



LJMU Research Online

Smith, RJ, Lucey, JR and Carter, D

The stellar initial mass function in red-sequence galaxies: 1- μ m spectroscopy of Coma cluster galaxies with Subaru/FMOS

<http://researchonline.ljmu.ac.uk/777/>

Article

Citation (please note it is advisable to refer to the publisher's version if you intend to cite from this work)

Smith, RJ, Lucey, JR and Carter, D (2012) The stellar initial mass function in red-sequence galaxies: 1- μ m spectroscopy of Coma cluster galaxies with Subaru/FMOS. *Monthly Notices of the Royal Astronomical Society*, 426 (4). pp. 2994-3007. ISSN 0035-8711

LJMU has developed **LJMU Research Online** for users to access the research output of the University more effectively. Copyright © and Moral Rights for the papers on this site are retained by the individual authors and/or other copyright owners. Users may download and/or print one copy of any article(s) in LJMU Research Online to facilitate their private study or for non-commercial research. You may not engage in further distribution of the material or use it for any profit-making activities or any commercial gain.

The version presented here may differ from the published version or from the version of the record. Please see the repository URL above for details on accessing the published version and note that access may require a subscription.

For more information please contact researchonline@ljmu.ac.uk

<http://researchonline.ljmu.ac.uk/>

The stellar initial mass function in red-sequence galaxies: 1- μm spectroscopy of Coma cluster galaxies with Subaru/FMOS[★]

Russell J. Smith,¹ † John R. Lucey¹ and David Carter²

¹*Department of Physics, Science Laboratories, University of Durham, South Road, Durham DH1 3LE*

²*Astrophysics Research Institute, Liverpool John Moores University, Twelve Quays House, Egerton Wharf, Birkenhead CH41 1LD*

Accepted 2012 August 10. Received 2012 August 7; in original form 2012 June 17

ABSTRACT

To investigate possible variations in the stellar initial mass function (IMF) in red-sequence galaxies, we have obtained infrared spectroscopy with Subaru/Fibre Multi-Object Spectrograph (FMOS) for a sample of 92 red-sequence galaxies in the Coma cluster. Velocity dispersions, ages and element abundances for these galaxies have been previously determined from optical data. The full range of velocity dispersions covered by the sample is $\sigma = 50\text{--}300\text{ km s}^{-1}$. By stacking the FMOS spectra in the rest frame and removing sky-subtraction residuals and other artefacts fixed in the observed frame, we derive composite spectra in the 9600–10500 Å range for galaxies grouped according to their velocity dispersion or Mg/Fe ratio. We measure the Wing–Ford band of FeH and a new index centred on a Ca I line at 10345 Å; these features are strong in cool dwarf stars, and hence reflect the form of the IMF at low mass ($<0.5 M_{\odot}$). The Ca I line, unlike the Wing–Ford band and other ‘classical’ IMF indicators (Na I doublet and Ca II triplet), is unaffected by the abundance of sodium. We compare the measured indices against predictions from spectral synthesis models matched to the element abundances estimated from the optical data. Binning galaxies by velocity dispersion, we find that both the Wing–Ford and Ca I index measurements are best reproduced by models with the Salpeter IMF. There is no clear evidence for an increase in dwarf-star content with velocity dispersion over the range probed by our sample (which includes few galaxies at the highest velocity dispersions, $\sigma > 250\text{ km s}^{-1}$). Binning the observed galaxies instead by Mg/Fe ratio, the behaviour of both indices implies a trend of IMF from Chabrier-like, at abundance ratios close to solar, to Salpeter or heavier for highly α -enhanced populations. At face value, this suggests that the IMF depends on the mode of star formation, with intense rapid starbursts generating a larger population of low-mass stars.

Key words: stars: luminosity function, mass function – galaxies: elliptical and lenticular, cD – galaxies: stellar content.

1 INTRODUCTION

The stellar initial mass function (IMF) is a crucial ingredient in interpreting extragalactic observations: without constraints, or assumptions, for the IMF, observed luminosities cannot be converted into estimates of assembled stellar mass. For old galaxies with a Salpeter (1955) IMF, the great majority of output luminosity is contributed by giant-branch and main-sequence stars with mass $\sim 1 M_{\odot}$. However, some 80 per cent of the stellar mass is locked

up in cool dwarfs with mass $<0.5 M_{\odot}$, which collectively provide less than 10 per cent of the bolometric luminosity. The total stellar mass-to-light ratio is hence very sensitive to the form of the IMF at low masses.

In the Milky Way, resolved star counts indicate that the IMF follows a Salpeter-like power law $dN(M) \propto M^{-2.35} dM$ for $M \gtrsim M_{\odot}$, but becomes shallower at lower masses. This form can be well represented either by a broken power law (e.g. Kroupa, Tout & Gilmore 1993) or by a lognormal distribution (Chabrier 2003). Extensive searches for variation in the IMF in different environments within the Milky Way have not found conclusive evidence for non-universality (e.g. see the review by Bastian, Covey & Meyer 2010). For extragalactic systems, low-mass stars are not directly observable, and the application of a Chabrier-like IMF to distant galaxies must be regarded with caution until empirically tested. Limits on

[★] Based on data collected at Subaru Telescope, which is operated by the National Observatory of Japan.

†E-mail: russell.smith@durham.ac.uk

the contribution of dwarf stars to the mass budget can be inferred from total mass estimates from dynamical modelling (e.g. Bell & de Jong 2001; Cappellari et al. 2012) or gravitational lensing (e.g. Treu et al. 2010), but this approach cannot unambiguously discriminate dark mass in cool dwarfs (bottom-heavy IMF) from that in stellar remnants (top-heavy IMF) or in non-baryonic dark matter.

An alternative method to constrain the mass locked in low-mass stars is to exploit stellar spectral features that depend strongly on surface gravity, and hence distinguish giant from dwarf stars at similar effective temperature. In particular, red optical and near-infrared spectral features that are strong in M dwarfs and not present in cool giants (or vice versa) can potentially discriminate the low-mass stellar content in unresolved galaxies. The Wing–Ford band (WFB) of FeH at 9915 Å (Wing & Ford 1969) and the Na I doublet (8190 Å) were identified as possible giant-to-dwarf star indicators and exploited in early works (e.g. Spinrad & Taylor 1971; Whitford 1977; Cohen 1978; Faber & French 1980; Carter, Visvanathan & Pickles 1986; Couture & Hardy 1993). Some of these papers [e.g. Spinrad & Taylor (1971) and Faber & French (1980), both of which are based on Na I] suggested a substantial fractional contribution of dwarf stars to the integrated light of nearby galaxies, while others [e.g. Whitford (1997), based on the WFB] inferred a much lower dwarf-to-giant ratio. Similarly, Carter et al. reported tension between the two indicators for a sample of 14 galaxies, with Na I requiring a larger fraction of dwarf light than the WFB.

The work of the 1970s to 1990s was hampered by the limitations of spectral synthesis models at that time. In particular, there were no large empirical infrared stellar libraries covering the full range of spectral types relevant to the problem. Moreover, knowledge of the range of elemental abundance variations in galaxies was not yet well constrained; nor was the machinery required to account for such variations through synthetic spectral modelling yet developed. Emphasis in the study of the IMF from resolved populations shifted somewhat to the Ca II triplet (which is strong in giants and weak in dwarfs), with Cenarro et al. (2003) concluding that either calcium was underabundant in giant ellipticals or a dwarf-enriched IMF was required.

A new impetus was given to such studies by the publication of the IRTF spectral library by Rayner, Cushing & Vacca (2009), which finally provided empirical library spectra for cool stars across the red/infrared regime. Building models from these stars to predict the integrated spectra of old populations, and comparing them to observations of eight giant elliptical galaxies, van Dokkum & Conroy (2010) argued that a strongly dwarf-enriched IMF was required to account for the strength of the Na I and (to a lesser extent) the WFB. Their spectra could be fitted only by adopting very bottom heavy IMFs, e.g. a power law with exponent $x = 3$ [where $dN(M) \propto M^{-x} dM$], in which a majority of the mass is in $\lesssim 0.15 M_{\odot}$ stars with $T_{\text{eff}} \lesssim 3000$ K. The stellar mass-to-light ratio M_{*}/L for the IMF favoured by van Dokkum & Conroy is a factor of 3–4 times larger than the Chabrier-like IMF often assumed in extragalactic contexts. Hence verifying this result, and determining the types of galaxies affected, has wide implications. Conroy & van Dokkum (2012a, hereafter CvD12a) have since generalized their spectral synthesis models to make predictions for simple stellar populations (SSPs) with a range of IMFs, ages, metallicities and abundance ratio patterns.

Spurred in part by the original van Dokkum & Conroy (2010) result, several groups have recently investigated the behaviour of IMF-sensitive indices in larger galaxy samples. Spiniello et al. (2012) studied the Na I doublet in luminous red galaxies from the Sloan Digital Sky Survey (SDSS) with velocity dispersions $\sigma =$

200–335 km s⁻¹), finding this feature to increase with σ , requiring a steepening IMF or an increase in Na/Fe enhancement (or both), with increasing galaxy mass. Subsequently, Conroy & van Dokkum have analysed new high signal-to-noise ratio spectroscopy including the WFB, the Na I doublet and the Ca II triplet, as well as blue/optical features, for a sample of 38 nearby galaxies (Conroy & van Dokkum 2012b, hereafter CvD12b; van Dokkum & Conroy 2012). In this work, they apply a sophisticated full-spectrum fitting method with many model components to describe abundance variations and possible confounding parameters. They conclude that dwarf enrichment increases with velocity dispersion and (perhaps more strongly) with [Mg/Fe], ranging from Chabrier-like to heavier-than-Salpeter IMFs. Although all three of the ‘classical’ indicators employed in this work seem to support dwarf-enriched IMFs in the most massive ellipticals, the details depend on which spectral features are included in the fit. Perhaps to compensate for this tension, the models require enhancements of Na/Fe by up to an order of magnitude over the solar ratio. Most recently, Ferreras et al. (2012) used a very large sample of SDSS early-type galaxies to confirm the strong increase in Na I with velocity dispersion in the $\sigma = 100\text{--}300$ km s⁻¹ range, interpreting the results as a steepening of the IMF slope. (They did not explicitly consider the effects of Na/Fe variations.)

For all but the most nearby galaxies, the WFB is redshifted beyond the limits of current ‘optical’ CCDs, and infrared detectors are required. In this paper we present new observations of the WFB for red-sequence galaxies in the Coma cluster, using the infrared Fibre Multi-Object Spectrograph (FMOS; Kimura et al. 2010) at the 8.2-m Subaru Telescope on Mauna Kea. Although galaxies at the distance of Coma are individually fainter than those studied by CvD12b, we can exploit the dense concentration of galaxies in the cluster, and the multiplex capability of FMOS, to observe dozens of galaxies simultaneously. Because the Coma galaxies have a range of radial velocities due to their motions within the cluster potential, stacking their spectra in the rest frame allows rejection of pixels contaminated by OH skyline residuals and other artefacts fixed in the observed wavelength frame. Our galaxy sample spans a range in velocity dispersion and other properties, derived in previous work from optical spectroscopy (Price et al. 2011; Smith et al. 2012), so that we can explore possible variations in the IMF with galaxy mass and control for element abundance effects. Besides the WFB, our spectral range also includes a Ca I line which we exploit as an IMF indicator for the first time.

