 – 3 times larger than the aperture used in the analysis of particularly deep, and hence crowded, images. For these images, an aperture correction was applied to account for the differing aperture radii.

To search for afterglow emission and to properly account for the host galaxy contribution, we employ a modified version of the public POIS-IPP package\(^4\).

Using the last image in each instrument series as a template of the host, we find no evidence for residual afterglow emission. Subtractions employed on images spanning much smaller time scales and taken with other facilities likewise show no signs of variability between observations. A subtraction between the coaddition of all R-band Keck observations taken 18 days after trigger and a 30 minute exposure taken 1 year after the event can be seen in Figure 3.

We estimate the limiting magnitudes on these observations by placing 1000 blank apertures at random positions on the images. We take the standard deviation of the resulting photometry distribution to represent the counts associated with the sky noise of that image and assume that the faintest observable object is 3 times this value for a 3σ limiting magnitude. The limiting magnitudes as a function of position on the host galaxy were calculated by taking this theoretical limiting magnitude due to the sky and adding in quadrature the error associated with the counts from the

\(^3\) http://terapix.iap.fr/rubrique.php?id_rubrique=49

\(^4\) http://pan-stars.ifa.hawaii.edu/project/IPP/software/
Figure 3. (Left) Stacked Keck R-band image (taken 18 days after the event) of the region around the putative host galaxy. The 3 sigma XRT error circle is plotting in red. (Right) The subtraction of the same field observed one year after the event showing no significant sources at the position of the XRT position. Similar subtractions were performed for all multi-epoch observations that spanned more than one day.

galaxy at a given pixel, which were assumed to be Poissonian and therefore \( \propto \sqrt{N} \). The resulting spatially resolved limiting magnitudes are shown as a contour plot in Fig. 2. We check these estimates by placing fake stars with a magnitude near our estimated limiting magnitudes in a subset of our frames and then test whether they would be detectable at varying thresholds above background.

The summary of the optical/NIR observations and their associated upper limits can be found in Table 1. Figure 1 shows the upper limits at different epochs along with the BAT and XRT detections. Our early R-band observations taken with the NOT at \( \sim 0.73 \) days effectively rule out any emission originating from the forward shock in the context of the canonical fireball model.

2.4 Spectroscopy

We obtained a single 900 sec spectrum (henceforth called “slit 1”) of the host galaxy (Bloom 2007) of GRB 070724A with LRIS (low-resolution imager and spectrograph) at the Keck I telescope, Mauna Kea, Hawaii, on August 11, 2007. The spectrum covers the wavelength range between 3500 and 9400 Å. A slitwidth of 1′′ and grism 600/4000 were used in the blue, providing a resolution of \( 4.0 \) Å FWHM. In the red, a grism 400/8500 was used providing a resolution of \( \sim 6.5 \) Å FWHM. The spectra were extracted and wavelength calibrated using standard tools in IRAF. All wavelengths given are in air. Flux calibration was done using the spectrophotometric standard star BD+174708. The longslit covered both the host galaxy and a neighboring galaxy to the south-west of the host galaxy at similar redshift (henceforth called G3). Note that the narrow slit of 1′′ does not contain the entire flux for both galaxies, which affects later analyses derived from the flux calibrated spectra.

Furthermore, we obtained spectra with LRIS at three other slit positions on October 10 and 11, 2007, with the same setting and calibrated with the standard star Feige 110. The three slits (“slit 3, 4 and 5”) cover a number of galaxies of which only three show strong emission or absorption lines allowing for redshift determinations.

We also obtained a 600 sec spectra with FORS 1 (called “slit 2”) at the VLT on Cerro Paranal, Chile, on July 25, 2007, starting 09:52 UT using grism 300V and a 1′′ slit, which has a resolution of 11 Å FWHM. Due to heavy fringing, the emission lines of the host are not detected with high significance. The slit was positioned in the north-south direction and therefore covered a slightly different part of the host galaxy. The other objects in the slit do not show any obvious emission lines. Reduction and calibrations were also performed with IRAF standard tasks and flux calibration performed using the spectrophotometric standard LTT9239. In contrast to the LRIS spectra, wavelengths for the FORS spectrum are in vacuum.

The position of the slits and the field around the host of GRB 070724A can be seen in Figure 4.

