

## SN 2010MB: DIRECT EVIDENCE FOR A SUPERNOVA INTERACTING WITH A LARGE AMOUNT OF HYDROGEN-FREE CIRCUMSTELLAR MATERIAL

SAGI BEN-AMI<sup>1</sup>, AVISHAY GAL-YAM<sup>1,16</sup>, PAOLO A. MAZZALI<sup>2,3,4</sup>, ORLY GNAT<sup>5</sup>, MARYAM MODJAZ<sup>6</sup>, ITAY RABINAK<sup>1</sup>, MARK SULLIVAN<sup>7</sup>, LARS BILDSTEN<sup>8</sup>, DOVI POZNANSKI<sup>9</sup>, OFER YARON<sup>1</sup>, IAIR ARCAVI<sup>1</sup>, JOSHUA S. BLOOM<sup>10,11</sup>, ASSAF HORESH<sup>12</sup>, MANSI M. KASLIWAL<sup>13</sup>, SHRINIVAS R. KULKARNI<sup>12</sup>, PETER E. NUGENT<sup>10,11</sup>, ERAN O. OFEK<sup>1</sup>, DANIEL PERLEY<sup>12</sup>, ROBERT QUIMBY<sup>14</sup>, AND DONG XU<sup>15</sup>

<sup>1</sup> Department of Particle Physics and Astrophysics, The Weizmann Institute of Science, Rehovot 76100, Israel; [sagi.ben-ami@weizmann.ac.il](mailto:sagi.ben-ami@weizmann.ac.il)

<sup>2</sup> Astrophysics Research Institute, Liverpool John Moores University, Liverpool L3 5RF, UK

<sup>3</sup> Max-Planck-Institut für Astrophysik, Karl-Schwarzschildstr. 1, D-85748 Garching, Germany

<sup>4</sup> INAF-Osservatorio Astronomico, vicolo dell'Osservatorio, 5, I-35122 Padova, Italy

<sup>5</sup> Racah Institute of Physics, The Hebrew University, 91904 Jerusalem, Israel

<sup>6</sup> Center for Cosmology and Particle Physics, Department of Physics, New York University, 4 Washington Place, room 529, New York, NY 10003, USA

<sup>7</sup> Department of Physics (Astrophysics), University of Oxford, DWB, Keble Road, Oxford, OX1 3RH, UK

<sup>8</sup> Kavli Institute for Theoretical Physics and Department of Physics Kohn Hall, University of California, Santa Barbara, CA 93106, USA

<sup>9</sup> School of Physics and Astronomy, Tel-Aviv University, Tel Aviv 69978 Israel

<sup>10</sup> Department of Astronomy, University of California, Berkeley, CA 94720-3411, USA

<sup>11</sup> Physics Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA

<sup>12</sup> Cahill Center for Astrophysics, California Institute of Technology, Pasadena, CA 91125, USA

<sup>13</sup> Observatories of the Carnegie Institution for Science, 813 Santa Barbara Street, Pasadena, CA 91101, USA

<sup>14</sup> Kavli IPMU, University of Tokyo, 5-1-5 Kashiwanoha, Kashiwa-shi, Chiba 277-8583, Japan

<sup>15</sup> Dark Cosmology Centre, Niels Bohr Institute, University of Copenhagen, Juliane Maries Vej 30, DK-2100 Copenhagen, Denmark

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### ABSTRACT

We present our observations of SN 2010mb, a Type Ic supernova (SN) lacking spectroscopic signatures of H and He. SN 2010mb has a slowly declining light curve (LC) ( $\sim 600$  days) that cannot be powered by  $^{56}\text{Ni}/^{56}\text{Co}$  radioactivity, the common energy source for Type Ic SNe. We detect signatures of interaction with hydrogen-free circumstellar material including a blue quasi-continuum and, uniquely, narrow oxygen emission lines that require high densities ( $\sim 10^9 \text{ cm}^{-3}$ ). From the observed spectra and LC, we estimate that the amount of material involved in the interaction was  $\sim 3 M_{\odot}$ . Our observations are in agreement with models of pulsational pair-instability SNe described in the literature.

**Key words:** stars: mass-loss – stars: Population III – supernovae: general

**Online-only material:** color figures, machine-readable table

### 1. INTRODUCTION

A massive star with an initial mass above  $\sim 8 M_{\odot}$  ends its life in an explosion that destroys the star, leaving a neutron star (NS) or a black hole (BH) as a remnant (Heger et al. 2003). This explosion is triggered by iron photo-disintegration and loss of internal energy causing the star to undergo a gravitational core-collapse (CC) supernova (SN; e.g., Woosley & Janka 2005).

