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# Saturated VSD model of a six-phase induction machine

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**Abstract:** In this paper, a saturated model of an asymmetrical six-phase induction machine is presented. The model is based on the vector space decomposition approach, and it includes main and leakage flux saturation, as well as the mutual coupling between orthogonal planes, while using the Gamma equivalent circuit approach. The accuracy of the proposed model in unbalanced operating modes that are characteristic for multiphase machines, such as post-fault and power sharing operation, makes it advantageous compared to existing models. The model is developed assuming that the machine operates in asymmetrical conditions and that, therefore, there is fundamental frequency excitation in both planes. The inter-plane coupling effect is examined using finite element analysis and an experimental procedure for its quantification is developed. The model is verified by comparison with the experimental results obtained from a prototype asymmetrical six-phase induction machine, and its advantages compared to existing models are emphasized.

## 1 Introduction

Multiphase electrical machines are an opportune choice for a wide range of applications in modern power systems, such as high-power industrial drives [1], electric vehicles [2], more-electric aircraft [3, 4], and wind energy conversion [5]. Multiphase machines exhibit multiple advantages compared to their traditionally employed three-phase counterparts, the most noteworthy being increased torque density and torque/current ratio, reduced torque ripple, improved fault tolerance and lower current harmonic distortion [1, 6, 7].

Accurate modelling and parameter determination are essential for proper control and performance analysis of multiphase machines [8]. A straightforward approach would be to model the multiphase machine using phase-domain variables. However, the phase-variable model is highly computationally demanding due to time-varying inductance matrix coefficients [9]. This led to the development of alternative modelling approaches: the multiple  $dq$  [10], vector space decomposition (VSD) [11] and voltage behind reactance (VBR) [12] models. Each approach is more or less beneficial depending on the intended application.

Regardless of the model employed, proper modelling of magnetic circuit saturation is required for achieving satisfactory simulation accuracy and control performance. Electrical machines usually operate in the saturated region to increase torque density [13]. Extensive research was devoted to saturated models of three-phase machines over the past decades and is well documented [14–17]. Saturated multiphase machine models are still under development, with some notable research reported in [12, 18–21]. Most of the available models rely on including only main flux saturation, which relates directly to existing saturated three-phase machine models. Specific aspects of saturation in five-phase machines regarding the air-gap flux waveform have been studied in depth in [22] and [12].

In recent years, the VSD model has been a preferred choice for the control of multiphase machines [23–25]. However, an in-depth study of the saturation influence on the VSD model remains to be conducted. A coupling effect between the orthogonal subspaces (planes) of an asymmetrical six-phase induction machine (6PIM) was observed and described in [26]. This effect was termed “interplane cross-saturation”, or IPCS for short. This phenomenon represents a mutual coupling between fundamental frequency quantities of the  $dq$  and  $xy$  planes. In normal operating conditions there are no fundamental currents in the  $xy$  plane. Fundamental frequency

components in the  $xy$  plane are related to unbalanced operation, and therefore need to be accounted for in fault-tolerant control [25, 27] and arbitrary power-sharing algorithms for multiple three-phase winding machines [28, 29]. Saturated multiphase machine models which include the mutual influence between the  $dq$  and  $xy$  plane have been reported in literature [12], but only the influence of higher-order  $xy$  current harmonics was studied, which is not applicable for analysing unbalanced operation. This imposes the need to develop a new VSD model including the IPCS effect.

In this paper, the VSD model of a 6PIM with a distributed winding (shown in the Appendix, subsection 8.1), including main and leakage flux saturation, as well as the IPCS effect is presented. A parameter measurement procedure for determining the coupling flux components due to IPCS is developed. The proposed model provides improved accuracy compared to existing models when analysing unbalanced operating modes with pronounced saturation. The model validity and parameter measurement accuracy are verified experimentally. Verification is performed for operating modes with pronounced saturation and voltage unbalance.

This paper is organized as follows. The conventional linear VSD model is presented in section 2. The VSD model is formulated using a Gamma equivalent circuit in section 3. The development of the model including main flux saturation and IPCS is presented in the fourth section. The parameter measurement procedure and a comparison between simulation results and experimental results are given in section 5. The final section summarises the conclusions of the presented research.

## 2 Conventional VSD model

The VSD model voltage equations of a 6PIM in the  $dq$  and  $xy$  plane are given in space-vector form as follows [9]:

$$\vec{u}_{dqs} = R_s \vec{i}_{dqs} + \frac{d\vec{\psi}_{dqs}}{dt} + j\omega_e \vec{\psi}_{dqs} \quad (1a)$$

$$\vec{u}_{dqr} = R_r \vec{i}_{dqr} + \frac{d\vec{\psi}_{dqr}}{dt} + j(\omega_e - \omega_r) \vec{\psi}_{dqr} \quad (1b)$$

$$\vec{u}_{xys} = R_s \vec{i}_{xys} + \frac{d\vec{\psi}_{xys}}{dt}, \quad (1c)$$

where  $\omega_r$  (rad/s) is the rotor electrical angular speed, and  $\omega_e$  (rad/s) is the arbitrary angular speed of the rotating reference frame. Note that the zero-sequence equations are excluded from the given model, due to the fact that the neutral points of the 6PIM are considered isolated in the upcoming analyses. The machine could of course be configured with a single neutral point, in which case injection of the third stator harmonic current would be possible. This would undoubtedly impact significantly on the saturation-related behaviour of the machine. However, configuration with single neutral and the third harmonic stator current injection for torque enhancement purposes is never used in practice since it requires the connection of the machine's neutral point to either a seventh inverter leg or a mid-point of the dc-link capacitor bank [30, 31]. Hence this situation is not considered in the paper.

Additionally, since a machine with a squirrel-cage rotor is used, the rotor voltages are set to zero further on. The space vectors in the  $dq$  and  $xy$  subspaces are defined as:

$$\vec{\xi}_{dqs} = \xi_{ds} + j\xi_{qs} \quad (2a)$$

$$\vec{\xi}_{dqr} = \xi_{dr} + j\xi_{qr} \quad (2b)$$

$$\vec{\xi}_{xys} = \xi_{xs} + j\xi_{ys}, \quad (2c)$$

where  $\xi$  stands for an arbitrary 6PIM electrical quantity (voltage, current, or flux linkage). The flux linkages are given as:

$$\vec{\psi}_{dqs} = (L_m + L_{lsdq})\vec{i}_{dqs} + L_m\vec{i}_{dqr} \quad (3a)$$

$$\vec{\psi}_{dqr} = (L_m + L_{lr})\vec{i}_{dqr} + L_m\vec{i}_{dqs} \quad (3b)$$

$$\vec{\psi}_{xys} = L_{xy}\vec{i}_{xys}, \quad (3c)$$

where  $L_m$  is the magnetizing inductance,  $L_{lsdq}$  is the stator leakage inductance in the  $dq$  plane, and  $L_{xy}$  is the stator leakage inductance in the  $xy$  plane. The leakage inductances in the  $dq$  and  $xy$  planes are generally not equal due to mutual leakage inductance, as elaborated in [32]. Note that there is no mutual influence between the quantities of different subspaces, and that, with no saturation involved, all inductances in (3) are constant. The electromagnetic torque equation is given as:

$$T_e = \frac{1}{3K^2} p \operatorname{Im} \left( \vec{\psi}_{dqs} \cdot \vec{i}_{dqr}^* \right) \quad (4)$$

where  $*$  denotes complex conjugation,  $p$  is the pole pair number, and  $K$  is the VSD transformation coefficient. The value of  $K$  can be chosen arbitrarily, and is set to  $K = 1/3$  in all the analyses in this paper, thereby keeping the peak values of electrical quantities from the original phase-variable reference frame invariant. To complete the model, the electromechanical motion equation needs to be included:

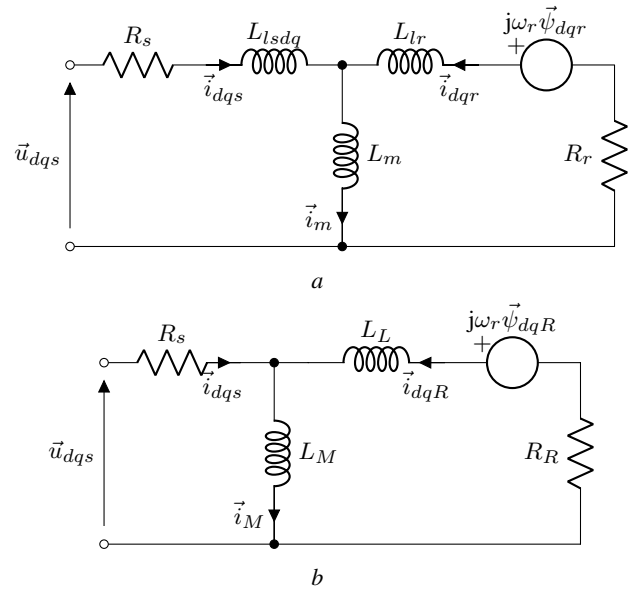
$$T_e - T_L = J \frac{d\Omega}{dt} + k_f \Omega, \quad (5)$$

where  $T_L$  is the load torque,  $J$  is the moment of inertia,  $\Omega = \omega/p$  is the mechanical angular rotor speed, and  $k_f$  is the friction coefficient.