3 THE HOST GALAXY AND ITS NEIGHBORS

The XRT position of GRB 070724A, with an error radius of 1.6 arcsec, is offset by roughly 0.8 arcsec from the center of a \( z = 0.456, \) \( 19.55 \) R-band mag, galaxy that we identify as the host galaxy. Following the formalism for small offsets (Appendix B in Bloom et al. 2002), we estimate that the chance alignment of a burst given this redshift, offset, and host magnitude is exceedingly small, roughly \( P = 0.002. \) At this redshift, the projected offset would transform to a physical distance of 5 kpc from the center of the galaxy although this value is largely unconstrained due to the relative size of the projected offset to the XRT error radius. Offsets of order
of a few kpc have been seen previously in some short bursts, though much larger offsets have frequently been reported (Bloom et al. 2007; Troja et al. 2008). As short GRBs are assumed to originate from the merger of two compact objects, the kick from the SNe explosions could lead to large offsets from their birthplace and their host galaxies, which seems to agree with the observations (e.g. Fryer, Woosley, Hartmann 1999; Bloom, Sigurdsson, & Pols 1999; Belczynski et al. 2006). No other bright sources have been found within or near the XRT position, however, two galaxies at a similar redshift have been discovered at a distance of several arcsec from the GRB position. We analyze the properties of these galaxies together with the host galaxy in this section (see also Sec. 3.2).

Figure 5 shows the spectra of the host as well as all galaxies where the redshift could be determined, named G2, G3, G4 and G5 in Figure 4.

3.1 Properties of the Host

The star-forming nature of the host galaxy is suggested by the detection of several emission lines in the spectra, including [O ii] λλ3727,3729, [O iii] λλ4959, 5007, Hβ and Hγ as well as the Ca H&K absorption lines. Unfortunately, the redshift of z=0.456 places Hα at ~9600Å outside the wavelength range of our spectra. In the Table 2, we give the emission line values for the two different slit positions, which probe slightly different regions, where “slit 1” indicates the Keck observation and “slit 2” indicates the position of the FORS observation. (in North-South direction). All other slits shown in Figure 4 were observed with Keck.

Using the detected emission lines in slit 1 and slit 2, we can derive a number of properties, including the extinction, star-formation rate and metallicity of the star-forming regions. As Hα is not available for extinction measurements using the Balmer line decrement (Osterbrock 1989), we use the ratio between Hγ and Hβ, which is 0.47 (for T_e = 10^4 K, n_e = 100 cm^{-3}) in the absence of any extinction. Unfortunately, Hγ is not observed in slit 2, likely due to the lower S/N, so we rely solely on the slit 1 for our extinction estimates. Using the measured line fluxes for slit 1 shown Table 2 and assuming R_V = 3.1 and the Cardelli et al. (1989) extinction curve, we obtain a reddening value of E(B−V)=1.2 ± 0.2 mag. The Galactic reddening along the line of sight towards the host is only E(B−V)=0.013 mag (Schlegel et al. 1998). From the detection of the Ca-break and the shape of the spectrum, we conclude that the galaxy has an underlying stellar older population which contributes some absorption in the Balmer lines. This affects Hγ more than Hβ, which, if we were able to correct for it, would lead to a lower value for the extinction. Due to the low resolution of the spectrum, we are not able to determine the strength of the stellar absorption.

The star-formation rate can be obtained from the Hα or the [O ii] λλ3727,3729 line flux using the conversion from Kennicutt (1998). Hα is outside of our spectral range and we therefore use [O ii], which has a more indirect connection to the ongoing SFR than Hα. Taking the unextincted values, we then obtain an absolute SFR of 0.83±0.03 M_☉/yr from the Keck slit (slit 1) and an specific SFR scaled with the g-band magnitude of 21.75 mag (at z = 0.456 equal to restframe u-band) of 1.64 M_☉yr^{-1}(L/L*)^{-1} assuming that an L* galaxy has an absolute magnitude of M_B = −21 mag (Christensen et al. 2004). The extinction corrected fluxes give a rather high absolute SFR or 129±4 M_☉/yr (the errors from the extinction correction are not propagated) and a SSFR of 253 M_☉yr^{-1}(L/L*)^{-1}. Due to the stellar absorption in the Balmer lines mentioned above, the true SFR lies inbetween the unextincted and the extinction corrected values.