The outline of this paper is as follows. Section 2 describes the observations and data reduction. The construction of composite spectra is detailed in Section 3. The results are presented in Section 4, starting with the global composite spectrum (Section 4.1), the measurement of spectral indices and their comparison to model predictions (Section 4.2), and then dividing the sample into four bins in velocity dispersion (Section 4.3) and Mg/Fe ratio (Section 4.4), to explore the dependence of the IMF on mass and star formation time-scale. In Section 5, we discuss our results in the context of relevant recent work, and summarize our conclusions in Section 6.

2 OBSERVATIONS AND DATA REDUCTION

We observed Coma cluster galaxies with FMOS on the nights of 2012 May 6–7, using high spectral resolution mode in ‘J-short’ configuration. This setting provides spectral coverage over the range 9130–11 270 Å (corresponding to 8920–11 010 Å at the average redshift of Coma), with resolution 4.6 Å full width at half-maximum

(FWHM), sampled at $1.25 \text{ \AA pixel}^{-1}$. FMOS features two separate (and somewhat different¹) spectrographs, IRS1 and IRS2, each receiving half of the 400 total fibres. The fibres project to 1.2 arcsec in diameter, corresponding to 0.6 kpc at our adopted distance for Coma (100 Mpc, i.e. $H_0 = 72 \text{ km s}^{-1} \text{ Mpc}^{-1}$).

Targets were selected from the sample of 362 Coma galaxies analysed by Smith et al. (2012), which spans a range from the most massive ellipticals ($r_{\text{petro}} \approx 12$, $M_* \approx 10^{12} M_{\odot}$ and $\sigma \approx 300 \text{ km s}^{-1}$) down to the dwarf galaxy population ($r_{\text{petro}} \approx 18$, $M_* \approx 10^9 M_{\odot}$ and $\sigma \approx 30 \text{ km s}^{-1}$). The Smith et al. sample excludes galaxies with significant ongoing star formation using a cut on $H\alpha$ equivalent width. FMOS deploys fibres within a 30 arcmin diameter field of view; however, each fibre can patrol only a small part of the field, so that the ability to pack fibres closely is quite limited. To maximize the number of galaxies observed, we used four fibre configurations with pointings offset by a few arcminutes and different field rotations (constrained by the availability of guide stars). Most galaxies in the sample were observed in more than one configuration, through different fibres, and in many cases with both spectrographs. In total, 121 galaxies were observed at least once. However, the faintest galaxies contribute little signal even if stacked. For the analysis in this paper, we restrict our attention to 92 target galaxies having $\sigma > 50 \text{ km s}^{-1}$.

The observations were made in cross-beam-switching mode, in which two fibres are allocated to each target and the telescope nodded (usually every 900 s) to move the galaxies from one fibre to the other between integrations. This allows the sky background for each target to be determined using the same fibre and detector pixels as for the galaxy light while also keeping each target within a fibre for all of the integrations (equivalent to nodding along the slit in single-object spectroscopy). The fibres in each pair are separated by 90 arcsec, which is sufficiently large to avoid contamination of the sky by the outer part of the target galaxy for all objects in our sample. The integration times for the four fibre configurations were 2.33, 1.50, 2.00 and 2.00 h, for a total of 7.83 h. Observing conditions were fairly good, with seeing 0.6–1.1 arcsec and occasional passing cirrus on the second night.

Initial data reduction was performed using the standard FMOS FIBRE-PAC pipeline (Iwamuro et al. 2012), yielding extracted, calibrated one-dimensional spectra for each fibre. At the redshift of Coma, the WFB is observable in a region of very clean atmospheric transmission. Typical telluric absorption features in this window are less than 1 per cent in amplitude. The FIBRE-PAC pipeline attempts to correct for atmospheric absorption using stars observed simultaneously with the science targets, but this requires correcting for the intrinsic spectral signatures of those stars, in particular the Pa δ line, lying $\sim 150 \text{ \AA}$ shortwards of the redshifted WFB. Despite taking particular care to improve the removal of this feature by manually tuning the spectral type assumed for the star, we were not able to remove these artefacts to our satisfaction. In practice, we opted instead to mask the affected region when stacking the spectra.

3 COMPOSITE SPECTRA

The results of this paper are based upon analysis of composite spectra created by stacking the spectra of many galaxies, after masking ‘bad’ intervals in observed wavelength and shifting to the rest frame. Because the sample spans a significant range in redshift

due to the large virial motions within the cluster (equivalent to $\text{FWHM} \sim 80 \text{ \AA}$), this process fills the ‘gaps’ imposed by residual sky contamination and other artefacts which are fixed in the observed wavelength frame.

FMOS features an OH suppression mask to suppress the wavelength regions affected by strong airglow lines. However, these regions must still be masked in software, since they otherwise appear as low-flux notches within the extracted spectra. The masks in the two spectrographs are different, with IRS1 masking many more weak lines than IRS2. Moreover, during the reduction of our data, we discovered that the mask in IRS2 was misaligned, causing masking to occur at wavelengths $\sim 60 \text{ \AA}$ shortwards from the target sky lines.

We impose software masking at the locations of the OH-suppression notches (in the case of IRS2, at their actual location, not at the position of the sky lines they were intended to mask, which are adequately subtracted in the reduced spectra). Additionally, we mask a 20- \AA region centred on (unredshifted) Pa δ to remove any residuals arising from poor atmospheric absorption corrections. After applying the mask, we correct the spectra to the rest frame, and rebin all of the data to a common wavelength scale of 9600–10 500 \AA , sampled at 1.25- \AA intervals (the native sampling of the original spectra). These spectra, of which there are up to four per galaxy, form the input to the compositing procedure.

To combine the spectra, we first apply a simple median over the inputs to obtain an initial estimate for the spectral shape. Each spectrum is then matched to this median using a fourth-order polynomial correction to remove residual throughput variations. After this we identify and reject outlying pixels (beyond $\pm 5\sigma$ from the average over all input spectra), and finally combine surviving pixels using a mean, weighted by the inverse of the variance spectrum. (For the latter, we use the background noise spectrum estimate from FIBRE-PAC together with a component representing Poisson noise in the object spectra.)

Errors in the composite spectra were estimated both by using the formal error on the weighted mean and by computing the composite using 100 bootstrap-resampled realizations of the set of input spectra. The two error estimates agree to within 5 per cent in the region of the WFB and within 20 per cent across the whole wavelength range. However, the rebinning of the spectra on to a common rest-frame wavelength scale inevitably introduces correlations in the errors between adjacent pixels. In our analysis we will generally make use of the bootstrap realizations of the composite spectrum to propagate the uncertainties appropriately.

In constructing the stacked spectra, we combine galaxies with a range of different velocity dispersions. Broadening all the spectra to the largest velocity width is not appropriate for these data, because it would spread systematic errors (e.g. from sky-subtraction residuals) and missing data into most of the pixels in each spectrum. In our subsequent analysis, we account for the effects of velocity broadening by convolving model spectra with a gaussian kernel to match the weighted mean velocity dispersion of galaxies in the stacks.²

² We have confirmed all of our results using a non-gaussian convolution kernel which accounts explicitly for the range of velocity dispersions entering the stack, with their corresponding weights. All measured line-strength indices are insensitive to which method is used, to within one-tenth of the quoted measurement error. We quote results from the simpler scheme throughout the paper.

¹ See Kimura et al. (2010).

4 RESULTS

4.1 Global massive-galaxy composite

We begin by examining qualitatively the spectrum created by stacking the 59 galaxies with $\sigma > 100 \text{ km s}^{-1}$. Within this sample, the mean velocity dispersion, weighted according to the contribution of each galaxy to the stack, is 180 km s^{-1} . The velocity dispersion range $\sigma = 125\text{--}250 \text{ km s}^{-1}$ contributes 90 per cent of the total weight. Using line-strength indices in the optical range (4000–5500 Å), Smith et al. (2012) derived metallicities and abundance estimates through comparison with the SSP models of Schiavon (2007). The average optically derived parameters, again weighted by contribution to the stacked FMOS spectrum are the following: $[\text{Fe}/\text{H}] = 0.00$, $[\text{Mg}/\text{Fe}] = +0.25$ and SSP-equivalent age 10 Gyr.

Fig. 1 shows the composite spectrum over the full spectral range observed. The total signal-to-noise ratio of this spectrum is $\sim 170 \text{ \AA}^{-1}$. A visual comparison with the SSP models of CvD12a shows that the observed features correspond closely to the expected absorption signatures of old stellar populations. We restrict all further analysis to the 9600–10500 Å spectral range which includes the IMF-sensitive WFB and Ca I line.

In Fig. 2 we divide the same set of input data into two, computing composite spectra for each of the FMOS spectrographs separately. Comparison between the IRS1 and IRS2 composites shows that despite their different characteristics (in particular the much heavier OH suppression mask in IRS1), there is excellent agreement between the average spectra obtained from the two spectrographs.

In Fig. 3, we compare the observed composite spectrum against SSP models from CvD12a, with old ages (13.5 Gyr), solar Fe/H and mildly enhanced α/Fe ratio. The publicly available CvD12a models cover either enhanced α/Fe ratio at 13.5 Gyr or solar abundance ratios at younger ages. We choose in this figure to match the enhanced α/Fe ratio rather than the slightly younger ages which may be relevant to our sample. (Note that for realistic, non-SSP star formation histories, the ages derived from optical spectra will be more affected by recent episodes of star formation, while the infrared spectra analysed here will be dominated by the oldest populations.) Each model has been broadened to match the weighted-mean velocity dispersion for the stack, and matched to the overall shape of the observed spectrum using a fourth-order polynomial correction. The smoothing scale has been verified by a direct χ^2 -fitting of the models to the observed spectrum; in practice, we prefer to fix this to the average velocity dispersion which is securely measured from optical data.

Qualitatively, there is very good agreement between observed and model spectra in regions that are not sensitive to the IMF, with all the main spectral features well reproduced. A systematic discrepancy is seen in the feature at $\sim 10050 \text{ \AA}$, which includes the Pa δ line as well as TiO absorption, with stronger absorption observed than in the models, causing residuals at the ~ 1 per cent level over an $\sim 50 \text{ \AA}$ interval. (A very similar discrepancy is also evident in fig. 1 of CvD12b.) More localized 1 per cent discrepancies are also seen at $\sim 9700\text{--}9800 \text{ \AA}$.

