

Magnetic dynamos in white dwarfs – I. Explaining the dearth of bright intermediate polars in globular clusters

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ABSTRACT

Recently, Bahramian et al. investigated a large sample of globular clusters (GCs) and found that bright intermediate polars (IPs) are a factor of 10 less frequent in GCs than in the Galactic field. We theoretically investigate here this discrepancy based on GC numerical simulations. We found that, due to disruptive dynamical interaction, there is on average a reduction of only half of bright IP progenitors, which is clearly not enough to explain the observed deficiency. However, if the rotation- and crystallization-driven dynamo scenario recently proposed by Schreiber et al. is incorporated in the simulations, the observed rareness of bright IPs in GCs can be reproduced. This is because bright cataclysmic variables (CVs) in GCs are typically very old systems ($\gtrsim 10$ Gyr), with white dwarfs that almost fully crystallized before mass transfer started, which does not allow strong magnetic fields to be generated. The observed mass density of bright IPs in GCs can be recovered if around one-third of the bright CVs dynamically formed through mergers have magnetic field strengths similar to those of IPs. We conclude that the observed paucity of bright IPs in GCs is a natural consequence of the newly proposed rotation- and crystallization-driven dynamo scenario.

Key words: methods: numerical – binaries: general – stars: magnetic field – novae, cataclysmic variables – globular clusters: general.

1 INTRODUCTION

The study of globular clusters (GCs) plays an important role in our understanding of the Universe since they are natural laboratories for testing theories of stellar dynamics and evolution, very old (~ 10 – 12 Gyr) and have high central stellar densities (up to $\sim 10^6$ stars per pc^3). These properties make stellar dynamical interactions very frequent in these environments, which turns GCs into factories of exotic systems such as low-mass X-ray binaries, red and blue straggler stars, millisecond pulsars, and black hole binaries, among others. One of the most abundant types of interacting binaries in GCs are cataclysmic variables (CVs), which are characterized by a white dwarf (WD) that stably accretes matter from a low-mass main-sequence (MS) star (see Knigge 2012; Belloni & Rivera Sandoval 2021, for comprehensive reviews on GC CVs).

Intermediate polars (IPs) are typically brighter (with X-ray luminosities $\sim 10^{33}$ erg s^{-1} being fairly common for IPs with periods longer than 3 h and rare for other classes) and have harder spectra than other classes of CVs (again with kT typically tens of keV for IPs,

and only rarely so high for other classes). Recently, Bahramian et al. (2020) thoroughly investigated 38 GCs with the *Chandra* satellite, provided an extensive catalogue containing more than 1100 X-ray sources in these GCs and analysed in detail those sources with X-ray luminosities greater than $\sim 10^{33}$ erg s^{-1} , which includes bright IPs. Based on the mass density of IPs in the Galactic disc, obtained from the space density (Schwope 2018) and assuming a local stellar density of 0.1 pc^{-3} , Bahramian et al. (2020) estimated that 16–34 bright IPs should be present in their sample of GCs. However, they found that only 2–3 sources in those GCs are likely bright IPs. This represents an underabundance of a factor of 10, at more than 3σ significance, in comparison with the Galactic field. This goes in the opposite direction of previous observational works, which suggested that most GC CVs should be IPs (e.g. Grindlay et al. 1995).

There are currently four scenarios that have been suggested for the origin of magnetic fields in WD populations in the Galactic field. In the *fossil field scenario*, the magnetic WD progenitor is most likely the product of a MS–MS binary merger. Such a merger can produce strong differential rotation and give rise to large-scale fields in the radiative envelopes of the merger outcome (Ferrario et al. 2009; Schneider et al. 2019). In the *common-envelope dynamo scenario*, the magnetic field is generated when two stars merge during common-envelope evolution. In this case, the dynamo is expected to be driven

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by shear due to differential rotation either in the envelope (e.g. Potter & Tout 2010), or in an accretion disc (Nordhaus et al. 2011), or in the hot outer layers of the degenerate core (Wickramasinghe, Tout & Ferrario 2014). In the *double degenerate scenario*, a magnetic field is generated during the merger of two degenerate cores (García-Berro et al. 2012). More recently, Schreiber et al. (2021) revised the model suggested by Isern et al. (2017) and proposed that a *rotation- and crystallization-driven dynamo* similar to those working in planets and low-mass stars can generate strong magnetic fields in the WDs in close binaries.