Using the R_{23} parameter, which takes the ratio of the [O ii], [O iii] and Hβ emission line fluxes (uncorrected for possible extinction in the host), we derive a metallicity for the host galaxy. The most frequently applied calibration for GRB hosts is the one taken from Kewley & Dopita (2002), which gives a rather high metallicity of 12+log(O/H)_{KDO2} = 9.1 for this galaxy as compared to the solar metallicity value of 8.66 (Asplund et al. 2004). Using the extinction corrected fluxes leads to the same value for the metallicity. The intrinsic error of this method for determing the metallicity is around 0.2 dex.

The Hβ-EW is dependent on the age of the stellar population and can be used as an upper limit for the age according to stellar evolutionary models, e.g. Leitherer et al.

Table 2. Photometric observations of the host galaxy of GRB 070724A

<table>
<thead>
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<th>Instrument</th>
<th>Filter</th>
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<th>Flux (μJy)</th>
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<td>21.12</td>
<td>9.120</td>
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<tr>
<td>Keck I/LRIS</td>
<td>g'</td>
<td>21.64</td>
<td>8.02</td>
</tr>
<tr>
<td>Keck I/LRIS</td>
<td>R</td>
<td>19.56</td>
<td>46.56</td>
</tr>
<tr>
<td>CT1.3m/ANDICAM</td>
<td>I</td>
<td>20.01</td>
<td>24.32</td>
</tr>
<tr>
<td>UKIRT/UFTI</td>
<td>J</td>
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</tr>
<tr>
<td>UKIRT/UFTI</td>
<td>H</td>
<td>20.05</td>
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</tr>
<tr>
<td>UKIRT/UFTI</td>
<td>K</td>
<td>19.56</td>
<td>10.000</td>
</tr>
</tbody>
</table>

Figure 4. Field around the host galaxy of GRB 070724A and positions of the slits. The galaxies in the field with measured redshifts are indicated according to the notation in Table ??.
A value of log(EW(Hβ))=1.04 or 1.22 for the two slit positions corresponds to an age of about 10 Myrs for an instantaneous starburst and about 100 Myr for continuous star formation depending slightly on the metallicity. The Balmer break, which is clearly visible in the spectrum of the host galaxy, furthermore indicates an age of at least 100 Myrs, indicating an older population of stars. This is consistent with the photometry for this host and quite interesting given the detection of strong emission features. It could indicate recent starburst activity, possibly due in part to a merger event with one of its nearby neighbors, revitalizing what was otherwise a old, red galaxy.

3.2 Is the Host Part of a Cluster?

Since the first well-localized SHB (GRB 050509b) was found to be in a cluster (Bloom et al. 2006), the frequency of SHBs in clusters has been a subject of study (Berger et al. 2006). To investigate the potential for diffuse X-ray emission from a cluster of galaxies, we search for extended sources using wavdetect (Freeman et al. 2002). We analyzed the X-ray data on scales of 0.8, 1.1, 1.6, 2.2, 3.1 arcmins and found no strong detections which are not centered on point sources in the 91.7 ksec exposure XRT PC mode image (Detection centered on point sources are considered suspect due to the broad wings of the XRT point spread function). Assuming a thermal Bremsstrahlung spectrum with $kT = 5.0$ keV (e.g. Bloom et al. 2006) at $z = 0.457$, any cluster must be fainter than $2 \times 10^{33}$ erg s$^{-1}$ (3σ). This is an order of magnitude fainter than the rich X-ray cluster associated with GRB 050509B.

We determined the redshift of the galaxies around the host in order to determine a possible group or cluster as found for two other SHB, (Berger et al. 2007). For most of the galaxies in the slit, we could not determine any redshift due to a lack of strong emission or absorption lines. We did not attempt to determine redshifts via SED modelling of the spectral shape for these sources, as such a method would not have yielded the necessary redshift accuracy to determine whether the galaxies could belong to a cluster or a group connected to the host galaxy.