In the WFB region, the bottom of the absorption feature lies close to the models with Chabrier or Salpeter IMF. However, it should be noted that the spectral region immediately redwards of the WFB

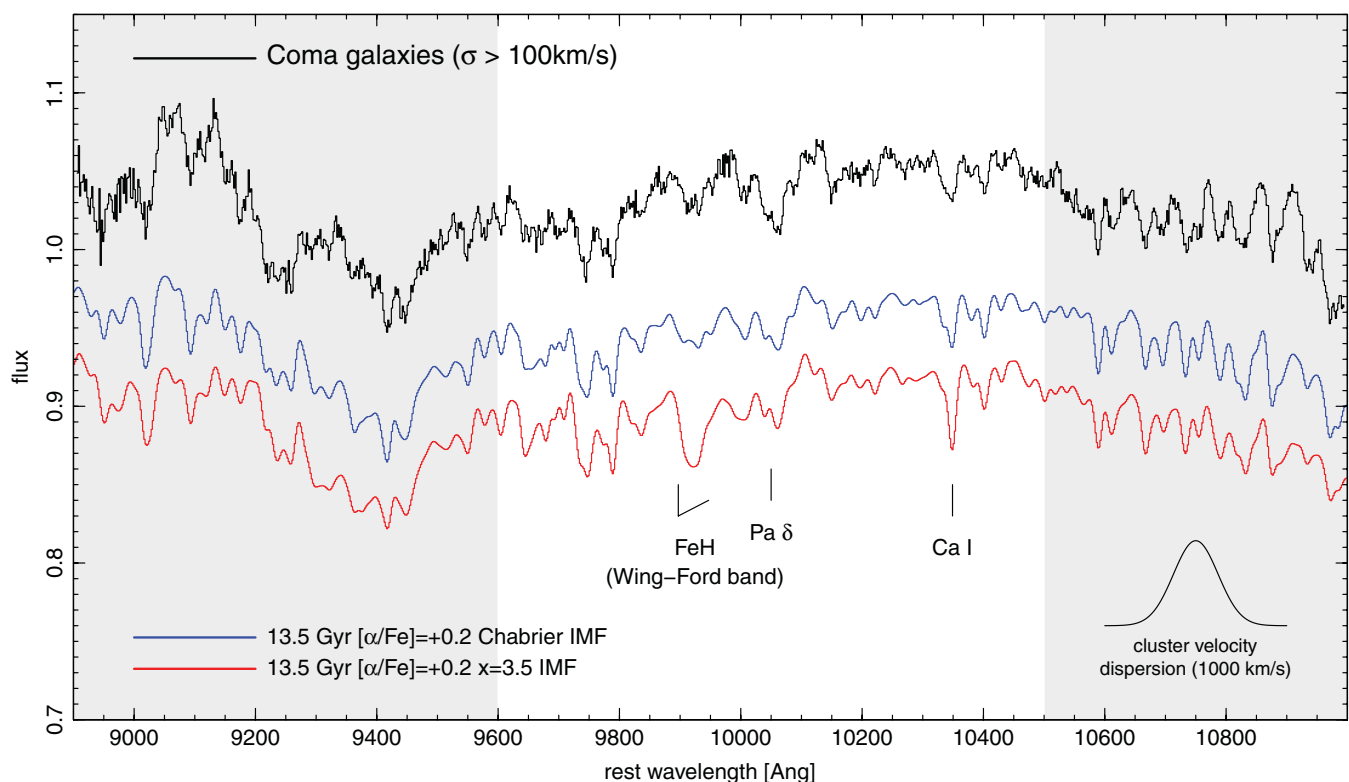


Figure 1. Composite spectrum constructed from the 59 observed galaxies with $\sigma > 100 \text{ km s}^{-1}$ (top, black). The lower spectra show SSP models from CvD12a, with different IMFs, smoothed to match the effective velocity dispersion of the stacked data. The $x = 3.5$ IMF is extremely bottom heavy and is plotted here to emphasize the IMF dependence of the WFB and Ca I line. Subsequent figures show only on the rest-frame 9600–10500 Å region (white background) where the atmospheric transmission is clean and which includes the IMF-sensitive spectral features. For reference, we show a Gaussian profile representing the shifts induced by the velocity dispersion within the cluster; any residual sky-subtraction artefacts are effectively smoothed over this scale in the composite spectrum.

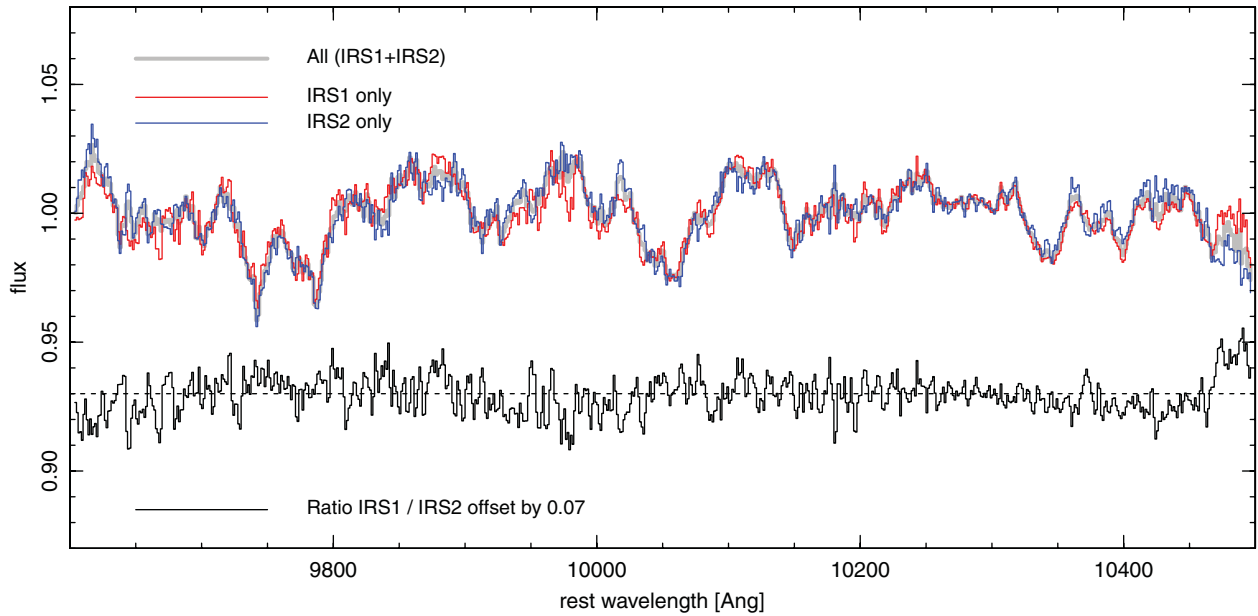


Figure 2. Comparison between composite spectra computed using each of the two FMOS spectrographs, IRS1 and IRS2. We have combined all $\sigma > 100 \text{ km s}^{-1}$ galaxies observed on each spectrograph; the two sets of raw data are independent, though some galaxies contribute to both stacks.

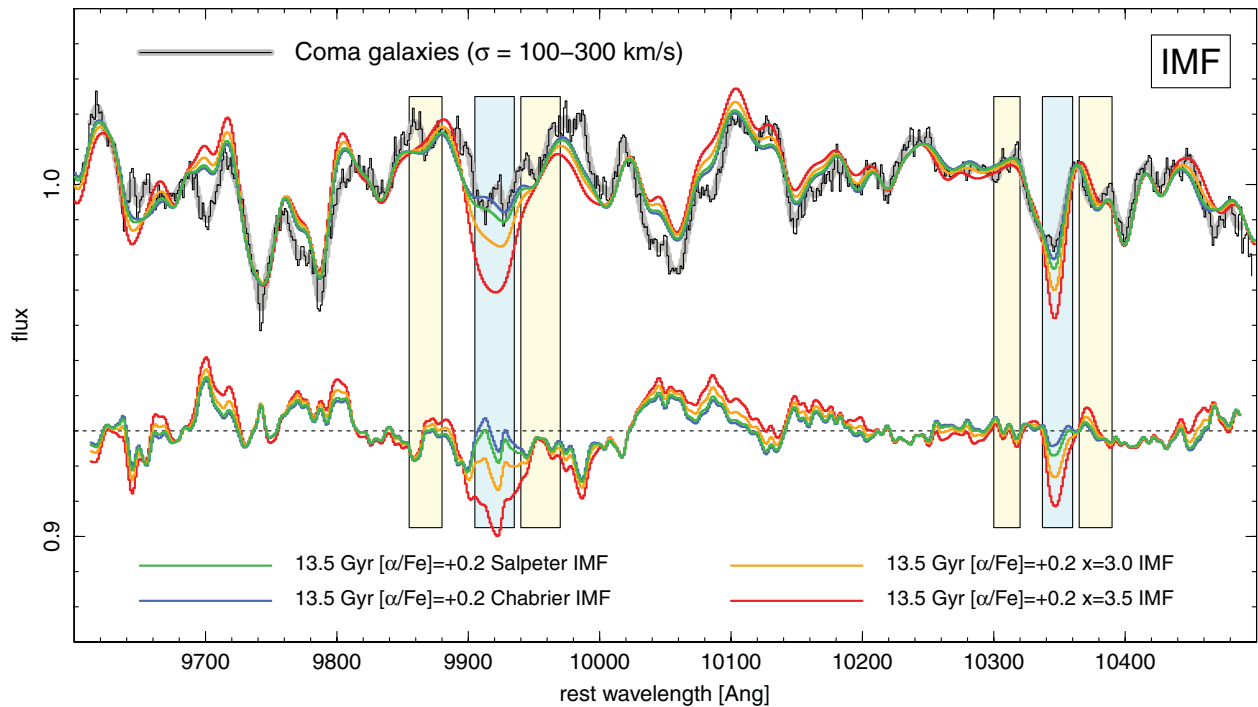


Figure 3. Composite spectrum constructed from the 59 observed galaxies with $\sigma > 100 \text{ km s}^{-1}$ (thin black line). The thick grey line shows the same spectrum smoothed by an amount corresponding to $\sigma \approx 100 \text{ km s}^{-1}$ to suppress small-scale noise while retaining the intrinsic resolution of the spectrum. The coloured lines are 13.5-Gyr models from CvD12a, with slightly enhanced α/Fe ratio (matching the optically determined Mg/Fe), and a range of different IMFs. The lower section shows the flux ratio (model/data) for the same set of models, offset by 0.07. The vertical bands show the feature and pseudo-continuum bandpasses defining the FeH0.99 index of CvD12a and our newly defined CaI1.03 index. Note that the red pseudo-continuum of FeH0.99 is higher in the data than the models; hence, although the ‘floor’ of the absorption matches the Chabrier IMF model, the measured index strength is closer to that of the Salpeter model (see Section 4.2).

is poorly matched by all models. Hence, inferring the IMF through this comparison is somewhat degenerate with the treatment of the continuum matching or, in the case of spectral indices, with the placement of the pseudo-continuum bandpasses. Furthermore, the detailed shape of the absorption band itself does not agree well

with any of the CvD12a models: the observed spectrum shows a bump at 9920 \AA , while the models at this resolution show a flat-bottomed absorption. This discrepancy does not seem to arise from a resolution mismatch; reducing the assumed velocity dispersion to introduce more structure in the WFB results in an unacceptable

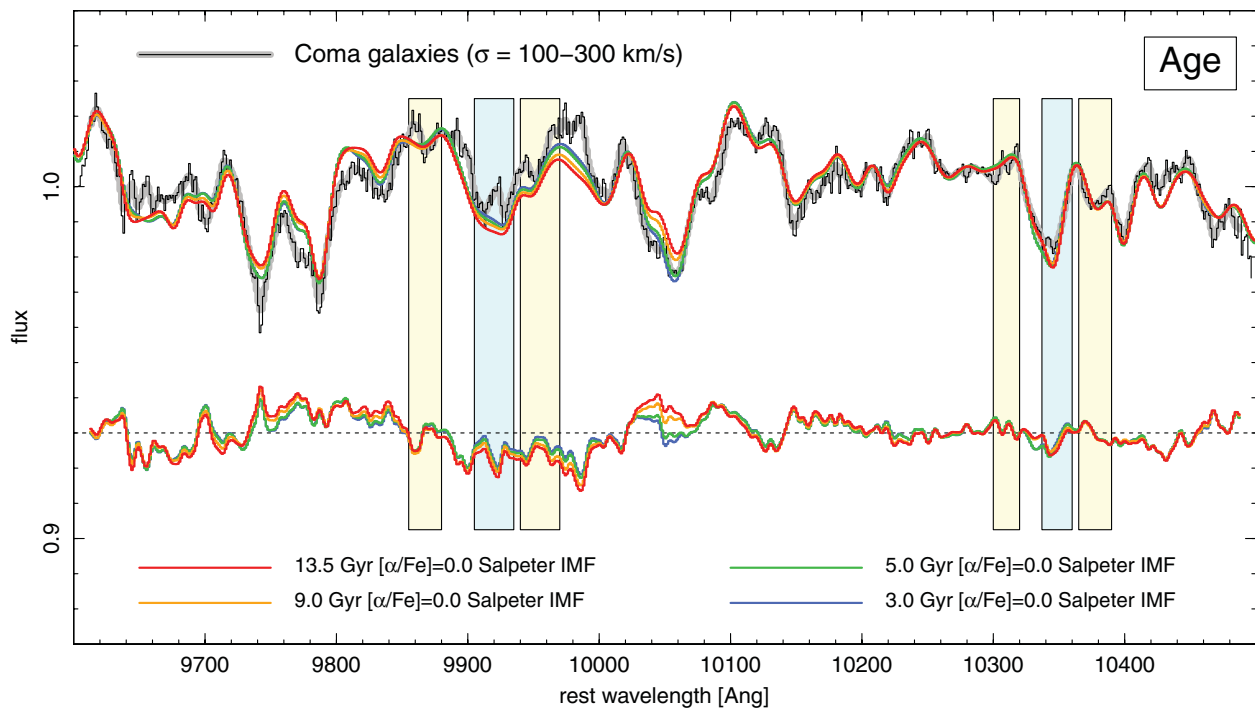


Figure 4. Same as Fig. 3, but compared to models for a range of different ages, all having solar α/Fe ratio and Salpeter IMF. The figure demonstrates the robustness of the models against age variations (except around Pa δ at $\sim 10\,050\text{ \AA}$). Note that none of these models has α/Fe appropriate to the observed galaxy sample.

fit elsewhere in the spectrum, and would be incompatible with the optically measured σ .