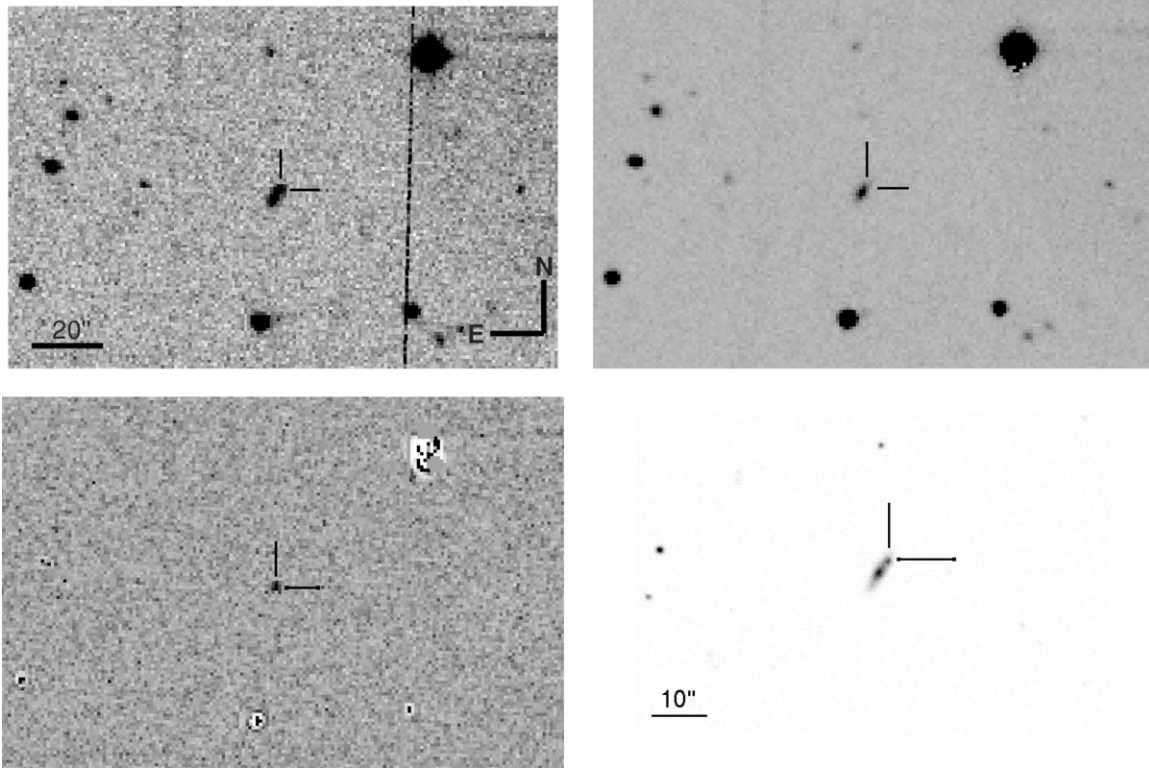
At an earlier stage in its evolution, an extremely massive star (initial mass above  $\sim 100 M_{\odot}$ ) will go through a phase of electron-positron generation in its core (Barkat et al. 1967; Rakavy & Shaviv 1968; Heger et al. 2003; Waldman 2008; Chatzopoulos & Wheeler 2012; Yusof et al. 2013). The pair production will render the star unstable, with two possible outcomes: (1) the star will end its life in an explosion that ejects the entire mass of the star and leaves no remnant at all, a pair-instability supernova (PISN; progenitor mass above  $\sim 140 M_{\odot}$ ). (2) The star will eject matter in a series of eruptions, a pulsation pair-instability (PPI) event (progenitor mass between  $\sim 100$ – $140 M_{\odot}$ ), thus reducing the core mass until it reaches hydrostatic equilibrium and returns to the normal evolution track of massive stars, ending its life due to CC (Heger & Woosley 2002; Heger et al. 2003; Woosley et al. 2007; Waldman 2008;

Chatzopoulos & Wheeler 2012). Several examples of possible PISN events have been observed in recent years (SN 2007bi and PTF10nmn; Gal-Yam et al. 2009; Gal-Yam 2012; O. Yaron et al. 2014, in preparation; though see Nicholl et al. 2013). Only indirect evidence has so far been presented for a PPI event (e.g., SN 2006gy and SN 2006jc; Woosley et al. 2007; Pastorello et al. 2008a; Chugai 2009).

For both CC SNe and PISNe, the energy source at late times is the radioactive decay of  $^{56}\text{Co} \rightarrow ^{56}\text{Fe}$ , with a half-life of 77 days. Other energy sources that can contribute to the optical display of an SN are hydrogen recombination in the expanding ejecta (Type II-P SN; Popov 1993), and an internal engine such as a nascent magnetar (Kasen & Bildsten 2010; Woosley 2010; Mazzali et al. 2009) or accretion by a stellar BH (Patnaude et al. 2011). In some cases, interaction of SN ejecta with circumstellar material (CSM), most commonly observed through narrow hydrogen lines in Type IIn SNe (e.g., Schlegel 1990; Kiewe et al. 2012), will inject additional energy into the optical display.

Here, we describe SN 2010mb (PTF10iue), an SN lacking signatures of either hydrogen or helium (Type Ic), with long-lasting emission powered by interaction of the SN ejecta with a large mass of hydrogen-free CSM. Section 2 describes our observations; and Section 3 presents our results and analysis. In Section 4, we discuss possible scenarios, with an emphasis on the PPI option. Conclusions are given in Section 5.

<sup>16</sup> Kimmel Investigator.



**Figure 1.** Top left: an image of SN 2010mb taken with the P48 on 2010 June 10 UT (MJD 55357.69). Top right: reference image for SN 2010mb. Bottom left: the SN 2010mb image after reference subtraction. Residuals from host galaxy contamination are negligible, as is apparent in the image. Bottom right: a zoomed image of SN 2010mb obtained on 2011 March 4 UT (MJD 55624.11) with LRIS mounted on the Keck-I 10 m telescope. The offset of the SN from the host center ( $1'3$  W and  $2'2$  N) is clearly seen.

## 2. OBSERVATIONS

### 2.1. Discovery

On 2010 April 10 UT (MJD 55296.76) the Palomar Transient Factory (PTF; Law et al. 2009, Rau et al. 2009) detected SN 2010mb at R.A. =  $16^{\text{h}}00^{\text{m}}23^{\text{s}}.103$  and decl. =  $37^{\circ}44'57''$  (detection magnitude  $21.47 \pm 0.2$  in  $r$  band) using the CFH12K survey camera mounted on the 48" Oschin Schmidt telescope at the Palomar observatory (P48). Analysis of previous images showed the SN was visible on 2010 March 18 UT (MJD 55273.96), at a magnitude of  $22.2 \pm 0.54$  in the  $r$  band. The object is located on the edge of the galaxy SDSS J160023.23+374454.8 at a redshift  $z = 0.1325$ . Figure 1 shows detection, reference, and subtracted images of the SN.

Prior to 2010 March 18 UT, the galaxy was imaged 10 times by the PTF survey between 2009 May 18 and August 21 UT (MJD 54969–55064), with no evidence for the SN or any activity in its vicinity.

We classified SN 2010mb as a Type Ic SN based on a spectrum lacking signatures of either hydrogen or helium taken on 2010 June 8 UT (MJD 55355. Figures 2 and 3; Ben-Ami et al. 2012).

### 2.2. Photometry

SN 2010mb, discovered near the P48 detection limit, was intermittently detected during the first 50 days after its discovery,<sup>17</sup> until it became continuously visible on 2010 May 31 UT

(MJD 55347.71) at an apparent magnitude of  $20.85 \pm 0.12$  in the  $r$  band.