### 3 Gamma model

As Slemon demonstrated in [33], the Gamma ( $\Gamma$ ) model is a suitable alternative to the conventional T-form model of a three-phase induction machine. In fact, the  $\Gamma$  model was shown to be advantageous in terms of including main flux saturation in motoring mode, when saturation occurs dominantly in the stator magnetic circuit. An obvious advantage of the  $\Gamma$  model is that the equivalent circuit parameters may be obtained directly from the conventional no-load and locked-rotor tests.

The 6PIM equations in the  $dq$  plane are identical to those of the  $dq$  model of a three-phase induction machine. This allows the  $dq$  subspace to be modelled using the  $\Gamma$  circuit. The T and  $\Gamma$  equivalent circuits corresponding to the stationary reference frame ( $\omega_e = 0$ ) are



**Fig. 1:** Equivalent circuits of a 6PIM in the  $dq$  plane in the stationary reference frame

a T circuit  
b  $\Gamma$  circuit

displayed in Fig. 1. The rotor quantities of the  $\Gamma$  model are related to those of the T model as follows:

$$\vec{\psi}_{dqR} = \gamma \vec{\psi}_{dqr} \quad (6a)$$

$$\vec{i}_{dqR} = \frac{\vec{i}_{dqr}}{\gamma}, \quad (6b)$$

where  $\gamma = L_s/L_m$  and  $L_s = L_m + L_{lsdq}$  is the stator no-load inductance. The relationships between  $\Gamma$  and T circuit parameters are given as:

$$L_M = \gamma L_m = L_s \quad (7a)$$

$$L_L = \gamma L_{lsdq} + \gamma^2 L_{lr} \quad (7b)$$

$$R_R = \gamma^2 R_r. \quad (7c)$$

All parameters of the  $\Gamma$  circuit can be obtained directly from the standard no-load and locked-rotor tests under  $dq$  voltage supply. The stator resistance can be determined either from the DC test or the  $xy$  sequence test [34]. The inductance  $L_M$ , which will be referred to as the “magnetizing inductance” further on, is equal to the no-load inductance, and is therefore obtained from the no-load test. After this, the locked-rotor test enables the determination of the two remaining parameters – the leakage inductance  $L_L$  and the rotor resistance  $R_R$ . Note that the coefficient  $\gamma$  can only be determined provided that the magnetizing inductance is known. However, this information is only required if the values of actual rotor quantities are of interest, which is usually not the case.

The  $\Gamma$  model is used in all of the following analyses. The  $dq$  plane equations of the conventional VSD model given in the previous section should be appropriately modified. The model in the  $xy$  subspace remains unaltered, as the  $\Gamma$  model relates only to the  $dq$  quantities. The voltage equations (1) and the torque equation (4) retain the same form, except that the indices “r” should be replaced with “R” in all instances. The rotor voltage equation is now given as:

$$0 = R_R \vec{i}_{dqR} + \frac{d\vec{\psi}_{dqR}}{dt} + j(\omega_e - \omega_r) \vec{\psi}_{dqR}, \quad (8)$$

whereas the torque equation is given as:

$$T_e = \frac{1}{3K^2} p \operatorname{Im} \left( \vec{\psi}_{dqs} \cdot \vec{i}_{dqR}^* \right) \quad (9)$$

The flux equations of the linear  $\Gamma$  model are given as follows:

$$\vec{\psi}_{dqs} = L_M (\vec{i}_{dqs} + \vec{i}_{dqr}) \quad (10a)$$

$$\vec{\psi}_{dqR} = L_M (\vec{i}_{dqs} + \vec{i}_{dqr}) + L_L \vec{i}_{dqR}. \quad (10b)$$

## 4 Proposed saturated VSD model

According to the results of [26], the  $dq$  and  $xy$  planes are coupled under saturated conditions. This phenomenon has been termed interplane cross-saturation. In the following subsections, a detailed analysis of IPCS and its effect on model variables is carried out using Finite element analysis (FEA). After that, a new model is proposed, incorporating main flux saturation and IPCS effect. Therein, the linear  $\Gamma$  model presented in section 3 is used as a foundation.

### 4.1 Finite element analysis of IPCS

The goal of this analysis is to determine the influence of IPCS on the flux linkages in the  $dq$  and  $xy$  planes. This is achieved by comparing the flux linkages obtained when only the current components in the corresponding plane exist to those obtained when current components in both planes are present. In other words, the difference between the flux linkage values corresponding to cases with and without mutual interaction between the orthogonal planes needs to be determined. These differences are expressed as:

$$\Delta\vec{\psi}_{dq}(\hat{i}_M, \hat{i}_{xys}) = \vec{\psi}_{dqs}(\hat{i}_M, \hat{i}_{xys}) - \vec{\psi}_{dqs}(\hat{i}_M, 0) \quad (11a)$$

$$\Delta\vec{\psi}_{xy}(\hat{i}_M, \hat{i}_{xys}) = \vec{\psi}_{xys}(\hat{i}_M, \hat{i}_{xys}) - \vec{\psi}_{xys}(0, \hat{i}_{xys}), \quad (11b)$$

where  $\Delta\vec{\psi}_{dq}$  and  $\Delta\vec{\psi}_{xy}$  are the flux changes caused by IPCS, and  $\hat{i}_M$  and  $\hat{i}_{xys}$  represent the magnitudes of the magnetizing and  $xy$  plane space vectors, respectively. It should be noted that the FEA was conducted under conditions corresponding to no-load operation, hence the  $\hat{i}_M$  is equivalent to the magnitude of the  $dq$  plane current space vector, i.e.  $\hat{i}_M \equiv \hat{i}_{dqs}$ .

The FE simulations were executed for multiple combinations of  $dq$  and  $xy$  currents, and the flux linkage changes attributed to IPCS were determined according to (11) for each combination. The machine model used in the FEA is based on the actual 6PIM prototype used in the upcoming experimental investigation. The construction details (machine dimensions and winding data) and rated data of the prototype are given in subsection 8.1 of the Appendix. Details regarding the FE model and an example of an FEA solution with the displayed finite element mesh and flux density distribution are also given in the Appendix, in subsection 8.2. In order to determine the dependence of the flux increments on the  $dq$  and  $xy$  current components, the FEA is performed for the following scenarios:

1. phase currents exciting only the  $dq$  subsystem are assigned to the windings (flux linkages  $\vec{\psi}_{dqs}(\hat{i}_M, 0)$  and  $\vec{\psi}_{xys}(\hat{i}_M, 0)$  are obtained);
2. phase currents exciting only the  $xy$  subsystem are assigned to the windings (flux linkages  $\vec{\psi}_{dqs}(0, \hat{i}_{xys})$  and  $\vec{\psi}_{xys}(0, \hat{i}_{xys})$  are obtained);
3. all combinations of phase currents from scenarios 1 and 2 are assigned to the windings, thereby exciting both the  $dq$  and  $xy$  subsystems (flux linkages  $\vec{\psi}_{dqs}(\hat{i}_M, \hat{i}_{xys})$  and  $\vec{\psi}_{xys}(\hat{i}_M, \hat{i}_{xys})$  are obtained).