At $10\,345\text{ \AA}$ there is another absorption feature which is sensitive to the IMF, according to the models. This feature is marked, as Ca I, among other possible IMF indicators, in fig. 10 of CvD12a, but was not pursued further in that paper nor (to our knowledge) in any other previous work. The Ca I line is in the red wing of a blend, with lines of Fe I and Sr II contributing to the blue wing (identified from the Arcturus atlas of Hinkle, Wallace & Livingston 1995). The observed spectrum shows weak absorption at Ca I, apparently compatible with the Chabrier/Salpeter IMF, and inconsistent with the dwarf-enriched models. The immediate environment of this feature is well fitted by the models.

According to the CvD12a models, the $1\text{-}\mu\text{m}$ spectra depend only weakly on age and element abundances, as we show in Figs 4 and 5. Fig. 4 compares the same observed composite spectrum against the CvD12a models of varying age, all having the Chabrier IMF and solar α/Fe . Relative to the oldest model, younger ages reduce the systematic discrepancies, but the improvement is substantial only in the $\sim 10\,050\text{ \AA}$ region (i.e. near Pa δ), and ages as low as 3–5 Gyr would be required to match the strong absorption here. Since the mean optical SSP-equivalent age is ~ 10 Gyr, and the equivalent age in the infrared should be older than in the optical,³ this does not seem to be a tenable solution, and may point to some incompleteness in the model. Fig. 5 similarly compares the observed composite spectrum against models with varying α/Fe abundance ratio. Increasing α/Fe strengthens the absorption at 9830 \AA , improving the agreement immediately bluewards of the WFB, but leads to weaker the

absorption at $\sim 9700\text{ \AA}$, increasing the discrepancy with the data here.

4.2 Absorption indices and inferred IMF

To quantify our results, we compute ‘Lick-style’ line-strength indices for the WFB and Ca I features. For the WFB, we adopt the index defined as FeH0.99 in CvD12a: feature band at $9905\text{--}9935\text{ \AA}$ and pseudo-continua at $9855\text{--}9880$ and $9940\text{--}9970\text{ \AA}$. To measure Ca I, we define a new index (CaI1.03) with feature band at $10\,337\text{--}10\,360\text{ \AA}$ and pseudo-continua at $10\,300\text{--}10\,320$ and $10\,365\text{--}10\,390\text{ \AA}$. The feature band is carefully placed to isolate the IMF-sensitive Ca I line, rather than the blended neighbouring absorption, and to minimize dependence on age and α/Fe ratio, through examination of Figs 3–5. When measuring the indices on the CvD12a models, we first match these models to the observed spectrum as before, using the weighted mean velocity dispersion of the galaxy stack, and allowing a fourth-order continuum correction to be fitted independently for each model. The results are insensitive to the order of the correction polynomial applied, within reasonable limits.

The index measurements made on the global composite spectrum are summarized in Table 1, and are compared to the model predictions in Fig. 6. The index errors are obtained by performing the index measurement on the bootstrap realizations of the composite spectrum.

The measured FeH0.99 is consistent with the Chabrier or Salpeter IMF for models with solar abundance mixtures. For higher assumed α/Fe ratios, more dwarf-enriched models are brought into closer agreement with the data. For the mildly α -enhanced model favoured by the optical data, the measured FeH0.99 is consistent with a Salpeter IMF, marginally inconsistent with Chabrier (at the 2.2σ level) and incompatible with the $x = 3$ model (at $>3\sigma$).

³ In principle, large dust-obscured young subpopulations could violate this expectation in some galaxies, but this seems unlikely to affect our sample *on average*.

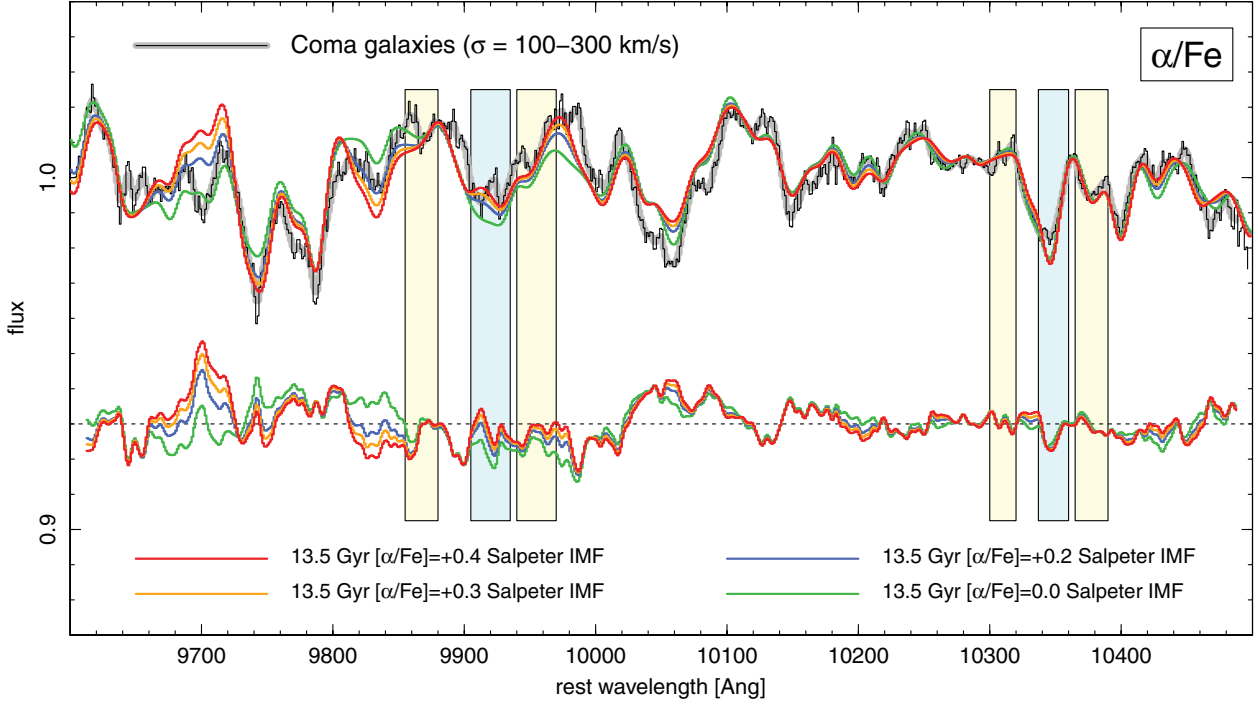


Figure 5. Same as Fig. 3, but compared to models for a range of different α/Fe ratios, all having an age of 13 Gyr and Salpeter IMF.

Table 1. IMF-sensitive line-strength indices measured from stacked spectra after binning as a function of velocity dispersion and Mg/Fe ratio. The first column indicates the parameter used to define the bin, while min and max are the limits for that parameter. N_{gxy} is the number of galaxies in the stack. Velocity dispersion σ is in units of km s^{-1} . We indicate the average SSP-equivalent stellar population parameters for each bin from the optical spectroscopy of Smith et al. (2012). These averages are weighted according to the contribution of each galaxy to the stacked spectrum.

Bin	Min	Max	N_{gxy}	$\langle \sigma \text{ (km s}^{-1}\text{)} \rangle$	$\langle [\text{Fe}/\text{H}] \rangle$	$\langle [\text{Mg}/\text{Fe}] \rangle$	$\langle [\text{Ca}/\text{Fe}] \rangle$	$\langle \text{Age (Gyr)} \rangle$	FeH0.99 (\AA)	CaI1.03 (\AA)
$\sigma \text{ (km s}^{-1}\text{)}$	100	300	59	180	+0.00	+0.25	+0.07	10.2	0.452 ± 0.055	0.274 ± 0.024
$\sigma \text{ (km s}^{-1}\text{)}$	50	130	50	109	+0.01	+0.20	+0.03	7.1	0.419 ± 0.053	0.339 ± 0.030
$\sigma \text{ (km s}^{-1}\text{)}$	130	170	18	150	+0.01	+0.19	+0.04	9.7	0.454 ± 0.045	0.279 ± 0.024
$\sigma \text{ (km s}^{-1}\text{)}$	170	210	15	190	-0.07	+0.31	+0.10	12.1	0.331 ± 0.053	0.311 ± 0.028
$\sigma \text{ (km s}^{-1}\text{)}$	210	300	9	246	+0.05	+0.27	+0.08	10.7	0.530 ± 0.053	0.231 ± 0.030
[Mg/Fe]	-0.05	+0.15	34	114	+0.08	+0.12	-0.03	6.9	0.391 ± 0.061	0.267 ± 0.032
[Mg/Fe]	+0.15	+0.22	21	168	+0.04	+0.19	+0.05	8.8	0.453 ± 0.044	0.314 ± 0.021
[Mg/Fe]	+0.22	+0.30	16	188	+0.01	+0.27	+0.07	9.4	0.529 ± 0.056	0.259 ± 0.029
[Mg/Fe]	+0.30	+0.46	21	201	-0.12	+0.36	+0.13	13.7	0.331 ± 0.049	0.300 ± 0.027

Only by adopting $[\alpha/\text{Fe}] = +0.4$ (which is inconsistent with $[\text{Mg}/\text{Fe}]$ measured from optical spectroscopy) can the $x = 3$ model be brought within 2σ of the observed index. Not unexpectedly, the predicted FeH0.99 is sensitive to the iron abundance, but this is measured to be close to solar from the fit to the optical indices. The index also depends on the sodium abundance, which is not constrained from our optical spectra. For the substantial sodium enhancements (up to +1.0 dex) suggested by the results of CvD12b, the predicted FeH0.99 could be reduced by 0.2 dex, which would yield agreement with the $x = 3$ model.