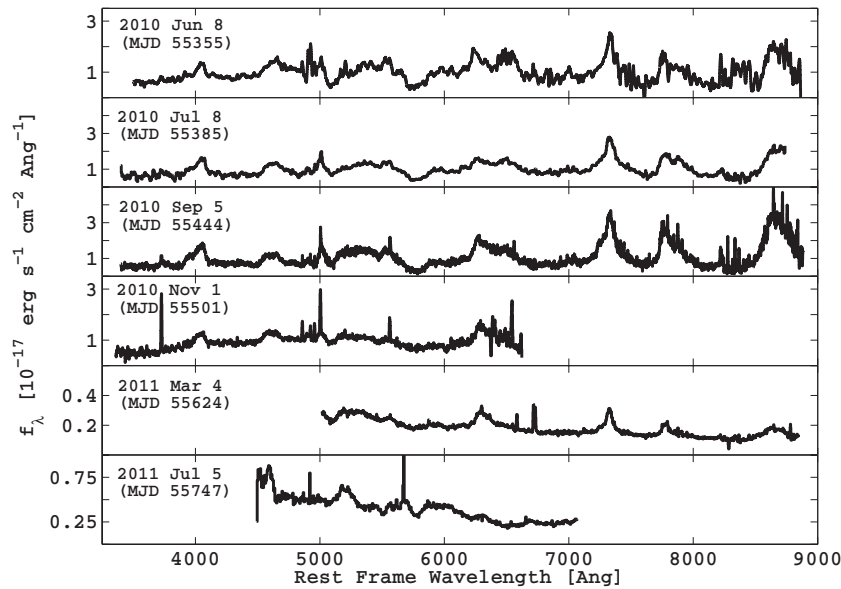
Photometry of SN 2010mb was obtained by the P48 (Law et al. 2009; Rau et al. 2009), the gamma-ray burst camera (Cenko et al. 2006) mounted on the Palomar 60" telescope (P60), the Large Format Camera mounted on the Palomar 200" Hale telescope (P200), and the Low Resolution Imaging Spectrograph (LRIS) mounted on the 10 m Keck-I telescope (Oke et al. 1995). Data were reduced using the MKDIFFLC photometry routine (Gal-Yam et al. 2004, 2008), except for P48 data reduced using point-spread function (PSF) photometry on image subtractions.<sup>18</sup> The PTF data processing is discussed in Laher (2014<sup>19</sup>), while the photometric calibration is discussed in Ofek et al. (2012a; 2012b). We adopt a distance modulus of 39.07 mag, based on measured redshift from the SN spectrum (assuming  $H_0 = 67.3 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $\Omega_m = 0.315$ ; Planck Collaboration et al. 2013), corresponding to a luminosity distance of  $\approx 650.46 \text{ Mpc}$ , a Galactic extinction correction of  $E(B - V) = 0.013 \text{ mag}$  ( $A_r = 0.033 \text{ mag}$ ; Schlafly & Finkbeiner 2011; see NED,<sup>20</sup> as well as Section 3.6.) Photometric results are given in Table 1 and plotted in Figure 4. No correction for the host extinction is applied. The lack of strong interstellar medium (ISM) absorption and the blue

<sup>17</sup> We define a detection as a  $5\sigma$  signal above the zero point photon count. In cases where the object did not pass this criterion, 2010 April 30 and May 12 UT (MJD 55316 and MJD 55328, respectively), the signals were  $4\sigma$  and  $2.6\sigma$  above the zero point photon count, respectively.

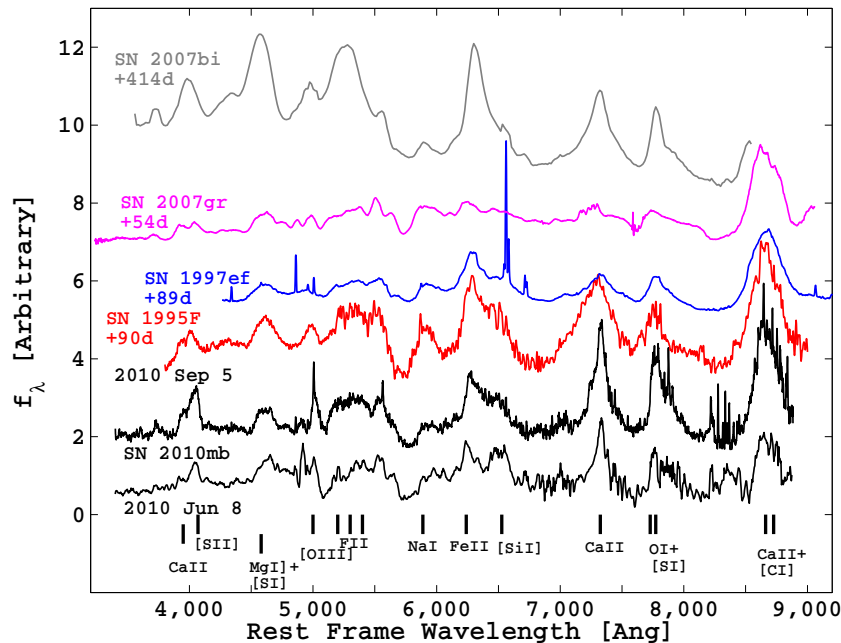
<sup>18</sup> Image subtraction using previous images of the same field of view taken with the same instrument during 2009 as references was carried out, followed by forced PSF photometry and absolute calibration to the SDSS catalog as done, e.g., in Ofek et al. (2013).

<sup>19</sup> See also Laher et al. 2012 for details on the Aperture Photometry Tool.

<sup>20</sup> NASA/IPAC Extragalactic Database (NED) is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.



**Figure 2.** SN 2010mb spectra. No signs of hydrogen lines (emission or absorption) or helium lines are seen, indicating a Type Ic SN. Late-time spectra (2011 March 4 and July 5 UT) are dominated by a blue quasi-continuum component. Spectra are available in digital form from the Weizmann Interactive Supernova Data Repository (WISereP; Yaron & Gal-Yam 2012; <http://www.weizmann.ac.il/astrophysics/wiserep/>).



**Figure 3.** SN 2010mb spectral classification. Early spectra of SN 2010mb resemble a Type Ic SN at the transition between the photospheric and nebular phases as late as  $\sim 170$  days after discovery (i.e., some photospheric features were still observed in the spectrum taken on 2010 September 5 UT, second spectrum from the bottom). Automatic classification of the spectra taken on 2010 June 8 and July 8 UT using Superfit (Howell et al. 2005) suggests that the best match is to the late-time spectra of SN 1997ef 89 days after peak magnitude (an energetic Type Ic SN showing late transition from photospheric phase to nebular phase; Mazzali et al. 2004), and SN 1995F 90 days after discovery (a Type Ic SN; Matheson et al. 2001), while automatic classification using SNID (Blondin & Tonry 2011) suggests that the best match is to SN 2007gr at 54 days after peak magnitude (a “normal” Type Ic; Hunter et al. 2009). This highlights the slow evolution of SN 2010mb, similar to SN 1997ef, and much slower than a “normal” Type Ic SN like SN 2007gr. The spectra also resemble nebular spectra of SN 2007bi (Gal-Yam et al. 2009).