Utilizing the simulations carried out by Belloni et al. (2019), which corresponds to the largest set of GC models used to study CVs to date, we here investigate the possible reasons leading to the observed low incidence of bright IPs in GCs, taking into account the effect of dynamics in creating and destroying IP progenitors as well as the above-mentioned scenarios for the origin of magnetic WDs.

2 GLOBULAR CLUSTER MODELS

In order to investigate bright IPs in GCs, we used the models simulated by Belloni et al. (2019) with the MOCCA code (Hypki & Giersz 2013, and references therein), which reproduces N -body direct simulations with great precision and accuracy (e.g. Madrid et al. 2017). One very important part of MOCCA for this work is the binary star evolution (BSE) code (Hurley, Tout & Pols 2002), which was properly upgraded and calibrated by Belloni et al. (2018) to deal with CV evolution.

Amongst our GC models, which have low metallicity ($Z = 0.001$) and assumed to be ≈ 12 Gyr old, we have a variety of different initial conditions spanning different values of the mass, the size, the King concentration parameter, the initial binary population, the binary fraction, and the Galactocentric distance. In addition, we also explored two parameters of stellar/binary evolution, namely inclusion or not of mass fallback for black hole formation and three different common-envelope efficiencies.

We compare the distributions of the central surface brightness and the core radius of our models with those of observed GCs in Fig. 1. These models are consistent with a substantial fraction of real GCs and can be considered as representative for the Galactic GC population. In addition, models assuming the so-called Kroupa initial binary population (Belloni et al. 2017b, and references therein) and with low common-envelope efficiency (≤ 0.50) better reproduce the observed numbers of CVs and CV candidates in NGC 6397, NGC 6752, and 47 Tuc. Taking this into account, from the 288 GC models investigated in Belloni et al. (2019), we here only consider this subset, which correspond to 96 GC models. Fig 1 also illustrates that, except for a very few bright compact models, this subset of models is consistent with the sample of GCs analysed by Bahramian et al. (2020).

3 BRIGHT CATAclysmic VARIABLES IN GLOBULAR CLUSTERS

In Belloni et al. (2019), we investigated the impact of dynamics in creating CVs and destroying CV progenitors, taking into account the entire population of detectable CVs in the GC models. We repeat here these analyses, but focusing only on *bright CVs*, which are assumed here to be those located above the orbital period gap, since the X-ray bright IPs are dominant in this region of the orbital period distribution (e.g. Suleimanov, Doroshenko & Werner 2019). In particular, we investigate whether effective dynamical disruption

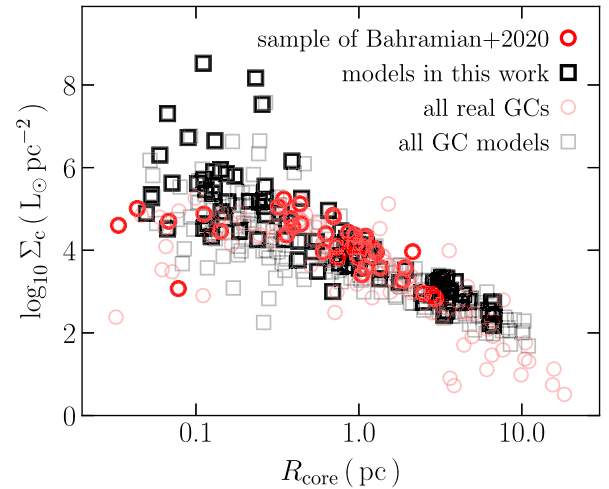


Figure 1. Central surface brightness (Σ_c) against the core radius (R_{core}), for the MOCCA models (squares) and real GCs (circles), which observational data was extracted from Harris (1996, updated 2010). The figure illustrates that the models considered in this work are consistent with the sample of GCs analysed by Bahramian et al. (2020), except for just a few very bright compact models.

of bright IP progenitors could explain the paucity of bright IPs in GCs.