As noted above, casual examination of Fig. 3 might suggest a better match for Chabrier models than for Salpeter models, in apparent disagreement with the index measurements. In fact, the index measurements implicitly force a more local continuum matching than the low-order polynomial correction we applied for the spectral comparison. Because the models lie lower than the observed spectrum in the red pseudo-continuum band of FeH0.99, the predicted index is reduced relative to the observed value. Fig. 7 (left-hand

panel) demonstrates this by plotting the WFB region after imposing a linear correction to normalize the pseudo-continua to unity in models and data alike. With this normalization it becomes clearer how Salpeter (or slightly more dwarf-enriched) models yield the closest agreement in the index value. However, this figure also highlights the failure of all models to match the detailed shape of the observed spectrum through the WFB region.⁴

The measured CaI1.03 index for the observed composite spectrum is consistent with the Chabrier IMF, for either the solar abundance ratio or the α -enhanced models, and marginally inconsistent (2σ) with the Salpeter models. The dwarf-enriched $x = 3$ models are excluded at the $>5\sigma$ level (see also Fig. 8, left-hand panel). However, the CvD12a α -enhanced models assume that calcium, being an α -element, is enhanced in lock-step with magnesium,

⁴ The same discrepancy in the red pseudo-continuum of FeH0.99 is seen in the van Dokkum & Conroy (2012) sample (their fig. 8).

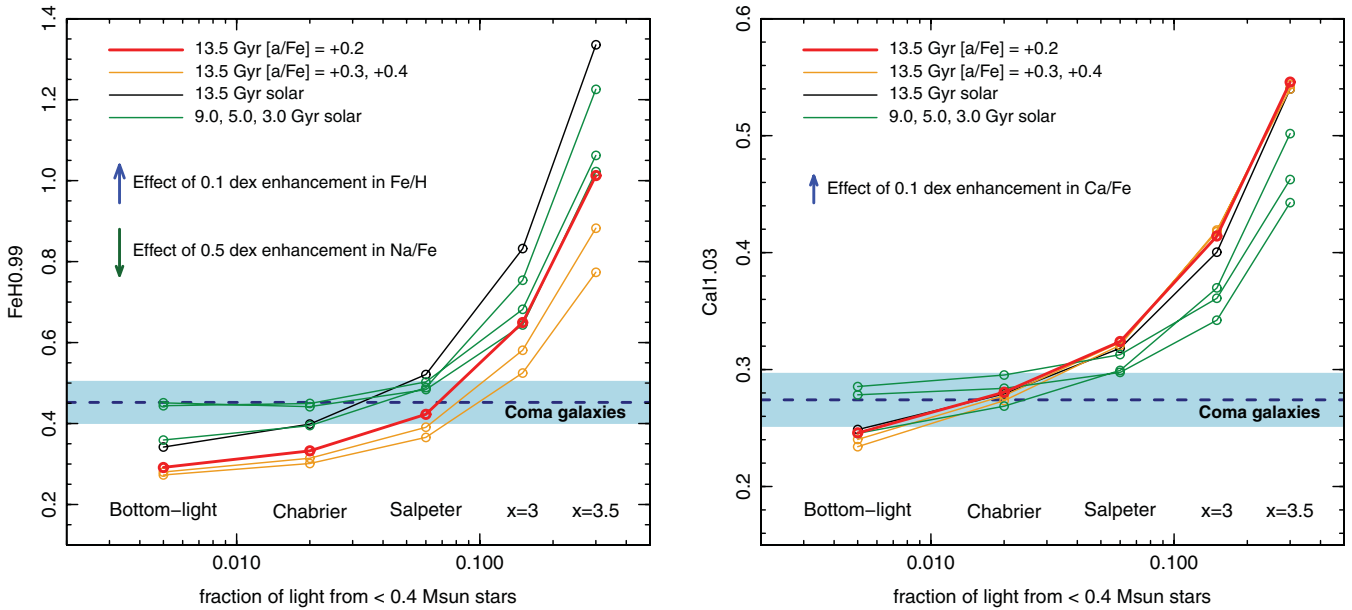


Figure 6. Absorption indices measured from the global composite (dashed line and blue-shaded 1σ error interval), compared to indices measured on the CvD12a model spectra. The WFB index FeH0.99 is defined as in CvD12a, while CaI1.03 is newly defined in this paper (see text). The indices are measured on the models after smoothing to the mean velocity dispersion appropriate to the data, and continuum matching with a low-order polynomial correction. Vectors indicate how the model predictions for FeH0.99 depend on Fe/H and Na/Fe, and how CaI1.03 is affected by Ca/Fe.

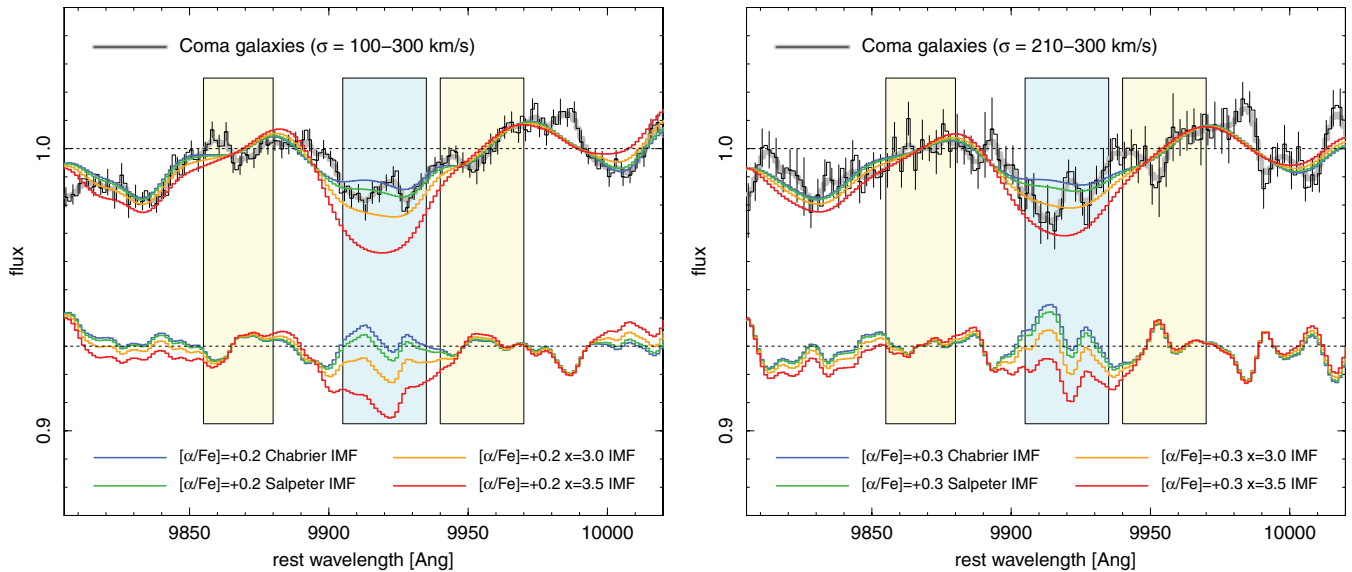


Figure 7. The WFB region of the composite spectra for the global composite and for the highest σ bin. The spectra have been normalized to unity at the FeH0.99 pseudo-continuum bands. Bootstrap-derived error bars are shown on every second pixel.

titanium, etc. Optical spectroscopic indices (primarily Ca4227) imply instead that the Ca/Fe ratio is lower than Mg/Fe on average (e.g. Vazdekis et al. 1997; Thomas, Maraston & Bender 2003; Smith et al. 2009), in which case the predicted CaI1.03 will be too high. For the $\sigma > 100 \text{ km s}^{-1}$ galaxies in our global composite, the contribution-weighted average [Ca/Fe] is +0.07, compared to [Mg/Fe] = +0.25. We estimate the effect of this by noting that for a Chabrier IMF (for which CvD12a provide a variety of models with variation in individual element abundances), an increase by 0.15-dex change in Ca/H increases the CaI1.03 index by 0.03 Å. Applying this as a linear correction to the indices predicted by the $[\alpha/\text{Fe}] = +0.2$ model, the Chabrier and Salpeter models are both consistent with

the data at the $\sim 1\sigma$ level. The $x = 3$ models would remain excluded ($> 4\sigma$).⁵

We conclude that for Coma red-sequence galaxies with $\sigma > 100 \text{ km s}^{-1}$, the strengths of the WFB and the IMF-sensitive Ca I line at 10345 Å each support an IMF with dwarf-star content similar to the Salpeter case, once the known abundance information is taken

⁵ Strictly of course, the low Ca/Mg ratios were derived for Salpeter IMF in comparison with the Schiavon (2007) models, rather than self-consistently with varying IMF. Figs 13 and 14 of CvD12a tend to confirm that Ca4227 and Mg5177 are much less sensitive to the IMF than to abundance variations, so this should not cause any serious bias.

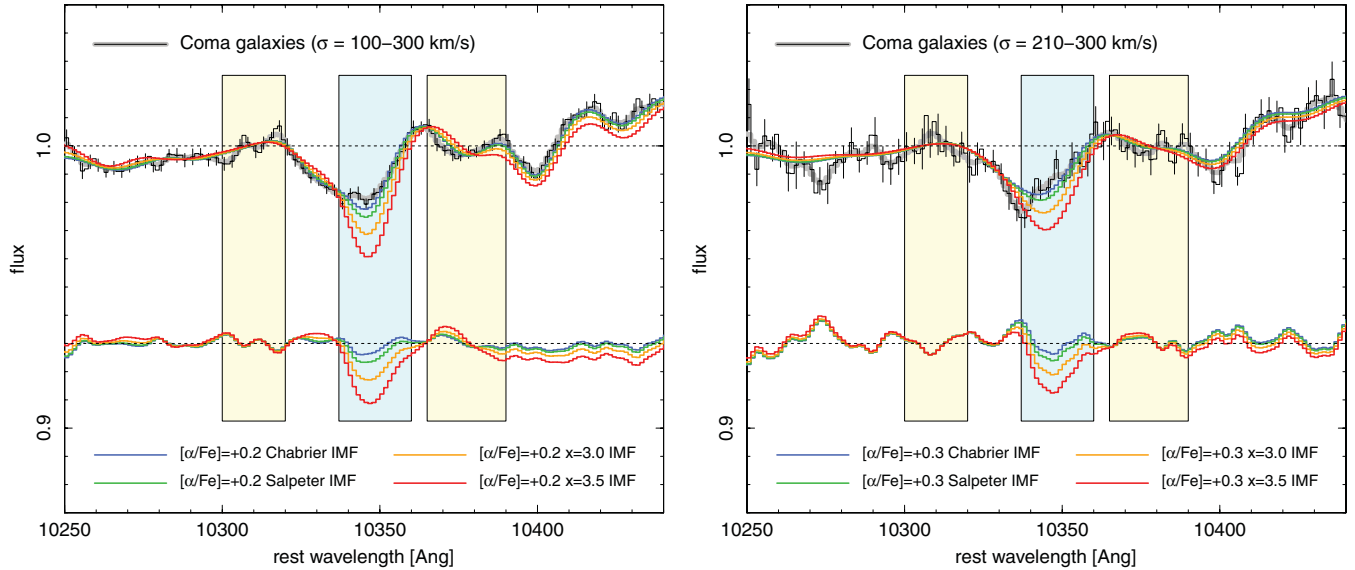


Figure 8. Same as Fig. 7, but for the CaI1.03 index.

into account. The concordance between the two indicators suggests that the average Na/Fe ratio, which influences the WFB, cannot be too large: $[\text{Na}/\text{Fe}] = +1.0$ would violate the consistency between FeH0.99 and CaI1.03, for *any* IMF, at the $>2.5\sigma$ level.