Each step of the analysis is conducted as follows:

- the values of currents in the  $dq$  and  $xy$  subspace ( $\hat{i}_M$  and  $\hat{i}_{xys}$ ) are selected;
- the phase current values to be assigned to the windings are calculated by applying the inverse VSD transformation [9] to the  $dq$  and  $xy$  current components;
- the FE simulation is conducted for the assigned current values and the values of per-phase flux linkages are obtained;

- the complex VSD transformation is applied to the per-phase flux linkage values to obtain the space vectors of flux linkages in the  $dq$  and  $xy$  subspace  $\vec{\psi}_{dqs}$  and  $\vec{\psi}_{xys}$ .

The described procedure is repeated for all three scenarios, with  $\hat{i}_M$  and  $\hat{i}_{xys}$  ranging from 0 to 6 A. The flux linkage changes are then determined by applying (11a) and (11b). The obtained flux changes were found to be shifted by approximately  $180^\circ$  with respect to the total flux linkages, and can therefore be expressed as follows (the negative sign is incorporated in the amplitudes of the flux increments):

$$\Delta\vec{\psi}_{dq}(\hat{i}_M, \hat{i}_{xys}) = \Delta\psi_{dq}(\hat{i}_M, \hat{i}_{xys})e^{j\angle\vec{\psi}_{dqs}} \quad (12a)$$

$$\Delta\vec{\psi}_{xy}(\hat{i}_M, \hat{i}_{xys}) = \Delta\psi_{xy}(\hat{i}_M, \hat{i}_{xys})e^{j\angle\vec{\psi}_{xys}}, \quad (12b)$$

where  $\Delta\psi_{dq}$  and  $\Delta\psi_{xy}$  are algebraic values of the flux changes. The computed values are displayed in Fig. 2. The presented results clearly demonstrate that the flux changes are negative, i.e. that the IPCS effect causes a reduction of flux linkage values in both planes. Considering the flux values in the corresponding planes (rated  $dq$  plane flux  $\sim 0.7$  Wb;  $xy$  plane flux at rated current  $\sim 35$  mWb), the IPCS effect is much more pronounced in the  $xy$  plane than in the  $dq$  plane. This can be attributed to the main flux saturation, which has a much greater influence on the  $dq$  plane flux linkage compared to IPCS, particularly at large values of magnetizing current, as will be demonstrated next.

The main flux saturation may be accounted for by expressing the magnetizing flux as a non-linear function of the magnetizing current:

$$\vec{\psi}_M(\hat{i}_M) = \vec{\psi}_{Mu} + \Delta\psi_M(\hat{i}_M)e^{j\angle\vec{\psi}_M}, \quad (13)$$

where  $\vec{\psi}_{Mu} = L_{Mu}\vec{i}_M$  is the unsaturated value of magnetizing flux linkage,  $L_{Mu}$  is the unsaturated value of the magnetizing inductance and  $\Delta\psi_M$  is the algebraic flux change due to main flux saturation. According to the  $\Gamma$  model relationships, it follows that  $\vec{\psi}_M \equiv \vec{\psi}_{dqs}$  and  $\Delta\vec{\psi}_M \equiv \Delta\vec{\psi}_{dqs}$ , i.e. the value  $\Delta\psi_M$  actually represents the change in the  $dq$  plane flux linkage. This allows the total difference between the values of actual and unsaturated  $dq$  plane flux linkages to be expressed as:

$$\begin{aligned} \Delta\vec{\psi}'_{dq}(\hat{i}_M, \hat{i}_{xys}) &= \Delta\vec{\psi}_M(\hat{i}_M) + \Delta\vec{\psi}_{dq}(\hat{i}_M, \hat{i}_{xys}) \\ &= \left[ \Delta\psi_M(\hat{i}_M) + \Delta\psi_{dq}(\hat{i}_M, \hat{i}_{xys}) \right] e^{j\angle\vec{\psi}_{dqs}}. \end{aligned} \quad (14)$$

The values of  $\Delta\psi'_{dq}$  obtained from FE simulations are displayed in Fig. 3. By comparing Figs. 2a and 3, it can be observed that the contribution of main flux saturation exceeds that of IPCS by up to two orders of magnitude. This indicates that the influence of IPCS on the  $dq$  plane flux linkage can be neglected from a practical standpoint. This possibility will be re-evaluated after obtaining the experimental results.

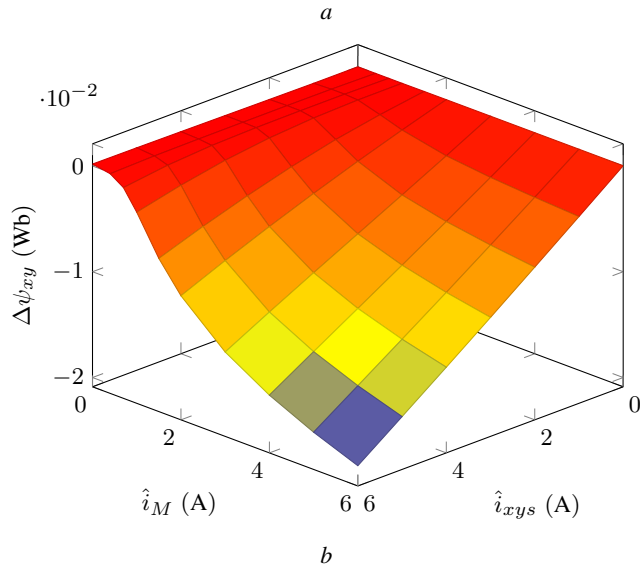
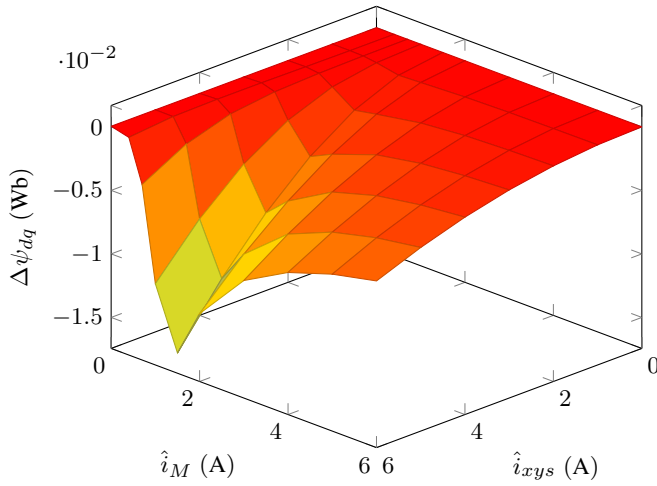
### 4.2 Model formulation

Based on the results obtained in the previous subsection, a saturated model of a 6PIM is proposed. The new model is based on the  $\Gamma$  model presented in section 3. Flux linkages in the  $dq$  and  $xy$  planes are used as the state variables. The flux equations in the  $dq$  and  $xy$  subspaces need to be modified in order to model the IPCS effect and main flux saturation. The following form of the flux equations is proposed:

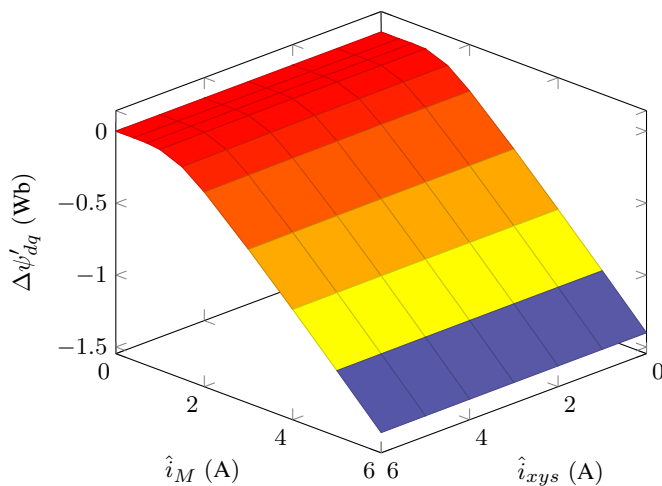
$$\vec{\psi}_{dqs} = L_{Mu}\vec{i}_M + \Delta\vec{\psi}'_{dq}(\hat{i}_M, \hat{i}_{xys}) \quad (15a)$$

$$\vec{\psi}_{dqR} = L_{Mu}\vec{i}_M + \Delta\vec{\psi}'_{dq}(\hat{i}_M, \hat{i}_{xys}) + L_L\vec{i}_{dqR} \quad (15b)$$

$$\vec{\psi}_{xys} = L_{xy}\vec{i}_{xys} + \Delta\vec{\psi}_{xy}(\hat{i}_M, \hat{i}_{xy}). \quad (15c)$$



**Fig. 2:** Changes of flux linkages due to IPCS determined by FEA  
a flux change in the  $dq$  plane  
b flux change in the  $xy$  plane



**Fig. 3:** Total change of  $dq$  plane flux linkage due to IPCS and main flux saturation determined by FEA

The voltage, torque and electromechanical motion equations remain unaltered compared to the original linear model. The following equations constitute the saturated 6PIM VSD model:

- voltage equations (1a), (8), and (1c),
- flux-current relationships (15),
- algebraic torque equation (9), and
- electromechanical motion equation (5).