3.1 Dynamical disruption of progenitors

Belloni et al. (2019, see their fig. 2) found a strong correlation between the initial GC stellar encounter rate and the fraction of destroyed CV progenitors. Since bright IPs are mostly concentrated above the orbital gap, these systems have to come from MS–MS binaries with relatively long initial orbital periods ($\sim 10^3 - 10^5$ d). This is because the pre-CV formation rate is much higher at the beginning of the GC evolution (within ~ 2 Gyr after the GC formation) and only systems that start the CV phase close to the present day (i.e. spend ~ 10 Gyr in the pre-CV phase) will likely be found above the orbital period gap. These systems have to leave the common-envelope evolution with longer orbital periods than those identified below the orbital period gap. Their progenitors are therefore easier to be disrupted in dynamical interactions.

We found that the destruction fractions range from ~ 10 per cent (sparsest model) to ~ 80 per cent (densest model), and the average is ~ 60 per cent, which corresponds to an average reduction of a bit more than a factor of 2. Even though such a destruction rate of bright CV progenitors considerably reduces the incidence of these systems in GCs compared to the field, it is clearly not enough to explain the observed deficit of a factor of 10 of bright IPs. Therefore, *the reason behind the observed low incidence of bright IPs in GCs must be related to the different characteristics of GC CVs, in comparison with their counterparts in the Galactic disc.*

3.2 Present-day population: formation channels and mass densities

As dynamical disruptive interactions of CV progenitors cannot explain the small number of observed bright IPs in GCs, we now take a look at the different formation channels of the predicted bright CVs with the aim of identifying which of these channels might perhaps not produce magnetic systems. Of all bright CVs, common-envelope evolution of binaries undergoing no or very weak dynamical

Table 1. Relative frequency of bright CVs averaged over all our 96 GC models, separated by their formation channels. In the last column, we provide their mass density, having in mind that the observed mass density of bright IPs in GCs is $\sim 2 \times 10^{-7} M_{\odot}^{-1}$.

Formation channel	Number	Fraction	Mass density (M_{\odot}^{-1})
Common envelope	272	88.9 per cent	7.10×10^{-6}
Exchange	11	3.6 per cent	2.87×10^{-7}
Merger	23	7.5 per cent	6.00×10^{-7}
All	306	100.0 per cent	7.99×10^{-6}

interactions is the dominant CV formation channel, producing on average ~ 89 per cent of bright CVs. The remaining ~ 11 per cent were formed because of at least one strong dynamical interaction.

This predicted low-average fraction of dynamically formed bright CVs can be compared with observational X-ray investigations considering the GC stellar encounter rates, which depends on the central density, the core radius, and the central velocity dispersion, and can be used to address whether a particular type of X-ray source is predominantly created in GCs through dynamical interactions. Our results seem to be in agreement with the *Chandra* study performed by Cheng et al. (2018), who used a sample of 69 GCs, but to contradict previous *Chandra* investigations, which considered much smaller GC samples (e.g. Pooley & Hut 2006). We restrict ourselves here to only mentioning that this sort of observational investigation should be taken with a grain of salt, since (i) the number of detected X-ray sources is very sensitive to the exposure time, (ii) these investigations are strongly X-ray biased and contamination from other types of sources are not rare, (iii) the properties of a particular GC might significantly change over its lifetime, (iv) the size of the investigated sample does matter. We can then conclude that even though such X-ray studies are helpful, follow-up researches involving multiwavelength observations, as homogeneous as possible, of a large sample of GCs are needed to fully investigate the dynamical origin of CVs in GCs. A detailed discussion on this topic is provided in Belloni & Rivera Sandoval (2021, their sections 2.4 and 4.2).

Amongst the several dynamical formation channels for bright CVs in our simulations, the most prominent is three-body dynamical merger involving a single MS star and an MS–MS binary, leading to an MS–MS binary which later undergoes common-envelope evolution to become a WD–MS binary. This formation channel corresponds to ~ 47 per cent of the dynamically formed CVs. The second most efficient dynamical channel, corresponding to ~ 32 per cent of them, is three-body dynamical exchange in which a WD replaces the lighter MS star in an MS–MS binary, leading to a WD–MS binary. The third most frequent channel involves mergers of WDs with MS stars in three-body or four-body interactions, which provides ~ 12 per cent of the dynamically formed CVs. The last dynamical formation channel for bright CVs is either MS–MS or WD–MS merger in binary evolution, followed by dynamical exchange in three-body interactions, contributing to ~ 9 per cent of dynamically formed bright CVs. The number of bright CVs, separated by their main formation channels, as well as their mass densities are summarized in Table 1.