4.3 Dependence on velocity dispersion

Most properties of red-sequence galaxies correlate with proxies for their mass, and in particular with internal velocity dispersion, σ . In general, galaxies with larger σ have older SSP-equivalent ages and higher abundances of α -elements (including magnesium, calcium and titanium), and of carbon and nitrogen, although the iron abundance is flat or even declining with σ (e.g. Nelan et al. 2005; Graves et al. 2007; Smith et al. 2009; Johansson, Thomas & Maraston 2012). Recent work (CvD12b; Ferreras et al. 2012; Spiniello et al. 2012) suggest that the IMF may also depend on galaxy mass. For example, CvD12b find Chabrier-like IMFs favoured on average at $\sigma \approx 150 \text{ km s}^{-1}$, Salpeter-like IMFs at $\sigma \approx 250 \text{ km s}^{-1}$ and steeper still for the most massive galaxies. In this section we divide the FMOS sample according to velocity dispersion to examine the mass dependence of the indices and derived IMF for our sample.

In Fig. 9, we show composite infrared spectra computed for galaxies in four ranges of velocity dispersion. The bin boundaries were defined to yield a similar signal-to-noise ratio of $\sim 90 \text{ \AA}^{-1}$ in each stack. For these smaller samples, the stacked spectra show occasional noise spikes which are probably due to incomplete rejection of sky-subtraction residuals. It is clear that four is as many bins as can be justified by the quality of the present data. The effective mean velocity dispersion, age and abundances (weighted according to the contribution of each galaxy to the stack) for each bin are given in Table 1. Qualitatively, we find only weak variations in the spectra as a function of σ , with all the main features of the spectra reproduced in all four composites. There is a hint for slightly stronger WFB in the most massive bin, but lower in the second most massive bin, and hence no clear trend is apparent. The low σ galaxies appear to show stronger Ca I.

Fig. 10 shows the FeH0.99 and CaI1.03 indices measured from the σ -binned stacks, in comparison with the global composite

and with the CvD12a model predictions. Direct comparisons of observed and model spectra for the highest σ bin are shown in Figs 7 (right-hand panel) and 8 (right-hand panel) for FeH0.99 and CaI1.03, respectively. The predicted indices are derived after broadening the models to the contribution-weighted velocity dispersion of each stack, i.e. the Doppler broadening correction is applied to the models, rather than to the data. For the predicted indices, we track the variation in α/Fe ratio as a function of σ by interpolating between the CvD12a models, with $[\alpha/\text{Fe}] = 0.0, +0.2, +0.3$ and $+0.4$, to match the weighted-mean observed Mg/Fe. For FeH0.99, the models are adjusted to track the average measured Fe/H in each bin, and the predictions for CaI1.03 account for the average measured Ca/Fe.⁶

The behaviour of the measured FeH0.99 index in Fig. 10 is not monotonic, with an apparently significant dip in the $\sigma \approx 190 \text{ km s}^{-1}$ bin, as expected from the spectra. Comparison with the constant-IMF model predictions shows that this behaviour is expected, once allowance is made for the variation in Fe/H (by chance, this bin has lower average iron abundance than any other). With this correction, all bins with velocity dispersion $<210 \text{ km s}^{-1}$ are within 1σ from the predictions consistent with a common Salpeter IMF. The highest velocity dispersion bin has slightly higher FeH0.99 (1.5σ from the Salpeter prediction), but is still marginally inconsistent (2σ) with the $x = 3$ model.

The CaI1.03 index is consistent with Salpeter models across the whole range in σ , once corrected for the measured Ca/Fe. (If instead we were to assume $\text{Ca}/\text{Fe} = \text{Mg}/\text{Fe}$, the results favour a Chabrier IMF.) There is no evidence in this index for increased dwarf content at high-velocity dispersion: the $\sigma > 210 \text{ km s}^{-1}$ bin is consistent with Salpeter or Chabrier IMF, and inconsistent with the $x = 3$ model at the 3.6σ level.

As discussed above, the conclusions for the WFB are somewhat sensitive to the sodium abundance which is not well constrained from other data. If we had assumed that Na/Fe increases with velocity dispersion, then we would have recovered a trend of increasing

⁶ These differential corrections were derived from Chabrier IMF models and applied to all other IMFs.

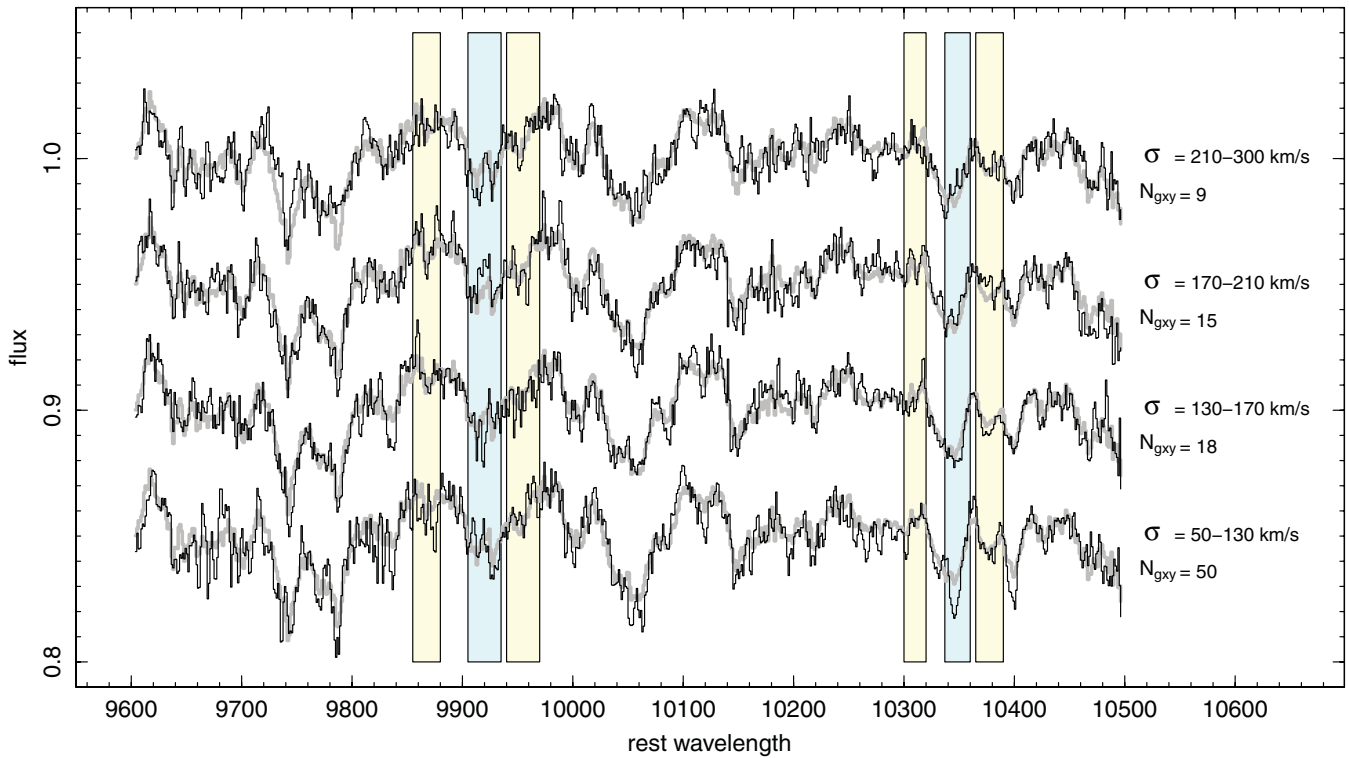


Figure 9. Composite spectra computed in four bins of velocity dispersion (black lines). In each case, the thick grey line shows the global composite of $\sigma > 100 \text{ km s}^{-1}$ galaxies for comparison. No smoothing has been applied to either the σ split or the global composite.

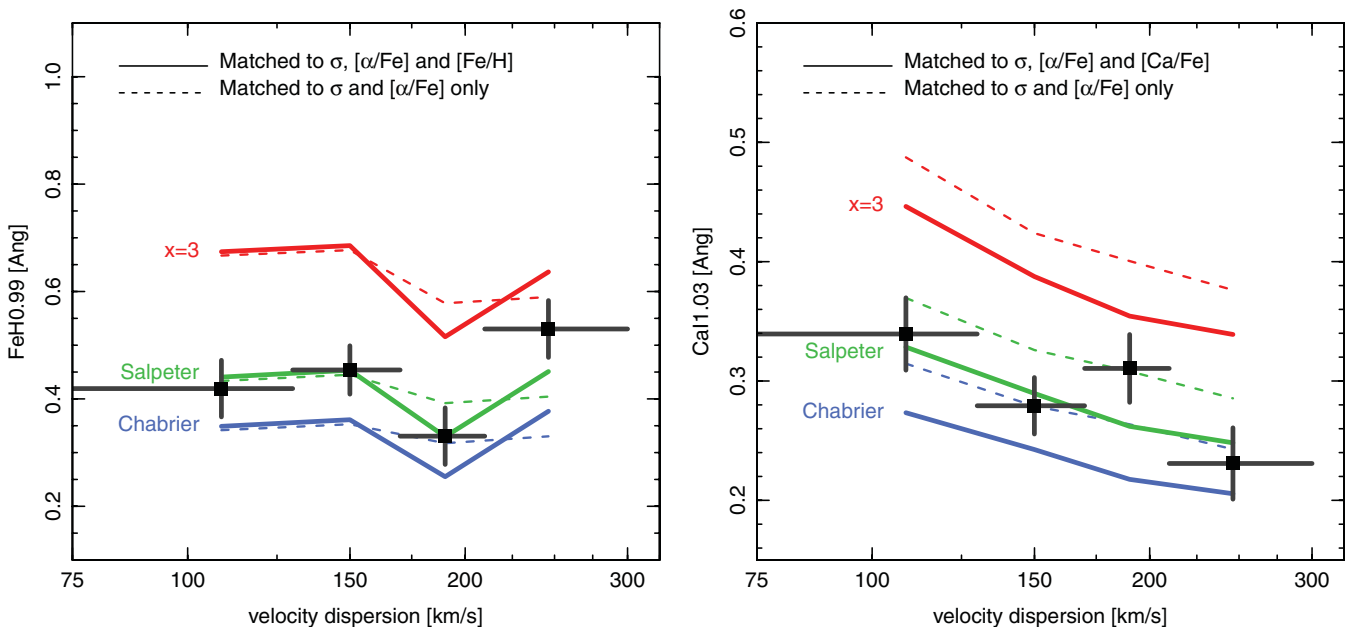


Figure 10. Variation in the FeH0.99 and CaI1.03 indices as a function of velocity dispersion. Bins in σ are as in Fig. 9. The coloured lines show the same indices measured from constant-IMF Cvd12a models tuned to match the properties of each bin. The models are smoothed to the contribution-weighted mean velocity dispersion, interpolating between models of different α/Fe to match the (similarly averaged) Mg/Fe . In the left-hand panel, the solid lines include a correction to the measured Fe/H (the dashed lines assume solar Fe/H). In the right-hand panel, Fe/H is assumed solar (the CaI1.03 index is insensitive to the iron abundance). Here, the dashed lines assume calcium behaves as a generic α element, i.e. $\text{Ca}/\text{Fe} = \alpha/\text{Fe} = \text{Mg}/\text{Fe}$. The solid line includes a correction to Ca/Fe as derived from the optical spectra; applying this correction changes the preferred IMF from Chabrier to Salpeter.

dwarf enrichment with σ from FeH0.99, though CaI1.03 would be unaffected.