The listed equations form a set of non-linear differential-algebraic equations (DAE), which can be readily solved by appropriate numerical integration methods. The non-linear flux increments can be included in the model by means of either lookup tables or suitable analytical functions.

## 5 Experimental results and discussion

### 5.1 Experimental setup overview

The experimental assembly used for 6PIM testing is depicted in Fig. 4. The core component is the 6PIM prototype, constructed by rewinding a three-phase squirrel-cage induction machine. As stated in section 4.1, the rated data and construction parameters of the prototype are given in the Appendix. The 6PIM is supplied from two identical converter groups, each comprised of an input rectifier stage, DC link, and an output inverter stage. The 6PIM is operated with two isolated neutrals.

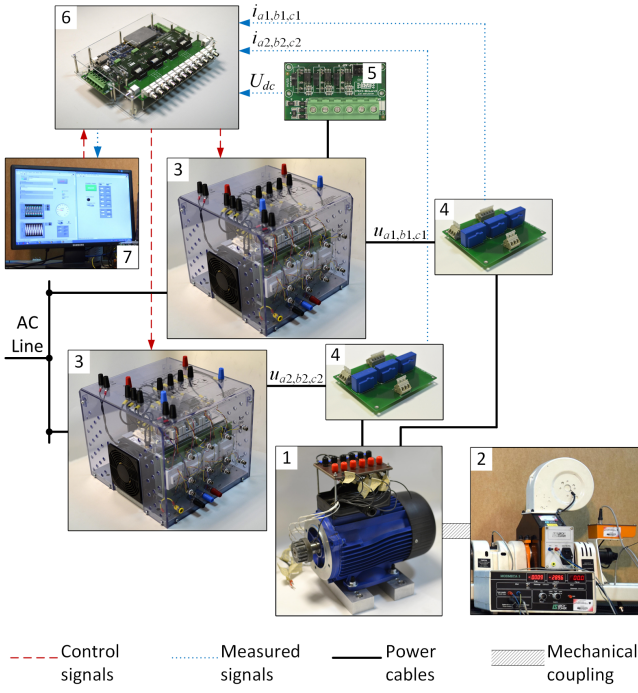
Motor currents are measured using six LEM current sensors, and the DC link voltage of the converters is measured using a high-voltage sensing board with input-output isolation and differential outputs. Converter control and analog-to-digital conversion (ADC) of measured signals are performed using the Power Electronics and Drives Board (PED-Board for short) controller based on an sbRIO-9651 platform by National Instruments. The sbRIO system incorporates an FPGA and processor, which allows simultaneous execution of multiple processes at various sampling rates. The PWM signal generation and high-speed A/D conversion are implemented on the FPGA, whereas the reference voltage values and other commands are executed via the processor. Sampled data are buffered and transferred to the PC for visualization and logging purposes, given the limited memory capacity of the FPGA and processor. The complete converter control and data acquisition algorithm was developed using the NI LabVIEW software package.

The converters are voltage controlled (open-loop control), and conventional (carrier-based) PWM is used. The switching period is set to  $120 \mu\text{s}$ , with the inverter dead-time of  $4.3 \mu\text{s}$ . The ADC sampling rate is set to 40 kHz. The measured signals are pre-filtered in order to avoid aliasing. Dead-time compensation is applied to the reference voltages, thereby eliminating the requirement for voltage measurement [35].

### 5.2 Parameter determination

The model parameters that need to be identified are: stator and rotor resistances ( $R_s$ ,  $R_R$ ), leakage inductances in the  $dq$  and  $xy$  plane ( $L_L$ ,  $L_{xy}$ ), unsaturated magnetizing inductance ( $L_{Mu}$ ), and flux changes due to main flux saturation and IPCS ( $\Delta\psi'_{dq}$ ,  $\Delta\psi_{xy}$ ). The stator resistance and leakage inductance in the  $xy$  plane are determined from the  $xy$  sequence test, as described in [32, 34]. The unsaturated magnetizing inductance,  $dq$  plane leakage inductance and rotor resistance are determined from the standard no-load and locked-rotor tests. The rotating speed of the 6PIM is maintained equal to synchronous speed by means of an auxiliary three-phase induction machine during the no-load and  $xy$  sequence tests.

The values of leakage inductances in the  $dq$  and  $xy$  subspaces are plotted against corresponding current magnitudes and displayed in Fig. 5. Both inductances are current-dependant, especially the one in the  $dq$  plane. This can be attributed to immediate saturation of rotor slot openings even at low current values, which is typical in the case of closed rotor slot [36]. This effect is very pronounced and needs to be accounted for. According to the guidelines for including leakage flux saturation given in [33],  $L_L$  is expressed as a non-linear function of  $\hat{i}_{dqs}$  (see the Appendix, subsection 8.3). The non-linearity of the  $xy$  plane inductance need not be explicitly modelled, but rather appended to the  $\Delta\psi_{xy}$  term, as demonstrated later in this section. It



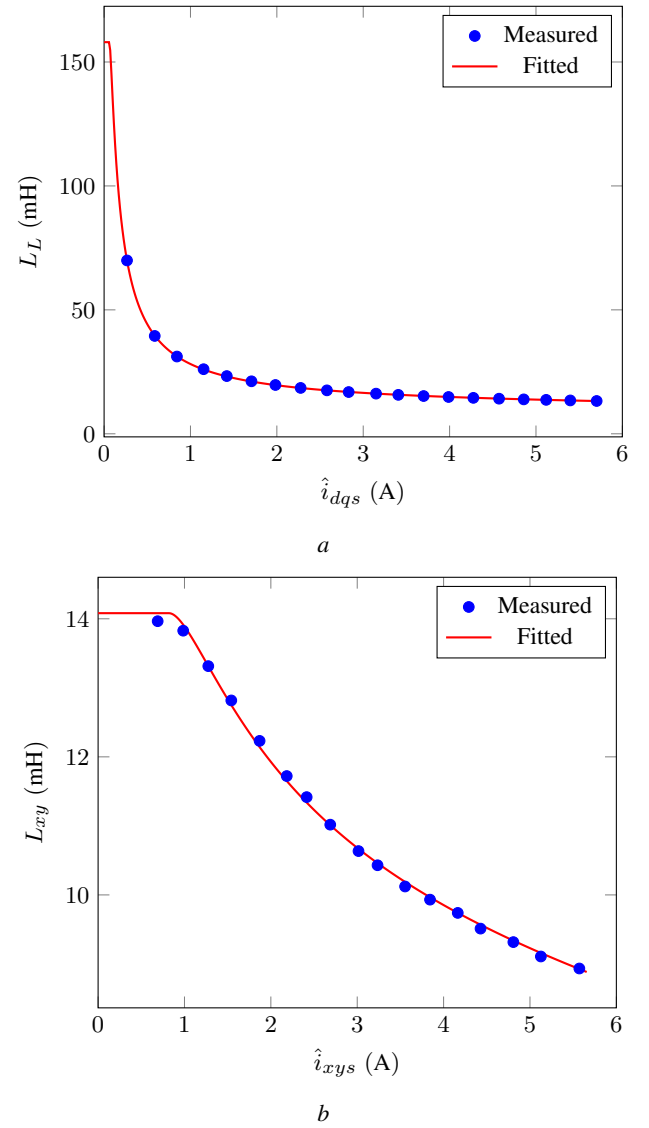
**Fig. 4:** Layout of the experimental setup: 1 – 6PIM, 2 – magnetic powder brake (load), 3 – three-phase AC-DC-AC converter, 4 – LEM sensors for measuring 6PIM phase currents, 5 – high-voltage sensing board for measuring DC link voltage, 6 – PED-Board controller for converter control and signal acquisition, 7 – PC