We find that the galaxy about 5 arcsec North-West of the host, called G3 in Fig. 4 and Table 7?, has the same redshift as the host. Also the galaxies G2 and G5, which are about 10–20 arcsec from the host both have a similar redshift of $z=0.473$ and 0.471. However, this would correspond to a distance of 57–115 kpc from the host to G2+G5, which makes it unlikely that the host+G3 are gravitationally bound to G2+G5. G4 only shows one emission line which, if associated with O II, would give a redshift of $z=0.860$. Together with the X-ray observations, we therefore conclude that the host may be a member of a smaller group, but does not belong to a larger cluster.

3.3 Properties of Surrounding Galaxies

A small number of the galaxies probed by the different slits show prominent emission or absorption lines for the remaining galaxies. The only galaxy which was identified to have the same redshift as the host galaxy, G3, has a young stellar population. Its spectral slope is rather flat and does not show any 4000 Å Balmer break. There is also no detection of the Ca II absorption lines. The Hβ equivalent width suggests an age between 4 and 20 Myrs and only one of the Ca II lines could be clearly identified. G2, which has a slightly higher redshift than the host, shares some of its properties, with a moderate Hβ EW indicating an age between 10 and 100 Myrs and a clear 4000 Å break. However, only one of the Ca II absorption lines could be clearly identified. G4 has only one emission line likely identified with the O II doublet, which indicates a younger stellar population. G5 has no emission lines, but clear Ca II absorption and a Balmer break, which suggests an older galaxy.

We utilize the same methods to determine the SFR, extinction and metallicity as those described in the previous section to determine the properties of the other galaxies probed by the different slits. For G2, we get an extinction from the Balmer line ratio between Hβ and Hγ of E(B−V)G2=0.96±0.7 mag. From the flux of [O II] λ 3729, in 12+log(O/H)G2 =9.05 (Kewley & Dopita 2002) or 8.9 using the extinction corrected fluxes. For G3, we find that the extinction is consistent with zero, which agrees with the indications of a young, star-forming galaxy, although dusty star-burst galaxies have been observed as well. The G3 SFR from [O II] is 3.4±0.1 M⊙ yr$^{-1}$ with a SSFR of approximately 21.2 M⊙ yr$^{-1}$ (L/L*)$^{-1}$ and the metallicity of either 12+log(O/H)G3 =8.23 for the lower branch solution or 12+log(O/H)G3 =8.61 for the upper branch, which in both cases would imply subsolar metallicities. For a clear determination of the metallicity, the detection of Hα or the temperature sensitive line [O II] λ 4363 would be necessary (Izotov et al. 2006). In case the emission line identified in G4 is [O II] λ 3727, 3729, we get a SFRG4 of 4.7±0.1 M⊙ yr$^{-1}$ or a SSFR of 47.0 M⊙ yr$^{-1}$ (L/L*)$^{-1}$. 

Figure 5. Spectra of the host galaxy and some of its neighbours with determined redshifts. The lines indicated correspond to the lines noted in Table 7. 

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4 CONSTRANTS ON EXCESS OPTICAL EMISSION

4.1 Constraints on an Associated Core-Collapse Supernova

Through the use of our non-detection of the optical afterglow, we can derive limits on the optical contribution from SN of Type Ib or Ic associated with this event. In Figure 6, we plot three different SN Ic templates together with the upper limits in the $R$-band. In order to convert our photometry to rest-frame values (k-correction), we fit the SED (in flux units) derived from data points and interpolations of the $UBVRI$ lightcurve for each day, redshifting the SED to $z = 0.456$, and determined the new fluxes at the effective $R$-band wavelength before transforming these flux values back into magnitudes. Furthermore, the templates were corrected for time dilation as well as for Galactic extinction. The comparison templates used are the “standard” long-duration GRB-SN 1998bw (Galama et al. 1998), an SN associated with an X-ray flash (XRF), SN 2006aj (Sollerman et al. 2006), and faint Ic supernovae 2002ap (Foley et al. 2003) and 1994i (Richmond et al. 1996). All observed SNe that have been associated with a long-duration GRB are within 0.5 mag of the peak luminosity of SN 1998bw, although some XRF-SNe have appeared to be significantly
f = 10

that goes into radioactive decay per \( \beta c \) function of three parameters: the mass, \( M \), and velocity, \( \beta c \), of the ejecta, and the fraction of the ejecta energy, \( f \), that goes into radioactive decay per \( c \)-fold in time over the relevant timescales. Following Li & Paczyński (1998), the peak luminosity, time to reach this value, and effective temperature of the ejecta, can be given respectively as:

\[
L_{pk} = 2.1 \times 10^{44} \text{ ergs s}^{-1} \times \left( \frac{f}{0.001} \right) \left( \frac{M}{0.01M_{\odot}} \right)^{1/2} \left( \frac{3\beta}{\kappa_e} \right)^{1/2} \left( \frac{\kappa}{\kappa_e} \right)^{-1/2},
\]

\[
t_{pk} = 0.98 \text{ days} \left( \frac{M}{0.01M_{\odot}} \right)^{1/2} \left( \frac{3\beta}{\kappa_e} \right)^{-1/2} \left( \frac{\kappa}{\kappa_e} \right)^{1/2},
\]

and

\[
T_{\text{eff}, pk} = 2.5 \times 10^5 \text{ K} \times \left( \frac{f}{0.001} \right)^{1/4} \left( \frac{M}{0.01M_{\odot}} \right)^{-1/8} \left( \frac{3\beta}{\kappa_e} \right)^{-1/8} \left( \frac{\kappa}{\kappa_e} \right)^{-3/8}.
\]

Here \( \kappa \) is the average opacity (\( \kappa_e \) is the opacity caused by electron scattering). The time of peak luminosity occurs when the photons can diffuse outward on the dynamical time, so that most of the thermal energy that is produced by the radioactive decay can be radiated efficiently before suffering significant adiabatic cooling. Since the photon diffusion time is \( \sim \tau \) times larger than the source light crossing time, \( R/c \), and \( \tau = \beta ct \), we can equate the dynamical time \( t_{\text{dyn}} \sim t = R/c \) with \( R/c \) at \( t_{pk} \) and therefore \( \beta^{-1} \sim \tau(t_{pk}) \propto \kappa M/(\beta \sigma T_{\text{eff}, pk})^2 \) and therefore \( t_{pk} \propto (\kappa M/\beta)^{1/2} \). The peak luminosity is approximately given by \( L_{pk} \sim f M c^2 / t_{pk} \propto f(\beta M/\kappa)^{1/2} \), and the effective temperature is obtained by equating this luminosity to that of a black body, \( L_{pk} \sim 4\pi R^2 (t_{pk}) \sigma T_{\text{eff}, pk}^4 = 4\pi (\beta c t_{pk})^2 \sigma T_{\text{eff}, pk}^4 \) and therefore \( T_{\text{eff}, pk} \propto T_{pk}^{1/2} (\beta t_{pk})^{-1/2} \propto f^{1/4} \kappa^{-3/8} (\beta M)^{-1/8} \).

As a result, the luminosity of the resulting light curve is directly proportional to the fraction of the ejecta energy that goes into radioactive decay (i.e. the ejecta efficiency). An increase in the assumed mass in the ejecta has the effect of both increasing the total luminosity of the emission as well as increasing the time, \( t_{pk} \), to peak luminosity, \( T_{pk} \). On the other hand, an increase of ejecta velocity leads to an
increase in the peak luminosity but a decrease in the peak duration.

The most efficient conversion of nuclear energy to observable luminosity is provided by elements with a decay timescale comparable to the timescale it takes the ejected debris to become optically thin, $t_{\tau}$. In reality, there is likely to be a large number of nuclides with a very broad range of decay timescales. The ejecta is optically thick at early stages but its optical depth falls rapidly as the ejecta expands. As this happens, radiative losses at the photosphere become important and the peak emission occurs when the optical depth reaches $\sim 1/\beta$. Current observational limits thus place interesting constraints on the abundances and the lifetimes of the radioactive nuclides that form in the rapid decompression of nuclear-density matter — they should be either very short or very long when compared to $t_{\tau}$ so that radioactivity is inefficient in generating a high luminosity.

Figure 7 shows the calculated $R$-band light curves for a range of ejecta mass $M$, velocity $\beta$ and the fraction of the ejecta energy $f$ that goes into radioactive decay, effectively the energy conversion factor. Plotted as black filled arrows are three $R$-band upper limits taken at roughly 0.74, 18.14, and 22.70 days from left to right. For comparison, we have also plotted the $R$-band upper limits for GRB 050509b presented by Hjorth et al. 2005b and Bloom et al. 2006. Unfortunately, many of our observations between days 2 and 18 are less constraining, so that we can only place limits on a very fast, massive, and highly efficient ejecta. Likewise, the deep upper limit at late time can only constrain a very slow, high mass ejecta.