4.4 Correlation with Mg/Fe ratio

CvD12b find that the apparent increase in dwarf content in massive ellipticals may be better correlated with the Mg/Fe ratio, rather than with velocity dispersion (their fig. 5). Such a correlation would be very interesting, as Mg/Fe is generally interpreted as an indicator for the time-scale of star formation (Thomas et al. 2005), and thus reflects the physical conditions in galaxies during their major formation epoch.

We test this result by binning our spectra according to (optically determined) Mg/Fe (Figs 11–12 and Table 1). For the WFB we measure an increase in the FeH0.99 index strength with increasing Mg/Fe, for the bins with $[\text{Mg}/\text{Fe}] < 0.3$. The models show that for constant IMF we would instead expect the index to fall with increased Mg/Fe (reduced Fe/H and increased velocity dispersion, as well as the direct effect of α/Fe). The inferred IMF drifts from Chabrier (or lighter) at $[\text{Mg}/\text{Fe}] = +0.12$ to Salpeter at $[\text{Mg}/\text{Fe}] = +0.19$ and close to $x = 3$ at $[\text{Mg}/\text{Fe}] = +0.27$. For the $[\text{Mg}/\text{Fe}] > 0.3$ bin, a lower index value is measured, but the recovered IMF would still be marginally heavier than Salpeter for the Fe/H-corrected models. The CaI1.03 index shows similar behaviour: the index value is fairly constant with Mg/Fe, within the errors, but fixed-IMF models predict a declining trend (mainly due to increased velocity broadening as the high Mg/Fe bins have higher σ on average). Hence, the inferred IMF becomes more dwarf enriched with increasing Mg/Fe, ranging from Chabrier to slightly heavier than Salpeter.

The observed trends with Mg/Fe are clearly only marginally significant in either index individually. Taken in combination, how-

ever, the similar pattern derived from the two indices lends tentative support to the claim that α -enhanced populations formed in rapid bursts harbour more low-mass stars than galaxies with extended star formation histories.

5 DISCUSSION

In this section we focus on comparing our results to those presented (during completion of this paper) by CvD12b. Their study is directly comparable to ours in probing a similar range in galaxy properties (although they have only two or three galaxies in the range covered by our lowest σ bin of 50–130 km s^{-1}), and in using optical spectra to constrain additional parameters (age and element abundances) so that possible IMF variations can be investigated.

Our analysis differs from that of CvD12b first in combining the data from many galaxies into composite spectra. This has the advantage of rejecting (or at least smoothing over) sky-subtraction residuals and other artefacts, but obviously limits the degree to which we can test IMF variations within the sample. Secondly, our analysis method is different from that of CvD12b: whereas they perform a full-spectrum fit (with a total of 22 free parameters), over their optical and far-red spectra simultaneously, we instead adopt most parameters from the best-fitting SSP model in the optical, and focus on the infrared indices to isolate the possible IMF variations.

A novel aspect of our work is that we introduce and exploit a new IMF-sensitive index measuring the Ca I line at 10345 Å. In contrast to *all* of the ‘classical’ IMF-sensitive features (WFB, Na I doublet and Ca II triplet), this index is nearly independent of the sodium abundance, and hence provides an independent test which complements results obtained from other features. We quantify this statement using the CvD12a models. Using the 13.5-Gyr, solar-abundance, Chabrier-IMF case as a fiducial point, we find that

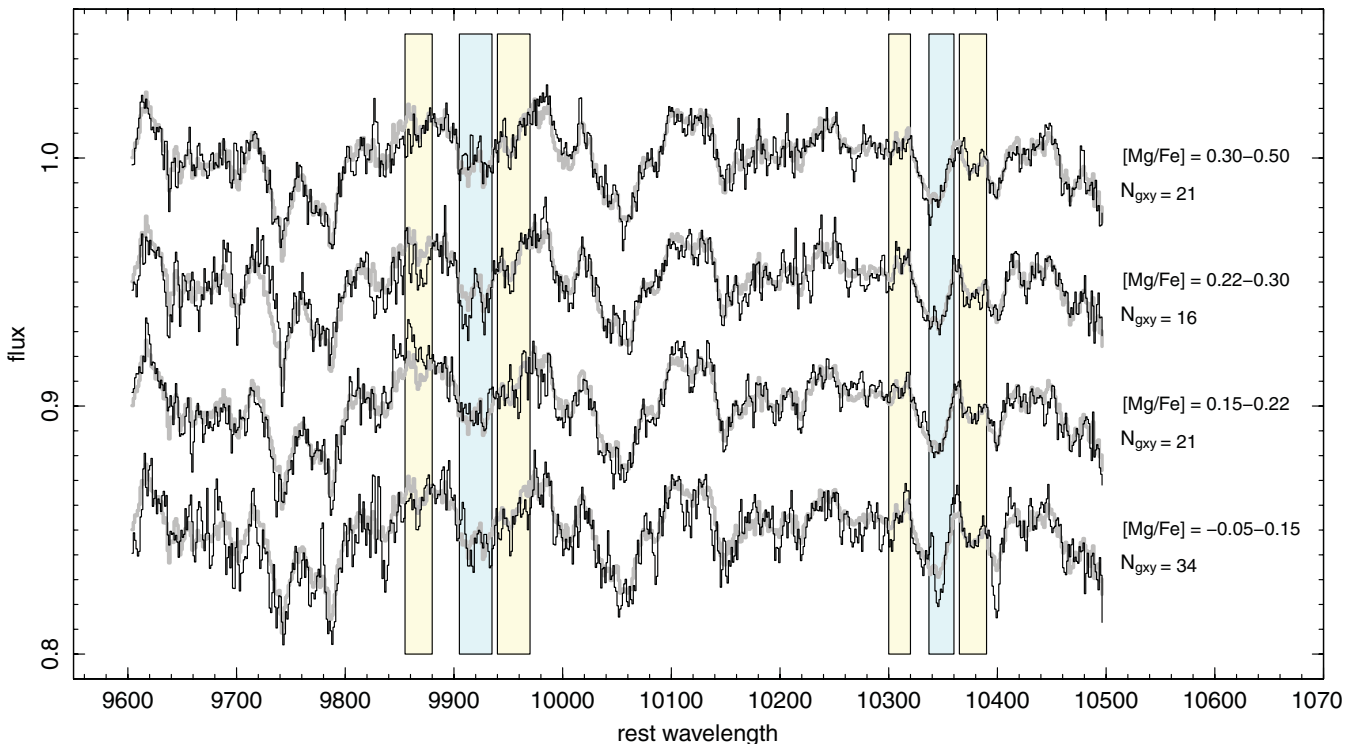


Figure 11. Composite spectra computed in four bins of Mg/Fe (black lines), for galaxies with $\sigma > 50 \text{ km s}^{-1}$. In each case, the thick grey line shows the fiducial composite of $\sigma > 100 \text{ km s}^{-1}$ galaxies for comparison. No smoothing has been applied to either the Mg/Fe split or the global composite.

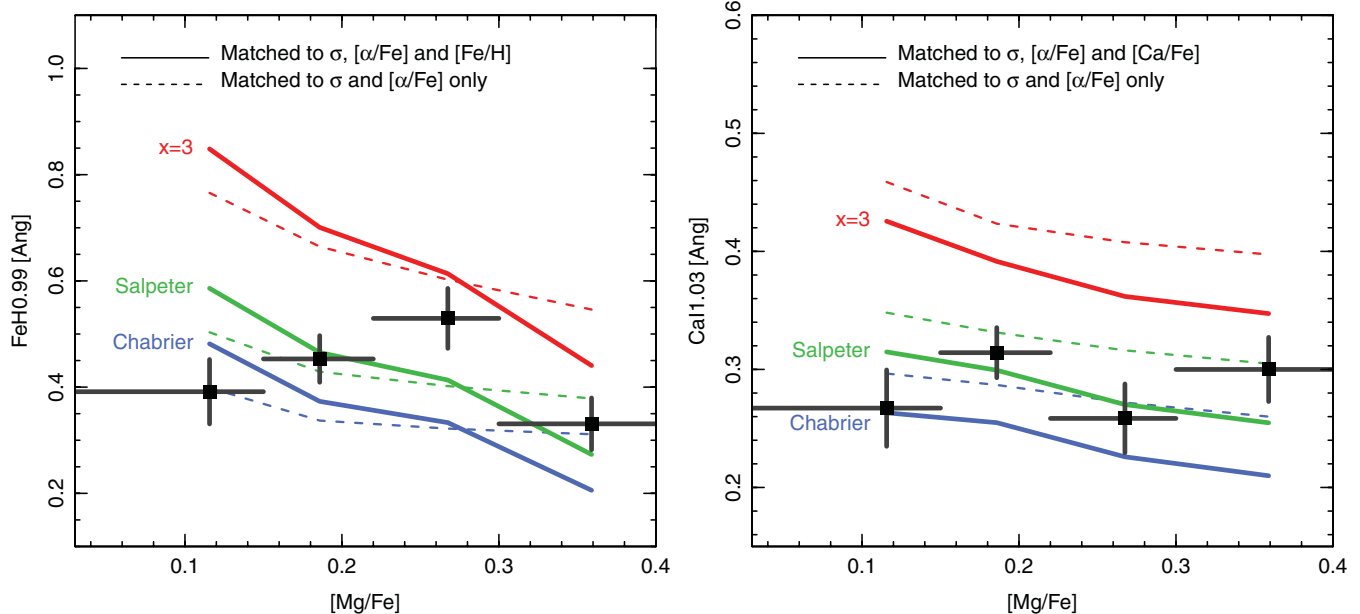


Figure 12. Variation in the FeH0.99 and CaI1.03 indices as a function of $[Mg/Fe]$. Bins in $[Mg/Fe]$ are as in Fig. 11. The models shown are as in Fig. 10.

increasing Na/Fe by 0.3 dex increases the NaI0.82 index by 0.14 \AA . If instead Na/Fe is held fixed, but the IMF is changed from Chabrier to Salpeter, this index increases by 0.12 \AA . The ratio of sodium sensitivity to IMF sensitivity for NaI0.82 is thus $0.14/0.12 = +1.17$. For the CaII0.86 index (measuring the Ca II triplet), the equivalent ratio is $+0.44$, while for FeH0.99 (measuring the WFB), the ratio is -0.56 . (As emphasized by CvD12a, sodium indirectly affects the strength of many spectral features, since it is one of the dominant electron donors in cool stellar atmospheres, and hence changes the ionization balance of other species.) In contrast, increasing Na/Fe by 0.3 dex reduces the CaI1.03 index by 0.002 \AA , while changing to a Salpeter IMF increases this index by 0.036 \AA and hence the sensitivity ratio is -0.06 , an order of magnitude smaller than the classical indices.

Our first-order result is that the WFB and Ca I measurements favour a Salpeter-like IMF on average in red-sequence/early-type galaxies, rather than the bottom-light Chabrier/Kroupa form found in the Milky Way. This is an important result, since it implies an ~ 50 per cent correction to the K -band mass-to-light ratios for early types relative to spirals. This conclusion is consistent with the average result from CvD12b in the relevant mass range, and also agrees with recent lensing and dynamical analyses (Spiniello et al. 2011; Cappellari et al. 2012).

We do not recover a clear signal of *increasing* dwarf enrichment with increasing mass (traced by velocity dispersion). The evidence for such a trend in CvD12b is driven mainly by the following: (a) their original stacked spectra of four Virgo galaxies with $\sigma \approx 300 \text{ km s}^{-1}$ from van Dokkum & Conroy (2010), which require an IMF significantly heavier than Salpeter; and (b) their ~ 10 galaxies at $\sigma \lesssim 180 \text{ km s}^{-1}$ which favour a Chabrier-like IMF on average. We have a few galaxies in the 300 km s^{-1} regime, and our highest σ bin is dominated by galaxies with $\sigma \approx 240 \text{ km s}^{-1}$. Hence, our results do not exclude a steeper IMF at the high-mass extreme. At lower σ , however, where our results are based on a much larger galaxy sample, our data clearly favour a Salpeter IMF, in apparent conflict with the Chabrier-like solutions recovered by CvD12b.