(A color version of this figure is available in the online journal.)

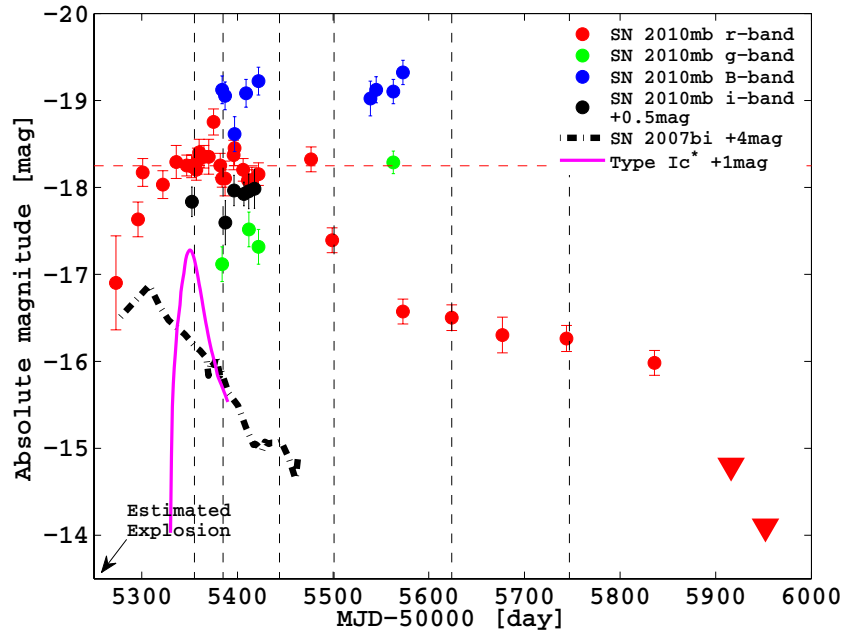
spectral energy distribution (SED) measured at late times (see Section 3) argue against a substantial host extinction.

The light curve (LC) displayed a long plateau in  $r$  band lasting  $\sim 180$  days, followed by a slow decline at a rate of  $0.004 \text{ mag day}^{-1}$ , while the  $g$  band and  $B$  band LCs are slowly rising during the first  $\sim 250$  days (Figure 4). The total amount of energy emitted in the  $r$  band over a period of  $\sim 600$  days is  $\sim 1.2 \times 10^{50} \text{ erg}$  (an average flux of  $1.8 \times 10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2}$ ). Assuming a bolometric correction of 10% at early times and

50% at later times based on our spectral decomposition (see Section 3.2), we estimate a total energy emission of  $3.7 \pm 0.4 \times 10^{50} \text{ erg}$  over a period of  $\sim 600$  days.

### 2.3. Spectroscopy

Six spectra of SN 2010mb were obtained (Table 2 and Figure 2). Data were reduced using standard IRAF and IDL routines (e.g., Matheson et al. 2000b, Gal-Yam et al. 2007)



**Figure 4.** Photometry of SN 2010mb. The *r* band LC rises by 1.3 mag during the first  $\sim 25$  days, and then settles onto a plateau lasting  $\sim 180$  days at an absolute magnitude of  $-18.25$  mag (deviations from the plateau, marked by a dashed red line, are not statistically significant). Following the plateau, the LC drops by 0.9 mag in less than 23 days, followed by a period of slow decline at a rate of  $0.004 \text{ mag day}^{-1}$ , which is less than half the decline rate of the  $^{56}\text{Ni}/^{56}\text{Co}$  powered SNe, during the next 350 days. Photometry in the *g* band and *B* band showed a slow increase in flux during the first 250 days of this event. The dashed vertical lines indicate a spectrum was taken on that day. The LCs of SLSN-R SN 2007bi in the *R* band (candidate pair instability event; Gal-Yam et al. 2009), and of an average Type Ic SN template in the *R* band (Drout et al. 2011), are shown for comparison. On 2011 December 21 UT (MJD 55916), the SN was no longer visible in deep images taken with the LFC. Deeper images taken with LRIS on 2012 January 26 UT (MJD 55952.09) set a limiting absolute magnitude of  $-14.1$  at that time (triangles).

(A color version of this figure is available in the online journal.)

**Table 1**  
SN 2010mb Photometry

Date (MJD)	Instrument	Exposure Time (s)	Apparent Magnitude (mag)	Magnitude Error (mag)	Filter
55273.96	P48 <sup>a</sup>	$2 \times 200$	22.20	0.54	<i>r</i>
55296.76	P48	$2 \times 200$	21.47	0.20	<i>r</i>
55384.84	P60 <sup>b</sup>	360	21.00	0.10	<i>r</i>
55477.64	LFC <sup>c</sup>	900	20.78	0.15	<i>r</i>
55624.11	LRIS <sup>d</sup>	120	22.60	0.14	<i>r</i>

**Notes.** Digital data are available from the Weizmann Interactive Supernova Data Repository (WiSeREP; Yaron & Gal-Yam 2012; <http://www.weizmann.ac.il/astrophysics/wiserep/>).

<sup>a</sup> Palomar 48" Oschin Schmidt telescope.

<sup>b</sup> Palomar 60" Oscar Mayer telescope.

<sup>c</sup> Large Format Camera mounted on the Hale 200" telescope.

<sup>d</sup> Low Resolution Imaging Spectrograph mounted on the Keck-I 10 m telescope.

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

and were smoothed with a three-pixel boxcar. The observed spectra are deredshifted by  $z = 0.1325$ , where the redshift was determined using a  $\chi^2$  minimization with an absorption-line template (Howell et al. 2005; Galactic lines identification in later spectra, e.g., 2012 February 20 spectrum, confirms our redshift determination). Spectra are calibrated to the *gri* P48/P60 photometry (2010 June 8, July 8, and November 1 UT), to an extrapolation of the *gri* P48/P60 photometry (2010 September 5 UT), and to LFC *r* band photometry (2011 March 4 and July 5 UT). A spectrum of the host galaxy nucleus was taken on 2012 February 20 UT using LRIS (see Section 3.6).