According to the available scenarios for the origin of magnetic WDs, we expect that a considerable portion of the bright CVs formed through mergers, which corresponds to ≈ 7 per cent of the bright CVs, could be magnetic (e.g. García-Berro et al. 2012; Wickramasinghe et al. 2014; Schneider et al. 2019). The mass density of bright CVs formed through dynamical mergers is $6 \times 10^{-7} M_{\odot}^{-1}$, which is about three times larger than the observed mass density of bright

IPs in GCs ($\sim 2 \times 10^{-7} M_{\odot}^{-1}$). Since potentially a large number of these systems are magnetic, this formation channel alone can explain the amount of observed bright IPs in GCs. On the other hand, mass densities of bright CVs not formed due to mergers is more than 30 times larger than the observed mass density of bright IPs. This implies that the bright CVs formed from either common-envelope evolution or exchange, which corresponds to ≈ 93 per cent of the bright CVs, need to be virtually all non-magnetic in order to reproduce the observed rareness of bright IPs in GCs.

4 IF NOT DYNAMICS, WHAT IS CAUSING THE SCARCITY OF BRIGHT INTERMEDIATE POLARS IN GLOBULAR CLUSTERS?

We have shown in Section 3 that on average the vast majority of bright CVs in all our GC models form through common-envelope evolution and that virtually all of them need to be non-magnetic in order to explain the observed dearth of bright IPs in GCs. Considering that the rotation- and crystallization-driven dynamo scenario is the only currently available scenario that can explain the characteristics of magnetic WDs in close binaries in the Galactic field, we investigate in what follows the implications of this model for the ≈ 93 per cent of GC bright CVs that were not formed through dynamical mergers. In other words, we will verify whether this scenario is also able to explain the characteristics of these systems in GCs, which implies that the model needs to predict that nearly all bright CVs not formed through dynamical mergers are non-magnetic. We start with a brief review of the recently proposed dynamo scenario.

4.1 The rotation- and crystallization-driven dynamo

Crystallization of carbon–oxygen WDs leads to the formation of an oxygen rich solid core surrounded by a convective carbon rich liquid. As first noticed by Isern et al. (2017), this configuration of a solid core embedded in a convective liquid could generate strong magnetic fields if the WD is rotating rapidly in full analogy to the mechanism responsible for generating the magnetic fields of planets and low-mass stars (Christensen 2010).

Schreiber et al. (2021) extended the work of Isern et al. (2017) by incorporating the dynamo in evolutionary tracks of CVs taking into account the WD spin period and orbital period evolution. The WD spin period is usually very long after the WD formation, but can become very short, near break-up spin period, when the WD is accreting mass and angular momentum, as in CVs.

The resulting scenario for the formation and evolution of magnetic CVs proposed by Schreiber et al. (2021) can be summarized as follows. The WDs in close post-common-envelope binaries are born non-magnetic. Due to angular momentum losses, this non-magnetic detached binary evolves towards shorter orbital periods, while the WD cools and its core may eventually starts to crystallize. When the secondary fills its Roche lobe, the onset of mass transfer occurs and the binary becomes a CV. Since the WD accretes mass and angular momentum, it spins up. If the WD core is crystallizing, the conditions for the rotation- and crystallization-driven dynamo are satisfied and the magnetic field is generated. Only if the WD is already substantially crystallized before mass transfer starts or if it is still too young for the onset of crystallization, the conditions are not mutually met and the CV will remain non-magnetic.

As described in detail in Schreiber et al. (2021), the new scenario can solve several long-standing problems in the Galactic field at once, namely the absence of magnetic WDs amongst young close detached binaries (Parsons et al. 2021), the higher/lower incidence