Our analysis supports the trend found by CvD12b for an increase in the dwarf content of red-sequence galaxies as a function of their α -abundance ratio. Notably, the only galaxies for which we infer a Chabrier-like IMF, similar to that of the Milky Way, are those which have abundance patterns closest to that of the Milky Way ($[Mg/Fe] \approx 0.1$). For galaxies with high α -enhancement ratios ($[Mg/Fe] \gtrsim 0.3$), thought to arise from formation in very rapid and intense starbursts, the indices require IMFs that are heavier than Salpeter, though not so extreme as an $x = 3$ power law.⁷

We emphasize that the results obtained from the two measured indices are consistent with each other. CvD12b note that their results depend somewhat on which combinations of IMF-sensitive features they include in their fits (although their overall conclusions are robust). For example, if the Na I region is removed from the fit, they recover lower dwarf enrichments, whereas if both Na I and the Ca II triplet are removed, so that only the WFB is constraining the IMF, then even heavier solutions are recovered. Their best-fitting models require large enhancements in Na/Fe, of up to an order of magnitude higher than the solar ratio in some cases. Since all three classical IMF indicators are affected (directly or indirectly) by the sodium abundance, it is important for future work to determine whether this high sodium enrichment is real, and not an artefact of inconsistencies between the various fitted features.⁸ In the meantime, it is encouraging that our sodium-independent CaI1.03 index yields results which are concordant with those from the WFB.

⁷ Our results refer strictly to the IMF at low stellar mass ($\lesssim 0.5 M_{\odot}$). Although a top-heavy IMF in bursts has been proposed as an explanation for the submillimetre galaxy counts (Baugh et al. 2005), this needs only to affect the IMF at much higher masses, $5\text{--}20 M_{\odot}$ (see Lacey et al. 2010). Hence, it is possible to construct an IMF satisfying both constraints, although we reserve judgement as to whether it is physically plausible.

⁸ Enhancements of $0.5\text{--}1.0$ dex in sodium have been reported for high-metallicity Milky Way bulge giants (Cunha & Smith 2006; Lecureur et al. 2007), but only mild (~ 0.1 dex) enhancements were found by Bensby et al. (2010) for bulge dwarfs and subgiants.

We finish by briefly comparing our results to two recent spectroscopic works based on the Na I doublet. Spiniello et al. (2012) analysed galaxies with velocity dispersions generally larger than average for our sample, but overlapping our two highest σ bins. They find a trend of increasing IMF slope from Salpeter-like at $\sigma \approx 200 \text{ km s}^{-1}$ (in agreement with our results at the same velocity dispersion) to $x \approx 3$ at $\sigma \approx 335 \text{ km s}^{-1}$ (a regime not probed by our sample). Ferreras et al. (2012) report consistency with a Salpeter-like IMF at $\sigma \approx 200 \text{ km s}^{-1}$ and $x \approx 3$ at 300 km s^{-1} . At lower mass, they infer a constant Kroupa-like IMF at $\sigma \lesssim 150 \text{ km s}^{-1}$, which may be marginally discrepant with the Salpeter IMF favoured by our results in this mass range. This comparison should be viewed with caution however, since Ferreras et al. have not yet taken into account the possible effect of Na/Fe variations as a function of σ . We note that our sample differs from both Spiniello et al. and Ferreras et al. in studying only rich galaxy cluster members. It is not yet known whether apparent IMF variations are correlated in any way with galaxy environment.

6 CONCLUSIONS

We have presented new infrared spectroscopy for red-sequence galaxies in Coma, and measured gravity-sensitive absorption features to probe their low-mass stellar content. Compared to other recent work in this field, our study differs in targeting cluster galaxies and in sampling ‘typical’ red-sequence galaxies, rather than only the most massive objects. By comparing line-strength indices against state-of-the-art spectral synthesis models, tuned to match element abundances estimated from optical data, we derive constraints on the average IMF and its variation with velocity dispersion and α -abundance ratio.

Our main conclusions are as follows.

(i) The average rest-frame 1- μm spectra for sizeable samples of galaxies at the distance of Coma can be recovered with Subaru/FMOS, thanks to its wide-field multiplex capability which is currently unique in the infrared.

(ii) The observed spectral features are generally in good qualitative agreement with the latest stellar libraries and synthesis models, although some localized discrepancies are seen.

(iii) The Ca I line at 10345 \AA is promising as a new gravity-sensitive feature, which unlike the ‘classical’ IMF indicators is independent of the sodium abundance.

(iv) The strength of the WFB and Ca I 10345 \AA line indicate that red-sequence galaxies in Coma have (on average) a dwarf-star content similar to that in a Salpeter IMF.

(v) There is no clear trend in the derived IMF as a function of velocity dispersion, with Salpeter models an adequate fit from $\sigma \approx 100$ to $\approx 250 \text{ km s}^{-1}$.

(vi) A more dwarf-dominated IMF cannot be ruled out at the highest velocity dispersions ($\sigma \gtrsim 300 \text{ km s}^{-1}$), which are not well sampled by our observations.

(vii) The derived IMF is correlated with Mg/Fe ratio, suggesting that galaxies which underwent intense rapid starbursts formed a larger number of low-mass stars per solar mass star.

Our work adds to a developing consensus that massive red-sequence galaxies have heavier IMFs, with a greater contribution from low-mass stars, than found in the Milky Way. A Salpeter-like IMF is favoured by recent spectroscopic (CvD12b; Ferreras et al. 2012; Spiniello et al. 2012; this paper), dynamical (Thomas et al. 2011; Cappellari et al. 2012) and lensing analyses (Spiniello et al.

2011; Sonnenfeld et al. 2012). For spiral galaxies, in contrast, recent lensing results are consistent with Chabrier-like IMFs similar to the Milky Way (Brewer et al. 2012), consistent with earlier dynamical results (Bell & de Jong 2001). Together with the apparent correlation of dwarf-star content with Mg/Fe ratio found here and in CvD12b, these results strongly suggest that the low-mass IMF may be dictated in part by the dominant mode of star formation: quiescent star formation yields Kroupa-like IMF, while starbursts lead to an excess of low-mass stars.

ACKNOWLEDGMENTS

We are grateful to Kentaro Aoki for expert assistance in our Subaru observations and to Charlie Conroy for providing model spectra in advance of publication. RJS was supported for this work by STFC Rolling Grant PP/C501568/1 ‘Extragalactic Astronomy and Cosmology at Durham 2008–2013’.

REFERENCES

- Bastian N., Covey K. R., Meyer M. R., 2010, *ARA&A*, 48, 339
 Baugh C. M., Lacey C. G., Frenk C. S., Granato G. L., Silva L., Bressan A., Benson A. J., Cole S. M., 2005, *MNRAS*, 356, 1191
 Bell E. F., de Jong R. S., 2001, *ApJ*, 550, 212
 Bensby T. et al., 2010, *A&A*, 512, A41
 Brewer B. J. et al., 2012, *MNRAS*, 422, 3574
 Cappellari M. et al., 2012, *Nat*, 484, 485
 Carter D., Visvanathan N., Pickles A. J., 1986, *ApJ*, 311, 637
 Cenarro A. J., Gorgas J., Vazdekis A., Cardiel N., Peletier R. F., 2003, *MNRAS*, 339, L12
 Chabrier G., 2003, *PASP*, 115, 763
 Cohen J. G., 1978, *ApJ*, 221, 788
 Conroy C., van Dokkum P., 2012a, *ApJ*, 747, 69 (CvD12a)
 Conroy C., van Dokkum P., 2012b, *ApJ*, preprint (arXiv:1205.6473) (CvD12b)
 Couture J., Hardy E., 1993, *ApJ*, 406, 142
 Cunha K., Smith V. V., 2006, *ApJ*, 651, 491
 Faber S. M., French H. B., 1980, *ApJ*, 235, 405
 Ferreras I., La Barbera F., de Carvalho R. R., de la Rosa I. G., Vazdekis A., Falcón-Barroso, Ricciardelli E., 2012, *ApJ*, preprint (arXiv:1206.1594)
 Graves G. J. M., Faber S. M., Schiavon R. P., Yan R., 2007, *ApJ*, 671, 243
 Hinkle K., Wallace L., Livingston W. C., 1995, *Infrared Atlas of the Arcturus Spectrum: 0.9–5.3 μm* . Astron. Soc. Pac., San Francisco
 Iwamuro F. et al., 2012, *PASJ*, 64, 59
 Johansson J., Thomas D., Maraston C., 2012, *MNRAS*, 421, 1908
 Kimura M. et al., 2010, *PASJ*, 62, 1135
 Kroupa P., Tout C., Gilmore G., 1993, *MNRAS*, 262, 545
 Lacey C. G., Baugh C. M., Frenk C. S., Benson A. J., Orsi A., Silva A., Granato G. L., Bressan A., 2010, *MNRAS*, 405, 2
 Lecureur A., Hill V., Zoccali M., Barbuy B., Gómez A., Minniti D., Ortolani S., Renzini A., 2007, *A&A*, 465, 799
 Nelan J. E., Smith R. J., Hudson M. J., Wegner G. A., Lucey J. R., Moore S. A. W., Quinney S. J., Suntzeff N. B., 2005, *ApJ*, 632, 137
 Price J., Phillipps S., Huxor A., Smith R. J., Lucey J. R., 2011, *MNRAS*, 411, 2558
 Rayner J. T., Cushing M. C., Vacca W. D., 2009, *ApJS*, 185, 289
 Salpeter E., 1955, *ApJ*, 121, 161
 Schiavon R. P., 2007, *ApJS*, 171, 146
 Smith R. J., Lucey J. R., Hudson M. J., Bridges T. J., 2009, *MNRAS*, 398, 119
 Smith R. J., Lucey J. R., Price J., Hudson M. J., Phillipps S., 2012, *MNRAS*, 419, 3167
 Sonnenfeld A., Treu T., Gavazzi R., Marshall P. J., Auger M. W., Suyu S. H., Koopmans L. V. E., Bolton A. S., 2012, *ApJ*, 752, 163
 Spiniello C., Koopmans L. V. E., Trager S. C., Czoske O., Treu T., 2011, *MNRAS*, 417, 3000

Spiniello C., Trager S. C., Koopmans L. V. E., Chen Y. P., 2012, ApJ, 753, L32
Spinrad H., Taylor B. J., 1971, ApJS, 22, 445
Thomas D., Maraston C., Bender R., Mendes de Oliveira C., 2005, ApJ, 621, 673
Thomas D., Maraston C., Bender R., 2003, MNRAS, 343, 279
Thomas J. et al., 2011, MNRAS, 415, 545
Treu T., Auger M. W., Koopmans L. V. E., Gavazzi R., Marshall P. J., Bolton A. S., 2010, ApJ, 709, 1195

van Dokkum P. G., Conroy C., 2010, Nat, 468, 940
van Dokkum P., Conroy C., 2012, ApJ, preprint (arXiv:1205.6471)
Vazdekis A., Peletier R. F., Beckman J. E., Casuso E., 1997, ApJS, 111, 203
Whitford A. E., 1977, ApJ, 211, 527
Wing R. F., Ford W. K., 1969, PASP, 81, 527

This paper has been typeset from a $\text{\TeX}/\text{\LaTeX}$ file prepared by the author.