### 3. ANALYSIS

#### 3.1. Classification and Evolution

The first three spectra of SN 2010mb (2010 June 8, July 8, and September 5 UT) resemble a Type Ic SN at the transition between photospheric and nebular phase as late as  $\sim 170$  days after discovery (i.e., photospheric features observed in the spectrum taken on 2010 September 5 UT, Figure 2). Automatic classification using Superfit (Howell et al. 2005) of the spectra taken on 2010 June 8 and July 8 UT suggests that the best match is to the late-time spectra of SN 1997ef 89 days after peak magnitude (an energetic Type Ic SN showing late transition



**Table 2**  
Spectroscopy Log

Date (MJD)	Telescope	Instrument	Exposure Time (sec)	Grism/Grating (lpm)	Slit
2010 Jun 8 (55355) <sup>a</sup>	Keck-I	LRIS	600	400/400 <sup>b</sup>	1''
2010 Jul 8 (55385)	Keck-I	LRIS	750	400/400	1''
2010 Sep 5 (55444)	Keck-I	LRIS	750	400/400	1''
2010 Nov 1 (55501)	Keck-I	LRIS	750	400/400	1''
2011 Mar 4 (55624)	Keck-I	LRIS	900	400/400	1''
2011 Jul 5 (55747)	Keck-II	DEIMOS <sup>c</sup>	1200	600	0.7''
2012 Feb 20 (55977) <sup>d</sup>	Keck-I	LRIS	1200	400/400	1''
Date (MJD)	Resolution (Å)	Airmass	Paralactic Angle	Position Angle	Spec. Standard
2010 Jun 8 (55355)	1.09/1.16	1.38	−75°	27°	BD+28 4211
2010 Jul 8 (55385)	1.09/1.16	1.01	134°	270°	BD+26 2606
2010 Sep 5 (55444)	1.09/1.16	1.24	−90°	19°	BD+28 4211
2010 Nov 1 (55501)	1.09/1.16	1.55	−90°	−8°	BD+28 4211
2011 Mar 4 (55624)	1.09/1.16	1.29	14°	270°	BD+28 4211
2011 Jul 5 (55747)	0.65	1.07	42°	256°	BD+28 4211
2012 Feb 20 (55977)	1.09/1.16	1.18	238°	328°	BD+26 2606

**Notes.**

<sup>a</sup> The CBET classifying this object (Ben-Ami et al. 2012) erroneously lists the Palomar 200'' Hale telescope as that used for the spectrum taken on 2010 June 8 UT. The details here supersede this publication.

<sup>b</sup> The Low Resolution Imaging Spectrograph has two arms (blue and red). Each arm is equipped with a different grating and, therefore, a different resolution.

<sup>c</sup> Deep Imaging Multi-Object Spectrograph Faber et al. (2003); filter: GG455.

<sup>d</sup> Host Galaxy spectrum.

from photospheric phase to nebular phase; Mazzali et al. 2004), and SN 1995F 90 days after discovery (a Type Ic SN; Matheson et al. 2001), while automatic classification using SNID (Blondin & Tonry 2011) suggests that the best match is to SN 2007gr at 54 days after peak magnitude (a “normal” Type Ic; Hunter et al. 2009), Figure 3. This highlights the slow evolution of SN 2010mb, similar to SN 1997ef, and much slower than a “normal” Type Ic SN like SN 2007gr. The spectra also resemble nebular spectra of SN 2007bi (PISN; Gal-Yam et al. 2009). Based on the resemblance between SN 2010mb spectrum taken on 2010 June 8 UT, and SN 1997ef spectrum taken 89 days after peak magnitude, and assuming a rise time of 20 days (Mazzali et al. 2004), we assume SN 2010mb exploded around 2010 February 23 UT, though this number is highly uncertain.

Early and intermediate spectra (i.e., 2010 June 8 until November 1 UT) are dominated by blended lines of Ca II  $\lambda\lambda 3933, 3968$  and [S II]  $\lambda 4069$ , Mg I  $\lambda 4571$  and [S I]  $\lambda 4589$ , [O III]  $\lambda\lambda 4959, 5007$ , Fe II blended lines (5200–5400 Å), [O I] and Fe II blended lines (6300–6400 Å), O I  $\lambda 7773$  and [S I]  $\lambda 7728$ , and Ca II  $\lambda\lambda 8498, 8542, 8662$ , and [C I]  $\lambda 8727$ . The Na I D line, [Si I]  $\lambda 6527$ , and [Ca II]  $\lambda\lambda 7291, 7324$  are clearly observed in the spectra as well (line identification is based on the models described in Section 3.2; Wavelengths are given in rest frame); see Figure 3. Some of the lines, such as the Mg I  $\lambda 4571$  line, seem to have an internal structure, while others, such as the [Ca II]  $\lambda\lambda 7291, 7324$  seem to have narrow cores, an indication for a different spatial distribution of these elements and, therefore, for a non-spherical geometry; see Figure 5. Late-time spectra (2011 March 4 and July 5 UT; Figure 2) are dominated by a blue quasi-continuum component, also reflected in *g* band and *B* band photometry that is slowly rising during the first  $\sim 250$  days, a behavior inconsistent with that of a purely hydrodynamic radioactive explosion that would show a monotonic decrease in temperature with time. We therefore conclude that