should be noted that the FEM results did not reveal non-linear phenomena in the  $xy$  plane aside from those related to IPCS, i.e. the  $xy$  plane inductance was constant. This is the consequence of insufficient information regarding the machine's magnetic circuit design. The values of resistances and unsaturated inductances are given in Table 1. The table also contains values of magnetizing current, flux and inductance corresponding to no-load operation at rated voltage. These values are denoted  $\hat{i}_{Ms}$ ,  $\hat{\psi}_{Ms}$  and  $\hat{L}_{Ms}$ , respectively, where the index  $s$  stands for “saturated”.

**Table 1** 6PIM parameter values

Parameter	Units	Value
$L_{Mu}$	mH	296
$\hat{i}_{Ms}$	A	2.4
$\hat{\psi}_{Ms}$	Wb	0.51
$L_{Ms}$	mH	210
$L_{Lu}$	mH	158
$L_{xyu}$	mH	14.1
$R_s$	$\Omega$	2.27
$R_R$	$\Omega$	1.83

Due to the IPCS effect, additional tests are needed for determining the corresponding flux changes in the  $dq$  and  $xy$  plane. The experimental procedure for obtaining these values is similar to that applied in the FE analysis of section 4.1. The values of  $dq$  and  $xy$  plane current components are adjusted to achieve different operating points. The motor is rotated at synchronous speed at all times during the test, hence it can be considered that  $\vec{i}_{dqs} \equiv \vec{i}_M$ . For each combination of currents,  $dq$  and  $xy$  plane flux linkage values are calculated



**Fig. 5:** Experimentally identified leakage inductances. The fitted values are obtained by curve fitting the measured values  
a  $dq$  plane leakage inductance  
b  $xy$  plane leakage inductance

as follows:

$$\vec{\psi}_{dqs} = \int (\vec{u}_{dqs} - R_s \vec{i}_{dqs}) dt \quad (16a)$$

$$\vec{\psi}_{xys} = \int (\vec{u}_{xys} - R_s \vec{i}_{xys}) dt \quad (16b)$$

After the flux linkages have been determined for all  $(\hat{i}_M, \hat{i}_{xys})$  combinations, the flux changes can be determined as described in section 4.1. However, in order to include the saturation of the  $xy$  plane inductance, the expression for the flux change in the  $xy$  plane needs to be modified. The total change of  $xy$  plane flux linkage with respect to its unsaturated value is expressed as follows:

$$\Delta \vec{\psi}'_{xy}(\hat{i}_M, \hat{i}_{xys}) = \vec{\psi}_{xys}(\hat{i}_M, \hat{i}_{xys}) - L_{xyu} \vec{i}_{xys}, \quad (17)$$

where  $L_{xyu}$  denotes the unsaturated value of the  $xy$  plane inductance. The flux change calculated by (17) includes the IPCS effect and the saturation caused solely by the  $xy$  plane current. The flux change in the  $dq$  plane is calculated according to (14).

The experimentally identified flux changes are presented in Figs. 6 and 7. Note that the blue dots in the figures represent the

values obtained from measurements, whereas the surfaces represent the interpolated values. The  $dq$  plane flux change caused by IPCS alone is displayed in Fig. 6. The experimental values are comparable to the FEM results shown in Fig. 2a. The total  $dq$  plane flux change displayed in Fig. 7a corresponds to the FEM results displayed in Fig. 3. As previously stated, the  $xy$  plane saturation observed from the experimental results was not present in the FEM model. Therefore, it can be considered that  $\Delta\psi_{xy} \approx \Delta\psi'_{xy}$  in terms of the FEM model, and the  $xy$  plane flux change shown in Fig. 7b can be related to Fig. 2b. The measurement results are fitted by appropriate analytical functions of  $(\hat{i}_M, \hat{i}_{xys})$  given in the Appendix (subsection 8.3). These functions are included in flux equations (15). Note that the  $xy$  flux linkage equation (15c) should now be rewritten as:

$$\vec{\psi}_{xys} = L_{xy}\vec{i}_{xys} + \Delta\vec{\psi}'_{xy}(\hat{i}_M, \hat{i}_{xys}), \quad (18)$$

which is actually an alternative formulation of (17).

The presented results confirm the conclusions regarding IPCS derived in section 4.1. It was therein assumed that the influence of the  $xy$  current component on the flux change in the  $dq$  plane can be neglected from a practical standpoint. The experimental results validate this assumption, as the flux change in the  $dq$  plane caused solely by IPCS was determined to be lower than the flux change due to main flux saturation by at least two orders of magnitude at each operating point. Therefore, it is sufficiently accurate to treat  $\Delta\psi'_{dq}$  as a function of the magnetizing current alone, as given in the Appendix (subsection 8.3).

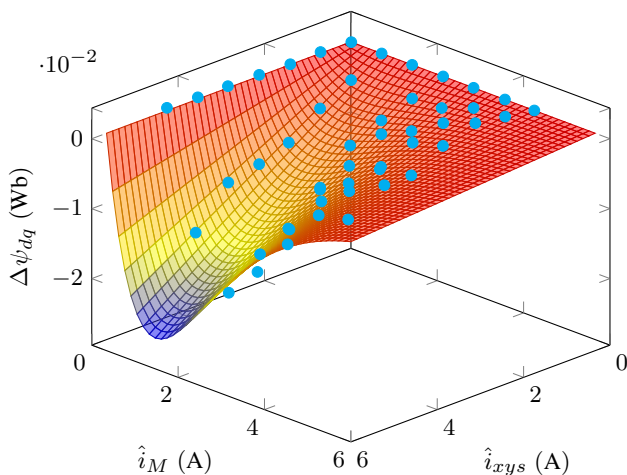
### 5.3 Experimental verification of the model

Based on the results of the FE analysis and the experimental parameter determination, a complete VSD model of the 6PIM including main flux saturation, leakage flux saturation, and IPCS effect, can be formed. The model equations are summarized as follows:

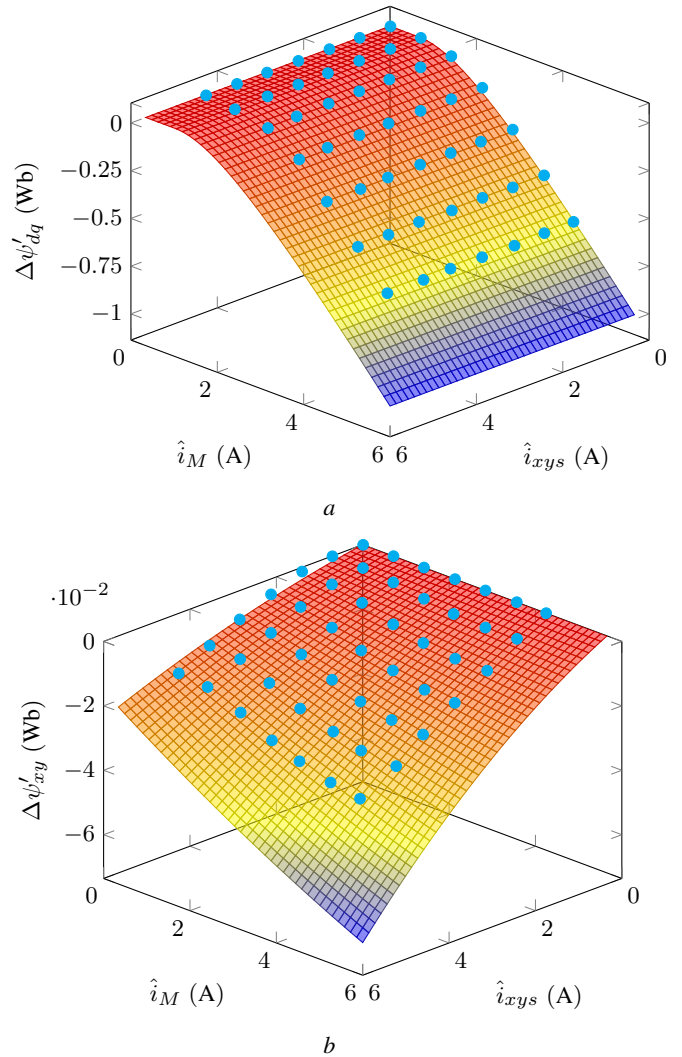
- voltage equations (1a), (8), (1c);
- flux-current relationships (15a), (15b), (18);
- torque equation (9);
- motion equation (5).