Also plotted as black hollow arrows are our equivalent limits if the optical emission from the LP-SN transient suffered the maximum $\sim 3$ magnitudes of rest frame $R$-band extinction that is implied by our measured reddening value of $E(B-V) = 1.2 \pm 0.5$ mag. The LP-SN could have been significantly brighter and have been missed by our search in this worst case scenario, allowing us to constrain only an extremely massive and highly efficient ejecta.

In Figures 8 & 9, we show the parameter space excluded by our observations for GRB 070724A as well as the upper limits on emission from GRB 050509b placed by Hjorth et al. (2005a). The areas in parameter space which are shown are divided into regions that are dominated by the constraint from a small number of upper limits. The dividing lines between these regions are demarcated by sharp changes in the direction of the contour lines.

In Figure 8, the right to lower right side of the plot is dominated by the upper limit of 23.72 mag in $R$-band at 0.5068 days, while the other half of the plot is dominated by the upper limit of 27.4 mag in $R$-band at 12.44 days, both epochs listed in the host frame. The darker blue regions of the plot represent deeper limits on the fraction of the ejecta energy $f$ that goes into radioactive decay. The parameter space that is least constrained represents the intermediate velocity ejecta at a wide range of masses. This corresponds to the time range for peak between days 2 and 18 for which we have only shallow upper limits in magnitude.

For comparison, the upper limits made by Hjorth et al. (2005a) place stricter limits on the presence of a LP-SN for GRB 050509b. The derived limits for GRB 050509b shown in Figure 9 are dominated by the upper limit of 26.6 mag in the $R$-band at 1.85 days (lower right) and by the upper limit of 27.5 mag in the $V$-band at 3.92 days (upper left). At the very upper left corner (high $M$ and low $\beta$) there is a small region which is dominated by the upper limit of 26.7 mag in $R$-band at 22.83 days. This is natural since $t_{pk} \propto (M/\beta)^{1/2}$ (see eq. 2).

Finally, there is some ambiguity as to the time of origin in the LP-SN model, as the onset of the GRB and the onset of the LP-SN transient may not necessarily coincide. Although there is currently no theoretical prediction for this
time-lag, we can constrain delays of up to 10 days even for the moderately massive efficient ejecta under consideration in Figure 7. Any such delay would have resulted in the detection of a bright transient during our deepest observation at 18.14 days post trigger. For an extremely inefficient ejecta (dash-dotted line), we can only constrain the delay in the onset of emission from 7-10 days post trigger. Alternatively, the late time R-band limit on the LP-SN emission from GRB 050509b placed by Hjorth et al. (2005a) rule out a delays up to 11 days for all models under consideration in Figure 7.

5 DISCUSSION

We present the results of a deep and extensive observing campaign of the short duration hard spectrum (SGRB) GRB 070724A. Although our observations do not reveal an optical or NIR transient associated with this event, the derived upper limits can constrain optical emission from a traditional forward shock as well as the physical parameters of SNe, or SNe-like models. Our early NOT and late time Keck observations show no sign of the optical emission from a forward shock in the slow cooling regime, as predicted by the fireball model. However, this optical contribution could have easily eluded detection given the weakness of the event combined with the potential extinction by the host galaxy.