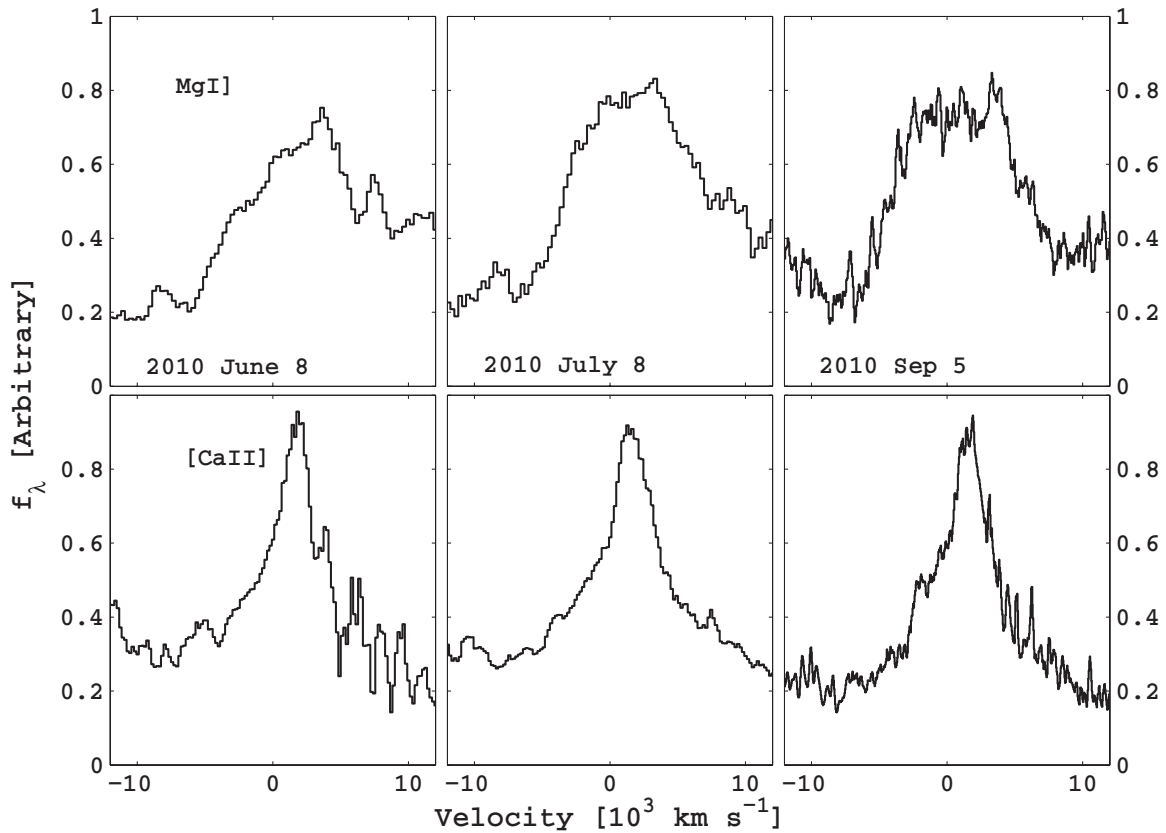
another energy source other than radioactive  $^{56}\text{Ni}$  decay, the common energy source powering Type Ic SNe, is driving the optical display of SN 2010mb at late times.

### 3.2. Modeling

We model the spectra obtained on 2010 June 8, July 8, September 5, and 2011 March 4 UT. We start by modeling the blue quasi-continuum,<sup>21</sup> which seems to be the strongest component in late-time spectra (2011 March 4 and July 5 UT). An attempt to fit blackbody profiles at temperatures between 5000–10,000 K to the observed blue quasi-continuum gave unsatisfactory results. We conclude that the quasi-continuum is nonthermal in the sense that it is not described well by a blackbody of any temperature, a behavior also known for Type II<sub>n</sub> SNe (e.g., Kiewe et al. 2012). We therefore use the spectral continuum derived from observations of SN 2005cl, a hydrogen-rich Type II<sub>n</sub> SN, with prominent blue quasi-continuum associated with CSM–ejecta interaction. We use the spectrum taken on 2005 August 13 UT, in which the CSM–ejecta interaction seems to be the strongest. The spectrum is dominated by overlapping Fe II lines, some of which are partially resolved (Kiewe et al. 2012). After removing the Balmer lines and fitting a smooth curve to the continuum, we scale the resulting curve to fit the continuum seen in each of the SN 2010mb observed spectra. The match between the SN 2010mb blue quasi-continuum at late times when the blue quasi-continuum is more prominent, and the model based on the SN 2005cl spectrum, lends credence to this approach.

After removing the blue quasi-continuum from the four spectra, we compose a photospheric model to match the observed spectra. SN 1997ef and SN 2010mb both show an extended

<sup>21</sup> As the spectra are calibrated to host-subtracted photometry, this component is unlikely to be residual host contamination.



**Figure 5.** SN 2010mb line profiles. Comparing the nebular emission lines from Mg I (top panels) and [Ca II] (bottom panels), we see a flat-topped profile in the former line, usually interpreted as a signature of non-spherical geometry (Mazzali et al. 2007; Maeda et al. 2008; Mazzali et al. 2005; Modjaz et al. 2008; Taubenberger et al. 2009). The different line shapes of [Ca II] and Mg I would, in this case, indicate a non-uniform spatial abundance structure.

LC and a late transition from photospheric to nebular emission (Figures 2, 3, and 6). We therefore use photospheric models based on SN 1997ef (Iwamoto et al. 2000; Mazzali et al. 2000) for SN 2010mb. After removing the photospheric component and following the methodology of Mazzali et al. (2004), we are left with nebular spectra.

To model the nebular component of SN 2010mb, we construct custom models, using nebular models of SN 1997ef (Mazzali et al. 2000, 2004; Mazzali & Lucy 1993; Lucy 1999; Mazzali et al. 2007) as a starting point with adjustments to each element mass, so that the nebular line strengths agree with the observations. We assume that the emitting nebula is spherical and has a sharp outer boundary, defined by the width of the emission lines. We adopt a boundary velocity of  $5000 \text{ km s}^{-1}$ , in agreement with the nebular O I lines, the Na I D line, the Ca II/[C I] blend near  $8600 \text{ \AA}$ , and the Fe II/Na I D emission near  $5200 \text{ \AA}$ , which is a blend of many different lines of different strength. The flat topped lines observed in some of the cases (e.g., Mg I) shown in Figure 5) indicate additional complexity, possibly a nonuniform spatial abundance structure. Based on the evolution of the LC and the spectra, e.g., the large flux fraction at late times from the blue quasi-continuum, we introduce an external energy source to the nebular model that increases with time, which we associate with high energy photons coming from interaction of the SN ejecta with hydrogen-free CSM.