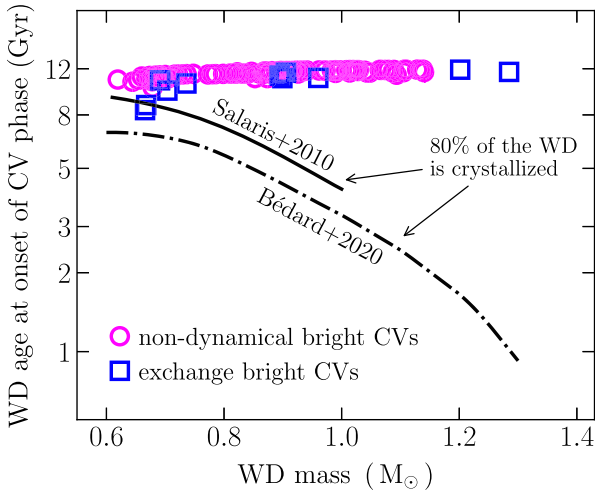


Figure 2. WD ages at onset of CV phase against their masses for bright CVs formed either without strong dynamical interactions (magenta circles) or due to dynamical exchanges (blue squares). The solid and dot-dashed lines indicate when 80 per cent of the WD matter has crystallized, according to the models computed by Salaris et al. (2010) and Bédard et al. (2020), respectively. Strong magnetic field generation due to the rotation- and crystallization-driven dynamo is expected for those systems lying well below these lines. Clearly, the new scenario does not predict any magnetic CV in this population.

of magnetism in CVs/detached binaries than in single WDs (Kepler et al. 2013; Pala et al. 2020; Parsons et al. 2021), the existence of the radio pulsar AR Sco (Marsh et al. 2016), and the relative frequency of magnetic and non-magnetic CVs (Pala et al. 2020). Given that the new scenario is successful for close detached WD binaries in the Galactic field, in what follows, we test whether this scenario can also explain the dearth of bright IPs in GCs.

4.2 Applying the rotation- and crystallization-driven dynamo scenario to cataclysmic variables in globular clusters

In the context of the rotation- and crystallization-driven dynamo scenario, two major differences between GCs and the Galactic disc need to be considered. First, while GCs have measured ages of ~ 11 – 12 Gyr, the Galactic disc is younger (~ 10 Gyr old). The second is related to the different history of star formation. In GCs, we expect one or more bursts of star formation within the first 10^8 yr of the cluster life (e.g. Bastian & Lardo 2018). On the other hand, in the Milky Way disc, assuming a nearly uniform star formation rate until the present time reasonably well explains several observational constraints (e.g. Schulz, Pflamm-Altenburg & Kroupa 2015). Taking into account the different ages and star formation rates, we should expect somewhat different CV populations in the Galactic disc and in GCs. In particular, CVs in GCs should be much older than those in the Galactic disc (see also Belloni et al. 2017a).

We show in Fig. 2, the WD age at onset of mass transfer for bright CVs formed either from isolated binaries or due to dynamical exchanges and compare it with the age at which 80 per cent of the WD matter has crystallized, according to the hydrogen-atmosphere WD models computed by Salaris et al. (2010) and Bédard et al. (2020). The differences between the ages in both works are mainly due to the different physics inputs and carbon–oxygen chemical stratifications, and the fact that the energy release due to phase separation upon crystallization is not included in Bédard et al. (2020) calculations.

According to the model proposed by Schreiber et al. (2021), a sufficiently strong magnetic field due to a rotation- and crystallization-driven dynamo is expected for those systems having partially crystallized WDs, with crystallized mass fractions below ~ 80 per cent. Fig. 2 clearly illustrates that all bright CVs formed without involving mergers are too old to develop any strong magnetic field. Despite the fact that accretion heating during nova cycles could in principle slowly melt a fraction of the crystallized core, this process is expected to take several hundred Myr (e.g. Epelstain et al. 2007). Given that this time-scale is longer than that of CV evolution above the orbital period gap, we expect a negligible number of magnetic bright CVs (if any) assuming the fields are generated by the rotation- and crystallization-driven dynamo scenario. Taking into account that these channels correspond to ≈ 93 per cent of all bright CVs, the scenario for the origin of magnetic WDs in close binaries predicts a significant underabundance of bright IPs in GCs.

Therefore, this new scenario for WD magnetic field generation naturally solves the issue investigated here, and even independently on the impact of dynamics in GC modelling, since even for initially sparse clusters, in which dynamics contribute to only ~ 10 per cent of reduction, the new scenario would still guarantee that the vast majority of the population of bright CVs would be non-magnetic. This implies that *the rotation- and crystallization-driven dynamo scenario is not only able to explain the characteristics of magnetic CVs in the Galactic disc, but also of those in GCs.*

5 BRIGHT INTERMEDIATE POLARS FROM MERGERS

So far, we have shown that no bright IP is expected in the population of bright CVs formed either from primordial binary evolution or through dynamical exchanges if the new dynamo model is correct. Therefore, *we expect that the observed population of bright IPs in GCs should be all dynamically formed through either MS–MS mergers or mergers involving WDs.*