Given the measured reddening value due to the host of E(B−V)=1.2±0.5 mag, any optical emission could have suffered as much as A_V ~ 4 magnitudes of extinction if the burst occurred well within the galaxy’s core. In this case, the optical limits on a mini-SN quickly become unconstrained and our late time observations would likewise barely constrain even most broad-lined Type Ic SNe. If the GRB exploded at a large offset to the host galaxy, as has been seen in other short GRB events, or if it the event occurred on the near side rather than the far side of the host, than much of the measured extinction would be largely irrelevant. The measured N_H could potentially help to resolve these different scenarios. Of the two Swift events which have solidly been shown to have A_V > 3, both have had large N_H column density at late times, of the order 5 x 10^{22} (Perley et al. 2009). We measure a similar column densities at t < 300 sec in association with large hardness variations in the X-ray detections, suggesting that the excess column density does not reflect the intrinsic absorption of the host. At late times (t > 300 sec), the N_H value becomes consistent with the Galactic value N_H,Galactic = 1.2 x 10^{20} cm^{-2} (Dickey & Lockman 1990). Furthermore, the extinction value is derived from emission line measurement, indicating the extinction towards the regions of the galaxy that harbor younger populations of stars. These regions may differ greatly in their dust content from the extended disk or halo in which one would expect to find the older population of stars which are thought to be the progenitors of SHBs. Therefore we conclude that the site of GRB 070724a was likely not affected by large amounts of extinction, although this remains a caveat in the proceeding discussion.

Because our limiting magnitudes are deepest at early and late times, our observations primarily constrain an efficient, very slow or a very fast ejecta with mass higher than 0.01 M_solar. Our observations do little to exclude a low mass ejecta at intermediate velocities resulting in emission reaching peak luminosity near 1 day. This can be seen in the case of GRB 050509b, where two limiting magnitudes in R-band are of similar depth although the earlier observation is far more constraining. This is due to the characteristic timescales associated with the peak emission predicted by the LP-SN model. This underscores the importance of deep optical observations hours to days after the short-hard GRB.

To be sure, our calculations are based on a very simplified analytic model, so it is not clear whether this basic result would hold for more realistic models, e.g. Kulkarni (2005).

We note that GRB 050509b and GRB 070724A represent the fainter end of the energetics distribution (Kann et al. 2008), at 2.40 x 10^{58} and 1.54 x 10^{59} ergs respectively, for short bursts thus far detected by Swift. If there is a correlation between the energy released in the prompt gamma-ray emission and the velocity and/or mass of the resulting ejecta, not an unreasonable scenario, then one would expect more energetic bursts to result in brighter emission from a mini-SN. Therefore, deep optical observations of more energetic SHB events may yet yield detectable emission.

To date, only 14 short bursts (GRB 050709, GRB 050724, GRB 051221A, GRB 060121, GRB 060313, GRB 061006, GRB 061201, GRB 070707, GRB 070714B, GRB 070809, GRB 071227, GRB 080503, GRB 080905A, GRB 090510) have resulted in the detection of an optical afterglow and therefore a position determined at subarcsecond level. Only four short GRBs detected thus far were proposed to be associated with early type galaxies such as elliptical or early type spiral galaxies, namely GRB 050509B (Gehrels et al. 2005; Bloom et al. 2006), GRB 050724 (Barthelmy et al. 2005; Berger et al. 2005; Gorosabel et al. 2006; Malevansky et al. 2007), and possibly GRB 050813 (Prochaska et al. 2006; Ferrero et al. 2007; Berger 2006) and GRB 060502B (Bloom et al. 2006). The lack of emission lines in these galaxies implies low limits on any ongoing star-formation, the age of their stellar populations is of the order of Gyrs, and their metallicity is around solar. All other short GRB hosts are late type or irregular, star forming galaxies, although with lower star formation rates as compared to their long GRB host counterparts. The best studied examples of late type galaxies associated with short bursts to date are the irregular host of GRB 050709 (Hjorth et al. 2005a; Covino et al. 2006; Fox et al. 2005) and the host of GRB 051221A (Soderberg et al. 2006).

The host galaxy of GRB 070724a is one of the growing number of star forming galaxies to harbor a short GRB. Despite the lack of an optical detection, the position of the XRT error circle covering most of the galaxy makes an association very likely. The host galaxies of short bursts have been found to be a rather diverse population, in contrast to the more uniform sample of star forming galaxies or at least star-forming regions within larger galaxies for long-duration GRB hosts (Christensen et al. 2004; Fruchter et al. 2006). Recently, Berger (2009) published properties for the full set of short GRB host galaxies, whose properties could be determined. They find that the range of SFRs is moderate with 1–10 M_{solar} yr^{-1}, with some outliers with very low SFR, which is slightly lower than found for long duration GRBs (e.g. Christensen et al. 2004). The metallicity is in general 0.6 dex (Modjaz et al. 2008) higher than for long-duration GRB hosts, however, this is based on a rather small sam-
Figure 10. Left: The $R_c$ magnitude of the afterglow (or upper limit thereto) at 1 day after the GRB in a common $z = 1$ frame adopted from Kann et al. (2008), vs. the specific star formation rate of the host galaxy as presented by Rhoads (2008). GRB 070724A represents the deepest limit on the existence of short time scale supernova-like emission for a short GRB in a moderately star forming galaxy. The color scheme represents the isotropic equivalent energy of the event. Right: The ratio of the extended emission fluence to that of the initial spike (or limit thereof) vs. the host’s specific star formation rate of a subset of short GRBs with known redshift. There is no clear connection between the events with significant extended emission and the properties of the host galaxy.