The good match between the models and the observed spectra confirms our assumption regarding the three components that comprise SN 2010mb observed spectra, as well as the use of SN 2005cl spectra and SN 1997ef models as starting points. The decomposition results are shown in Figure 6, while Table 3

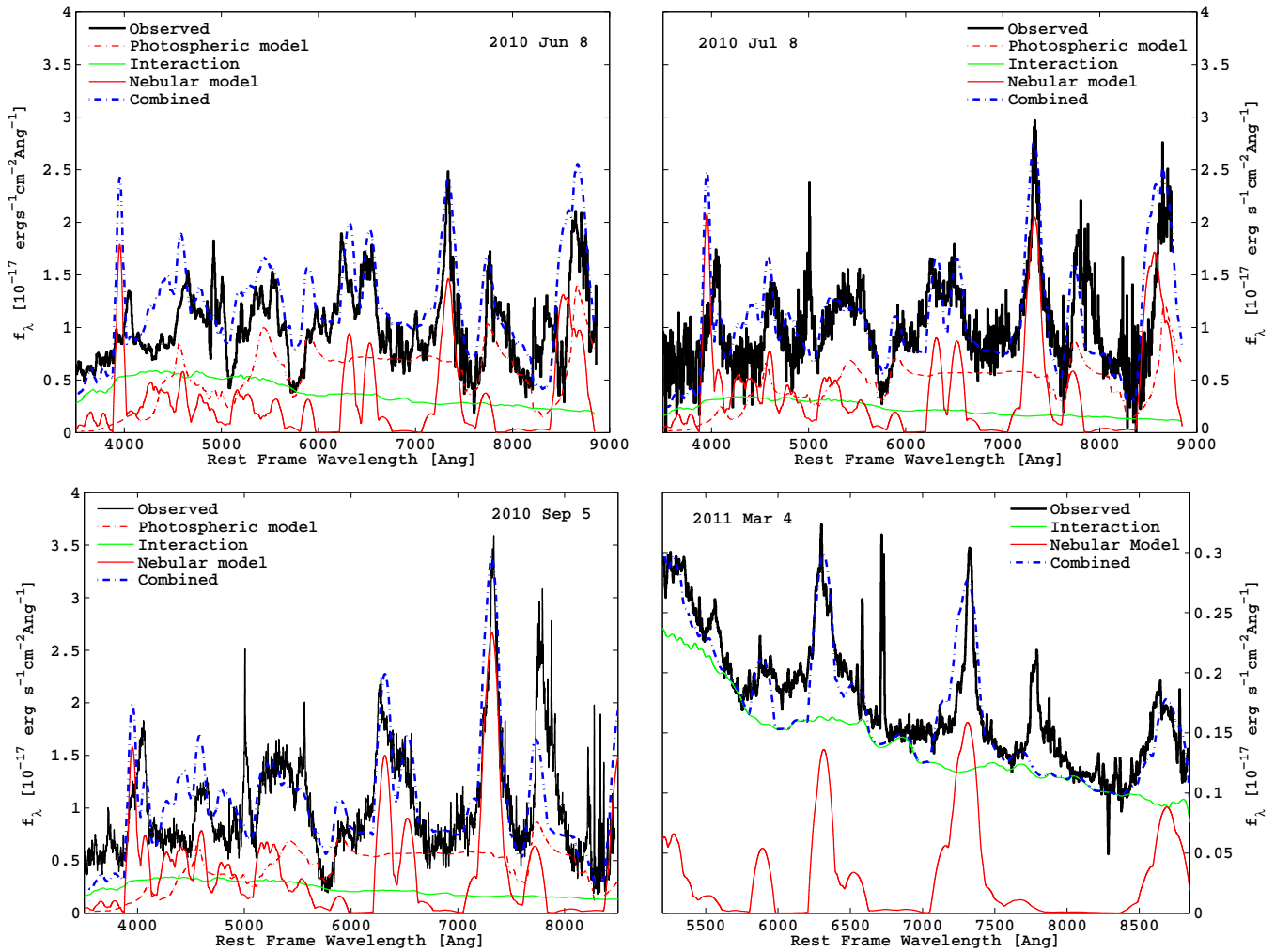
**Table 3**  
Nebular Model Composition

Element	2010 Jun 8	2010 Jul 8	2010 Sep 5	2011 Mar 4
Si	7	6.5	10	8,
Ca	0.27	0.19	0.67	0.1
O	3	2.4	5	1.6
C	0.8	0.7	3	1.2
Na	0.0001	0.00015	0.00015	0.0013
Mg	0.0006	0.0008	0.0013	0.02
S	3	2.5	5	1
Total	14.0	12.4	23.7	11.0

**Note.** Masses are given in solar mass ( $M_{\odot}$ ).

gives our results for the chemical composition and emitting mass from the nebular phase at each epoch. The mass of each element is scaled linearly with the line strengths, and from the match quality, we derive a  $\leq 50\%$  uncertainty in the derived masses for most of the elements.

The derived masses indicate we have observed the explosion of a massive stripped-envelope star (for comparison, the ejecta mass in an average Type Ic SN is  $1.7 M_{\odot}$ ; Drout et al. 2011). This is further supported by the long transition phase from photospheric to nebular emission; delayed disappearance of photospheric (optically thick) emission indicates a large total mass (Mazzali et al. 2004). The simultaneous detection of photospheric/nebular SN emission and radiation from the ejecta-CSM interaction suggests either that the outer hydrogen-free material is clumpy, or a non-spherical geometry of the CSM



**Figure 6.** SN 2010mb spectral decomposition. We decompose each spectrum to three components: photospheric and nebular models based on SN 1997ef (Mazzali & Lucy 1993; Lucy 1999; Mazzali et al. 2000, 2004, 2007); and an interaction component based on comparison to the continuum shape (H lines excluded) of the spectrum of SN 2005cl, a hydrogen-rich Type II SN, with a prominent blue quasi-continuum originating from the SN ejecta interacting with hydrogen-rich CSM (Kiewe et al. 2012). Some of the lines, such as the [O III]  $\lambda 5007$  and the O I  $\lambda 7773$  recombination lines, are not handled by our modeling code.

(A color version of this figure is available in the online journal.)

(e.g., in a disk or torus). Such a geometry allows us to observe the ejecta directly, in addition to ongoing interaction in the inner circumference of the CSM structure. The mass estimates based on the nebular model should be increased, correcting for the covering fraction of the CSM.