In case of an MS–MS merger, Schneider et al. (2019) showed that strong, large-scale surface magnetic fields can be produced in coalescing MS binaries. In cases of mergers involving WDs, mostly WD–MS and WD–WD mergers, a dynamo can be produced due to differential rotation in the outer layers of the merger product, resulting in strong magnetic fields (e.g. García-Berro et al. 2012; Wickramasinghe et al. 2014). However, we shall emphasize that it is not clear under which condition the magnetic field strengths generated in mergers would be in the range of those observed amongst IPs (~ 1 – 10 MG), since the fields could be either stronger or weaker, which would not make them as bright in X-rays as IPs. For instance, a correlation between the differential rotation and the field strength is expected in mergers involving WDs (e.g. Wickramasinghe et al. 2014). While objects differentially rotating near break-up velocities are expected to develop the strongest fields, only weak fields should be generated with less differential rotation. The situation is similar for MS–MS mergers. It is not clear how strong the magnetic field generated in MS–MS mergers could be in the resulting WD.

An example for the formation of weakly magnetized WDs in close binaries most likely through MS–MS merger is V471 Tau. This post-common-envelope binary hosts a K-type MS star paired with a magnetic WD, which is the hottest, youngest, and most massive WD in the Hyades open cluster. In order to explain this clear evolutionary contradiction, O’Brien, Bond & Sion (2001) proposed that such a WD is descended from a blue straggler, with mass about twice the turnoff mass, which was the product of an MS–MS merger. When this star evolved and became a red giant, it filled its Roche lobe

and the binary underwent common-envelope evolution, which led the binary to shrink to its present orbital period. The weak magnetic field of the WD in V471 Tau clearly excludes this system to become a bright IP. This example illustrates that, despite mergers likely lead to magnetic field generation, many of the emerging magnetic WDs may not become IPs. Therefore, the predicted mass density of bright CVs formed through mergers ($6 \times 10^{-7} M_{\odot}^{-1}$) should be treated as an upper limit for the predicted mass density of bright IPs.

Also the observed mass density of bright IPs in GCs of $\sim 2 \times 10^{-7} M_{\odot}^{-1}$ (Bahramian et al. 2020) should be taken with a grain of salt. These authors only analysed 38 GCs, which is a relatively small sample of the entire Galactic GC population, and it has been already shown that results from X-ray investigations can strongly depend on the size of the sample (Cheng et al. 2018). In addition, they only provide bright IP candidates in their catalogue, lacking spectroscopic confirmation (or, at least, proper characterization of these candidates, e.g. by obtaining orbital periods), which might change the estimated mass density upon confirmation or not of some of these candidates. Taking into account these uncertainties and given that the mass density estimated from observation is only a factor of 3 below the upper limit predicted by our simulations, we conclude that model predictions and observations are in very good agreement. The observed mass density could be explained, provided that roughly one-third of the bright CVs formed through mergers have field strengths consistent with IPs, which appears reasonable.

6 SUMMARY AND CONCLUSIONS

We investigated bright CVs and IPs formed through either isolated binary evolution or dynamical exchange/merger interactions, in 96 GC models evolved with the MOCCA code with the aim of explaining the observed underabundance (of a factor of 10) of bright IPs in GCs compared to the Galactic disc.

We showed that dynamical disruption of bright IP progenitors in GCs is not enough to explain the observed dearth of bright IPs, since dynamics only provide on average a reduction of a bit more than a factor of 2. Many progenitor binaries survive in these dense environments and successfully become CVs, in addition to other CVs that are formed because of dynamical exchange. These two groups correspond to ≈ 93 per cent of all bright CVs and belong to a predominantly old population, in which their WDs are almost fully crystallized before the onset of mass transfer. According to the recently proposed rotation- and crystallization-driven dynamo scenario, these characteristics do not allow strong magnetic field generation, which then explain the observed scarcity of bright IPs in GCs.

Finally, the remaining ≈ 7 per cent of bright CVs were dynamically formed through mergers and likely harbour magnetic WDs. The mass density of this group is about three times larger than the observed mass density of bright IPs. Thus, if about one-third of such CVs have magnetic field strengths comparable to those observed in IPs, the observed mass density of bright IPs can be recovered.

We conclude that the newly proposed rotation- and crystallization-driven dynamo scenario for the generation of magnetic fields in close WD binaries naturally explains the observed paucity of bright IPs in GCs. Additionally, the few bright IPs currently observed should be dynamically formed through mergers, involving either WDs or MS stars that later become WDs.