This diversity in both the host galaxy population and also the prompt energetics of short GRBs hints at the possibility of progenitors that may not be entirely restricted to the most widely favored scenario involving the merger of compact binaries. The dichotomy between the specific star formation rates of short GRBs hosts, in particular, has been proposed by Rhoads (2008) as a possible indication of multiple progenitors. Unlike the optical limits previously obtained for GRB 050509b (Hjorth et al. 2005a), which occurred within a quiescent galaxy with little or no ongoing star formation, our observations of GRB 070724A, along with those of GRB 050709, place the deepest constraints on the presence of short timescale supernova-like emission of a short GRB in a moderately star forming host galaxy.

This can be seen in the left panel of Figures 10, where we plot the $R_c$ magnitude of the afterglow (or corresponding upper limit) 1 day after the GRB in a common $z = 1$ frame, adopted from Kann et al. (2008), vs. the specific star formation rate of the host galaxy for a sample of short GRBs with known redshift estimated by Rhoads (2008). The color scheme in Figure 10 represents the isotropic equivalent energy $E_{iso}$ of the event. This effectively sets a limit to short timescale emission from a mini-SN, as any emission in excess to the afterglow light curve would have been detected. The constraints on short timescale supernova like emission are deepest for GRB 050509B and GRB 070724A for quiescent and moderately star forming galaxies respectively.

The extended emission observed to follow the initial gamma-ray spike in many events detected by Swift and BATSE has also been suggested as a possible indicator of progenitor diversity within the short GRB population. This phenomenon can generally be characterized by additional gamma-ray emission lasting several tens of seconds and fluence values that in some cases exceed that of the initial short GRB spike. The nature of this emission has recently been discussed in detail by Perley et al. (2008) who found that roughly 30% of short GRBs detected by Swift exhibit it in some form. They also find that the extended-to-prompt fluence ratio exhibits a large variance, with no clear correlation between the brightness or fluence of the initial spike and that of the subsequent extended emission.

We plot the ratio of the extended emission fluence to that of the initial spike (or limit thereof) adopted from Perley et al. (2008) vs. the specific star formation rate of a subset of short GRBs with known redshift, including GRB 070724A, in the right pane of Figure 10. There is no clear correlation between the events with significant extended emission and the type of galaxy harboring them, although the limited sample size makes the analysis far from conclusive. Despite this, placing such constraints on additional optical and gamma-ray components of short GRBs in galaxies of varying star formation rates and hence stellar populations begins to give us insights into the possible diversity of their progenitors.

Unless short GRBs are eventually found to be accompanied by tell-tale emission features like the supernovae associated with long-duration GRBs, the only definitive understanding of their progenitors will come from possible associations to direct gravitational or neutrino signals. Therefore, continued attempts to observe optical components from processes not associated with the standard afterglow emission from short GRBs remains vital to our understanding of the nature of these events.

The authors would like to note that Berger et al. (2009) reported the detection of a near-infrared and optical afterglow associated with GRB 070724A three weeks after the submission of this paper. They report the detection of the afterglow starting 2.3 hr after the burst at $K_s = 19.56 \pm 0.16$, consistent with our $K$-band upper limit from observations made at roughly the same time. The authors note that the optical to near-IR spectral index is much redder than expected from

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the standard afterglow model, pointing to the possibility of either significant dust exhibition along the line of sight or a non-afterglow origin to the infrared emission. The former scenario, although not supported by our X-ray $N_H$ column density estimates, would explain the lack of an optical afterglow detection and likewise significantly impact our constraints on the LP-SN model parameters.

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