The 6PIM parameters are summarized in Table 1, and the non-linear flux changes and inductances are expressed by analytical functions given in the Appendix.

The proposed model was implemented in *Matlab* environment. To enable more efficient simulation of stationary regimes, the steady-state model was formed as well. The dynamic model equations are solved using integrated *Matlab* functions for numerical integration



**Fig. 6:** Experimentally identified  $dq$  plane flux change due to IPCS. The surface is obtained by fitting the experimental values, which are represented by the blue dots

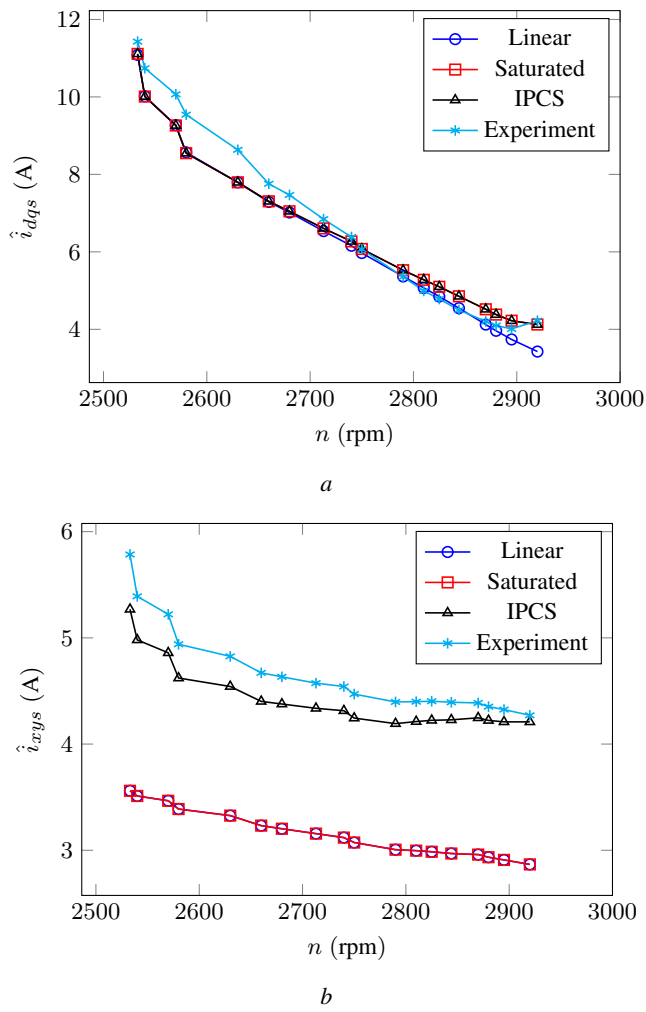


**Fig. 7:** Experimentally identified total flux changes due to main flux saturation and IPCS. The surfaces are obtained by fitting the experimental values, which are represented by the blue dots

a flux change in the  $dq$  plane  
b flux change in the  $xy$  plane

of DAEs, whereas the steady state model was solved using non-linear algebraic equation solvers. Operating modes with pronounced main flux saturation and unbalanced voltage supply were simulated, in order to emphasize the advantages of the new model compared to existing models. The results of the new model, referred to as the “IPCS model” further on, are compared to the results of the linear VSD model presented in section 3 and the model which includes main and leakage flux saturation, but ignores the IPCS effect. The latter will be referred to as the “saturated VSD model”. The linear model uses the constant value of magnetizing inductance equal to  $L_{Ms}$  (see Table 1) and the leakage inductance corresponding to rated current.

The steady state model was solved for different values of speed, i.e. mechanical load. First, the experiment was conducted where the 6PIM load was varied, and voltages, currents and the rotating speed values were obtained at each operating point. The obtained voltage and speed (slip) values were used as inputs to the steady-state model. The supply voltage components in the  $dq$  and  $xy$  planes were maintained at  $\hat{u}_{dqs} = 180$  V and  $\hat{u}_{xys} = 16$  V, respectively. The  $dq$  voltage magnitude is higher than the rated magnitude ( $\sqrt{2}U_{nf} = 168$  V), thereby insuring substantial saturation of the magnetic circuit. The selected value of the  $xy$  voltage component results in the  $xy$  plane current of approximately 70% of the rated value. Such test



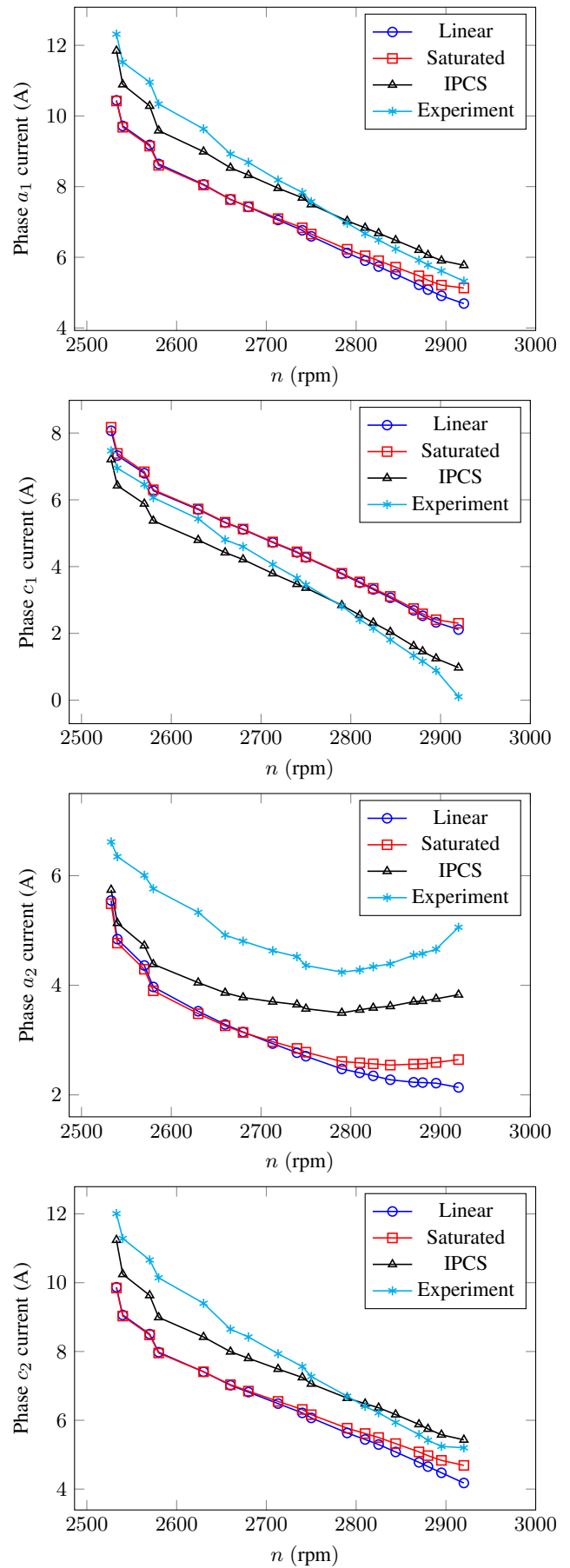
**Fig. 8:** Steady-state current magnitudes of the 6PIM obtained using the steady state model. The graphs display the variations of current magnitudes for different loading conditions

a  $dq$  plane current magnitude  
b  $xy$  plane current magnitude

conditions are chosen in order to verify the model performance under heavy unbalance.