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DATA AVAILABILITY

The data underlying this article can be obtained upon request.

REFERENCES

- Bahramian A. et al., 2020, *ApJ*, 901, 57
 Bastian N., Lardo C., 2018, *ARA&A*, 56, 83
 Bédard A., Bergeron P., Brassard P., Fontaine G., 2020, *ApJ*, 901, 93
 Belloni D., Rivera Sandoval L., 2021, in F. Giovannelli, L. Sabau-Graziati, eds, *The Golden Age of Cataclysmic Variables and Related Objects V.*, PoS-SISSA Palermo, Italy, p. 13
 Belloni D., Giersz M., Rocha-Pinto H. J., Leigh N. W. C., Askar A., 2017a, *MNRAS*, 464, 4077
 Belloni D., Askar A., Giersz M., Kroupa P., Rocha-Pinto H. J., 2017b, *MNRAS*, 471, 2812
 Belloni D., Schreiber M. R., Zorotovic M., Iłkiewicz K., Hurley J. R., Giersz M., Lagos F., 2018, *MNRAS*, 478, 5639
 Belloni D., Giersz M., Rivera Sandoval L. E., Askar A., Ciecielag P., 2019, *MNRAS*, 483, 315
 Cheng Z., Li Z., Xu X., Li X., 2018, *ApJ*, 858, 33
 Christensen U. R., 2010, *Space Sci. Rev.*, 152, 565
 Epelstain N., Yaron O., Kovetz A., Prialnik D., 2007, *MNRAS*, 374, 1449
 Ferrario L., Pringle J. E., Tout C. A., Wickramasinghe D. T., 2009, *MNRAS*, 400, L71
 García-Berro E. et al., 2012, *ApJ*, 749, 25
 Grindlay J. E., Cool A. M., Callanan P. J., Bailyn C. D., Cohn H. N., Lugger P. M., 1995, *ApJ*, 455, L47
 Harris W. E., 1996, *AJ*, 112, 1487
 Hurley J. R., Tout C. A., Pols O. R., 2002, *MNRAS*, 329, 897
 Hupki A., Giersz M., 2013, *MNRAS*, 429, 1221
 Isern J., García-Berro E., Külebi B., Lorén-Aguilar P., 2017, *ApJ*, 836, L28
 Kepler S. O. et al., 2013, *MNRAS*, 429, 2934
 Knigge C., 2012, *Mem. Soc. Astron. Ital.*, 83, 549
 Madrid J. P., Leigh N. W. C., Hurley J. R., Giersz M., 2017, *MNRAS*, 470, 1729
 Marsh T. R. et al., 2016, *Nature*, 537, 374
 Nordhaus J., Wellons S., Spiegel D. S., Metzger B. D., Blackman E. G., 2011, *Proc. Natl. Acad. Sci.*, 108, 3135
 O’Brien M. S., Bond H. E., Sion E. M., 2001, *ApJ*, 563, 971
 Pala A. F. et al., 2020, *MNRAS*, 494, 3799
 Parsons S. G., Gänsicke B. T., Schreiber M. R., Marsh T. R., Ashley R. P., Breedt E., Littlefair S. P., Meusinger H., 2021, *MNRAS*, 502, 4305
 Pooley D., Hut P., 2006, *ApJ*, 646, L143
 Potter A. T., Tout C. A., 2010, *MNRAS*, 402, 1072
 Salaris M., Cassisi S., Pietrinferni A., Kowalski P. M., Isern J., 2010, *ApJ*, 716, 1241
 Schneider F. R. N., Ohlmann S. T., Podsiadlowski P., Röpke F. K., Balbus S. A., Pakmor R., Springel V., 2019, *Nature*, 574, 211
 Schreiber M. R., Belloni D., Gänsicke B. T., Parsons S. G., Zorotovic M., 2021, *Nat. Astron.*
 Schulz C., Pflamm-Altenburg J., Kroupa P., 2015, *A&A*, 582, A93
 Schwope A. D., 2018, *A&A*, 619, A62
 Suleimanov V. F., Doroshenko V., Werner K., 2019, *MNRAS*, 482, 3622
 Wickramasinghe D. T., Tout C. A., Ferrario L., 2014, *MNRAS*, 437, 675