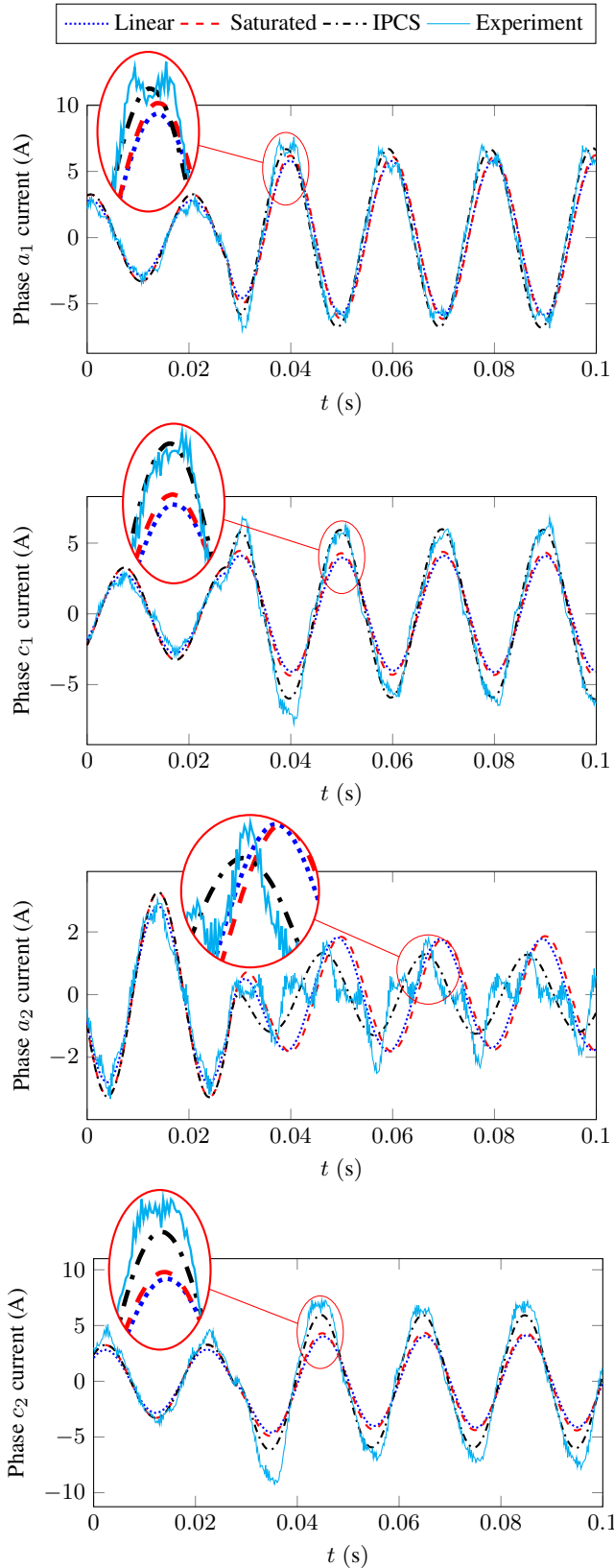
Values of  $dq$  and  $xy$  plane currents calculated from each model and the experimentally obtained values are plotted in Fig. 8. The  $dq$  current component calculated using the IPCS and saturated models shows better agreement with the experimental results compared to the linear model under light loads, where main flux saturation is the most pronounced. All three models exhibit an increased discrepancy between the measurements and calculated values under heavier loads, i.e. large slip values. This can be attributed to the non-linearity of the leakage inductance, which cannot be accounted for accurately enough using the  $\Gamma$  model. The  $xy$  current component calculated using the IPCS model is in much better agreement with the experimental values compared to the results of the linear and saturated models. Note that the values obtained using the linear and saturated models are identical, as the saturation effect in the  $xy$  plane is ignored altogether.

Steady state rms values of 6PIM phase currents are displayed in Fig. 9. Once again, the proposed IPCS model provides better agreement with the experimental results compared to the linear and saturated models. These results further emphasize the merits of the IPCS model. Note that the results of the linear model differ only slightly compared to the results of the saturated model. This is due to the fact that the linear model was formulated using the saturated (but constant) magnetizing inductance value, as previously stated. The



**Fig. 9:** Steady-state phase currents of the 6PIM (rms values) obtained using the steady state model. The graphs display the variations of current values for different loading conditions

traditional linear model which uses the unsaturated value of the magnetizing inductance would yield results which deviate significantly compared to the results of the saturated model.



**Fig. 10:** Instantaneous values of 6PIM phase currents under a sudden change of the  $xy$  voltage component. The enlarged details emphasize the accuracy of the IPCS model

The performance of the dynamic model is demonstrated by performing a sudden increase of the  $xy$  voltage component. The experimentally obtained voltages are used as inputs to the dynamic model. Fig. 10 displays the calculated and experimentally obtained phase current waveforms. The increase in the  $xy$  voltage component occurs at  $t = 28$  ms. The currents of the IPCS and saturated model are identical prior to the introduction of  $xy$  plane voltage, and are in good agreement with the experimental values. The linear model also provides a good approximation of experimental results under balanced supply, as it uses the values of magnetizing and leakage inductances corresponding to rated voltage and current, respectively. Following the voltage change, the IPCS model provides the most accurate results, both during the transient and in the new steady state. It should be noted that the currents obtained using the IPCS model best approximate the experimental results in terms of both phase and magnitude. The results of both the linear and saturated model deviate significantly from experimental results under unbalanced conditions.

The spectral analysis of experimental current waveforms reveals a presence of higher order harmonics, most notably the fifth and seventh. In a 6PIM, these components are projected into the  $xy$  plane [37]. As demonstrated in [12], the higher-order current harmonics in the  $xy$  subspace can be attributed to the air-gap flux distortion due to saturation. Another cause of low-order current harmonics is the inverter dead-time [35]. This effect is somewhat alleviated by the applied dead-time compensation. The presented research is aimed at analysing the fundamental frequency components of  $xy$  plane quantities, as stated in the Introduction. Therefore, the current distortion due to saturation and inverter dead-time was not accounted for.

## 6 Conclusion

A saturated VSD model of a 6PIM was presented. The model incorporated main and leakage flux saturation, as well as the mutual coupling between the  $dq$  and  $xy$  subspaces at the fundamental frequency. Fundamental frequency  $xy$  current components are not present in normal operation, but can be substantial in power sharing mode or in post-fault operation. The new model was aptly termed the IPCS model. The  $dq$  plane was modelled using the  $\Gamma$  model, similar to three-phase induction machines, which allowed determination of parameters directly from standard no-load and locked-rotor tests. An analysis of IPCS carried out using FEM revealed flux changes in the  $dq$  and  $xy$  plane. These flux changes were included in the flux-current relationships to account for the IPCS effect. An experimental method for determining the flux changes was proposed, and the obtained values were fitted using appropriate analytical functions. The IPCS model was then simulated in steady state and transient operating modes, under unbalanced supply conditions and with pronounced saturation. A comparison with the experimental results has demonstrated the advantages of the IPCS model compared to existing models.

Future research efforts will be aimed at improving the IPCS model by the inclusion of higher-order  $xy$  current harmonics inflicted by saturation. Specifically, in the case of 6PIM, the fifth and seventh harmonic should be included, as they are the most pronounced.

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## 8 Appendix

### 8.1 6PIM rated data and design parameters

The rated data of the 6PIM are given in Table 2. The rated data are estimated based on the design parameters. To accurately determine the rated current and power, a thermal test would need to be conducted. Table 3 contains the main dimensions and winding information of the 6PIM. The machine dimensions are denoted in the cross-section shown in Fig. 11. The rotor bars were inaccessible for measurement, therefore the rotor slot dimensions were estimated based on an analytical procedure proposed in [36]. The stator winding is a two-layer full-pitch distributed winding with two slots per pole and phase, as shown in Fig. 12.

### 8.2 FE model of the 6PIM

The FE model of the 6PIM developed according to the design information of Table 3 is used in the analysis of section 4.1. The material properties and maximal mesh sizes of individual blocks are given in Table 4. The highest mesh density is used in the air-gap area, where most of the magnetic field strength is localized. The mesh in the iron part is coarser compared to the air-gap, but is still limited due to potentially high field strength in saturated conditions. Due to the low flux density in the stator and rotor slot areas, it was acceptable to allow the solver to determine the mesh sizes in these areas automatically. Note that the mesh generator automatically reduces the element size where needed, e.g., in the vicinity of boundaries, corners and arcs.

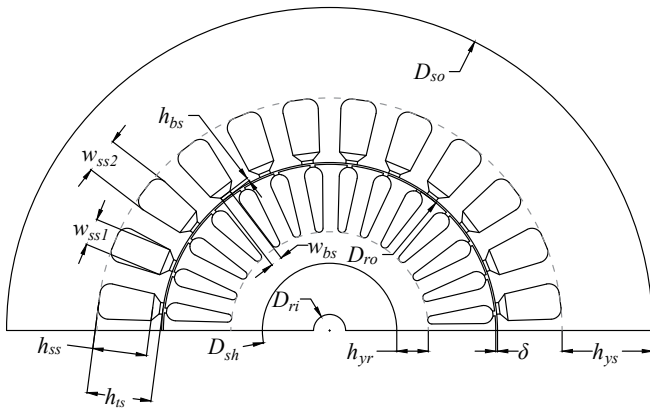
The cross-section of the FE model displaying the flux density distribution and the finite element mesh is given in Fig. 13. The displayed FE solution corresponds to the case where the main flux

**Table 2** 6PIM prototype rated data

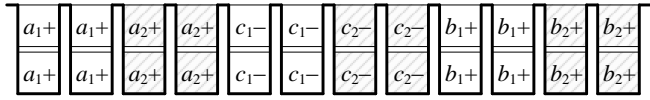
Parameter	Designation	Value
Rated phase voltage	$U_{nf}$	119 V
Rated frequency	$f_n$	50 Hz
Rated current	$I_n$	3.19 A
Rated power	$P_n$	1410 W

**Table 3** 6PIM prototype design parameters

Parameter	Designation	Value
Stator outer diameter	$D_{so}$	125 mm
Rotor outer diameter	$D_{ro}$	64.3 mm
Rotor shaft diameter	$D_{sh}$	26 mm
Rotor inner diameter	$D_{ri}$	6.2 mm
Air-gap length	$\delta$	0.35 mm
Stator yoke height	$h_{ys}$	17.5 mm
Rotor yoke height	$h_{yr}$	6.1 mm
Stator tooth height	$h_{ts}$	12.5 mm
Stator slot height	$h_{ys}$	10.7 mm
Stator slot width (inside)	$w_{ss1}$	5.2 mm
Stator slot width (outside)	$w_{ss2}$	6.9 mm
Stator slot opening width	$w_{bs}$	2.2 mm
Stator slot opening height	$h_{bs}$	0.6 mm
Axial length	$l_a$	90 mm
No. of stator slots	$Q_s$	24
No. of rotor slots	$Q_r$	28
No. of pole pairs	$p$	1
Stator winding connection	/	Y
Turns per phase	$N_{ph}$	104
Winding type	/	two-layer
Winding pitch	$y$	12



**Fig. 11:** Cross-section of the 6PIM magnetic circuit with the main dimensions denoted

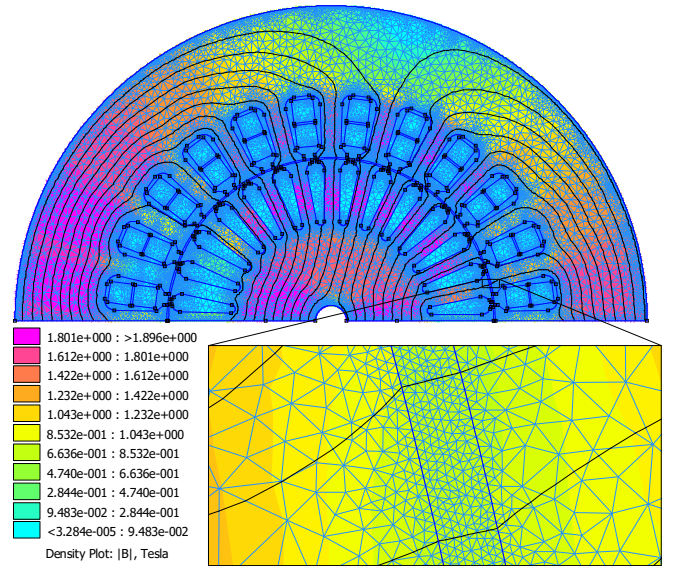


**Fig. 12:** Schematic display of the stator winding spanning one pole pitch. The + and - signs indicate the direction of the coil sides

**Table 4** Material properties and maximal mesh sizes of FEM block: M-19 Steel – laminated non-oriented steel, 1010 Steel – solid steel

Machine part	Material	Mesh size
Stator magnetic circuit	M-19 Steel	1 mm
Rotor magnetic circuit	M-19 Steel	1 mm
Rotor shaft	1010 Steel	1 mm
Stator winding	Copper ( $\varnothing$ 0.75 mm)	Auto
Rotor bars	Aluminium	Auto
Air-gap	Air	0.05 mm

saturation and IPCS are both highly pronounced. The FE simulation was conducted using 85000 nodes (168990 elements). The mesh in the air-gap area is shown in the enlarged detail of Fig. 13. To prove the accuracy of the results obtained using the selected mesh size, the mesh in the air-gap and the iron parts of the machine was refined, increasing the number of nodes to 190000. This modification resulted in an increase in accuracy of no more than 0.2% in terms of per-phase flux linkages, thus proving that the selected mesh density guarantees satisfying accuracy.



**Fig. 13:** Solution of 6PIM analysis in the FEA software for ( $\hat{i}_M = 6$  A,  $\hat{i}_{xys} = 6$  A) with the flux density distribution and mesh displayed. The enlarged detail shows the air-gap mesh

### 8.3 Non-linear model parameters

The  $dq$  subspace leakage inductance is given as a two-segment non-linear function of the stator current in (19). As noted, expressing the leakage inductance solely as the function of stator current is suggested in [33]. The  $dq$  flux change is expressed as a two-segment non-linear function of the magnetizing current in (20). The influence of IPCS on the  $dq$  flux change, albeit present, is negligible from a practical standpoint. The  $dq$  flux change is therefore expressed in terms of the magnetizing current alone, disregarding the influence of the  $xy$  current component. The flux change in the  $xy$  plane is expressed by (21), as a two-dimensional function of the magnetizing and  $xy$  plane currents. The values of constant parameters  $L_{Lu}$  and  $L_{Mu}$  featured in (19) and (20) are given in Table 1.

$$L_L(\text{mH}) = \begin{cases} L_{Lu}, & \hat{i}_{dqs} < 0.057 \text{ A} \\ -0.5219 \cdot \hat{i}_{dqs}^{-2} + 17.52 \cdot \hat{i}_{dqs}^{-1} + 11.37 - 0.2121 \cdot \hat{i}_{dqs}, & \hat{i}_{dqs} \geq 0.057 \text{ A} \end{cases} \quad (19)$$

$$\Delta\psi'_{dq}(\text{Wb}) = \begin{cases} 0, & \hat{i}_M < 0.679 \text{ A} \\ \left( (1.242 + 1.691 \cdot \hat{i}_M^{-1} + 0.5723 \cdot \hat{i}_M^{-2})^{-1} - L_{Mu} \hat{i}_M \right), & \hat{i}_M \geq 0.679 \text{ A} \end{cases} \quad (20)$$

$$\Delta\psi'_{xy}(\text{Wb}) = - \left( 5.56 \cdot \hat{i}_{xys} + 0.6733 \cdot \hat{i}_{xys}^2 \right) \cdot \left( 4.168 + 1.787 \cdot \hat{i}_M - 0.0516 \cdot \hat{i}_M^2 \right) \cdot 10^{-2} \quad (21)$$