### 3.3. Light Curve Decomposition

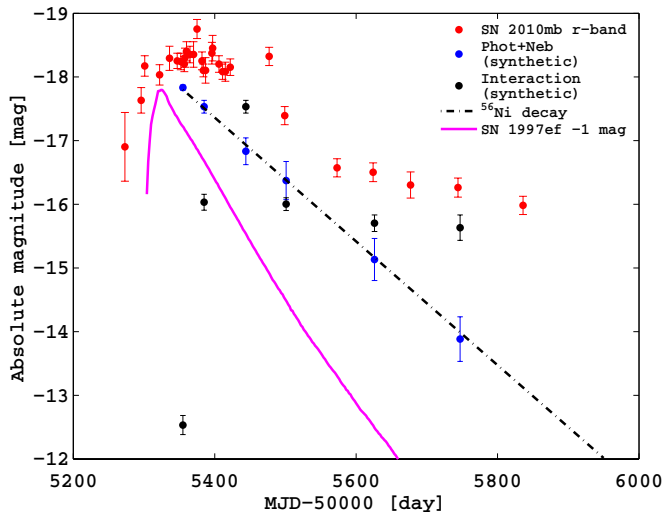
Based on our spectral decomposition, we estimate for each point in the LC the fraction of SN flux (photospheric+nebular), as well as that coming from the ejecta–CSM interaction. We thus produce two synthetic LCs—one representing the SN and the second representing the ejecta–CSM interaction, plotted in Figure 7. The synthetic SN LC decline is in agreement with that expected from  $^{56}\text{Ni}$  decay, assuming full trapping of the gamma-rays as was observed for SN 1987A, indicating a large ejecta mass (the estimated ejecta mass for SN 1987A is  $15 M_{\odot}$ ; e.g., de Freitas Pacheco 1989). The amount of radioactive nickel ( $0.1 \pm 0.01 M_{\odot}$ ) suggested by the model is also consistent with the nickel mass–peak magnitude relation for SNe (Perets et al. 2010), predicting  $\sim 0.1\text{--}0.15 M_{\odot}$  of  $^{56}\text{Ni}$ , using our photometry. The synthetic LC for the interaction component shows a maximum at around 2010 September 1 UT (MJD 55440) and

a much slower decline; and it is dominating the optical display from around 2010 November 20 UT (MJD 55520).

### 3.4. Interaction Lines

Some of the spectra show a prominent emission line on top of the [O I]  $\lambda 5577$  nebular feature, Figure 8. We detected and measured this emission line flux by fitting a function to the line profile using the standard IRAF procedure *splot*, as well as a custom MATLAB script, which removes nearby continuum by using a spline function at an area of  $\pm 100 \text{ \AA}$  around the line peak. We verified that the line is not sensitive to varying the parameters in the nebular modeling code (i.e., changing the amount of radiating mass in the nebular phase) and the blue quasi-continuum flux, and we found a variability of  $\sim 5\%$  for the line flux. We ran Cloudy, a spectral synthesis code designed to predict the physical properties in the ISM under a broad range of conditions<sup>22</sup> (Ferland et al. 1998). We modeled slabs of optically thin and collisional gas, in order to compute the

<sup>22</sup> Cloudy is designed to compute the ionization and thermal properties, as well as the emitted spectrum, of plasma in collisional or photoionization equilibrium.



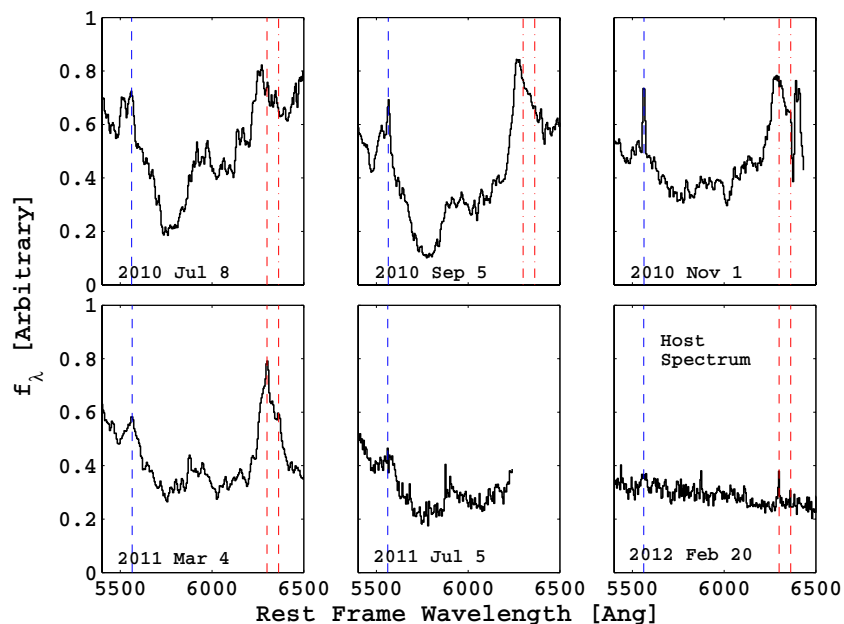
**Figure 7.** SN 2010mb synthetic LC. A synthetic LC (blue), showing the decay of the photospheric and nebular components only (based on the spectral decomposition, and interpolated between spectral epochs), behaves as expected of a  $^{56}\text{Ni}$  powered SN (dashed black curve), assuming full trapping of the gamma-rays, indicating a large ejecta mass. A synthetic LC of the interaction component is given as well (black dots), as is a scaled LC of SN 1997ef, a slowly declining Type Ic SN (magenta; Mazzali et al. 2004). The results suggest that a large fraction of the observed flux is not internal to the SN explosion but originates from interaction of the SN ejecta with CSM.

(A color version of this figure is available in the online journal.)

emission-line intensities associated with specific elements, as functions of the gas temperature and density. Two models are presented in Figure 9, both with parameters chosen so that the flux of the [O I]  $\lambda 5577$  line is similar to/greater than that of

the [O I]  $\lambda\lambda 6300, 6363$  lines, as observed: (1) a composition of S, C, O, Si, Ca, and Fe, with ratios in accordance with the 2011 March 4 UT nebular model results (Table 3) in order to check whether the [O I]  $\lambda 5577$  line is coming from the SN ejecta (i.e., a slow component in the nebular matter); and (2) a composition of 90% He, 5% C, 5% O, and traces of iron—the expected shell composition from a PPI event (Chatzopoulos & Wheeler 2012). Comparing to the observed spectra, we rule out the nebular model, which would lead to strong lines of Ca II that are not seen. The presence of only [O I] lines is consistent with the PPI model. The property of the [O I]  $\lambda 5577$  line as a tracer of high densities of oxygen is shown in Figure 9, bottom panel. Only at densities  $\sim 10^7 \text{ cm}^{-3}$  does the [O I]  $\lambda 5577$  line have a flux similar to the flux in the [O I]  $\lambda\lambda 6300, 6363$  lines. For the [O I]  $\lambda 5577$  to be the dominant line among the [O I] lines as observed, a density of  $\sim 10^9 \text{ cm}^{-3}$  is required. Further examination of the data show that the [O I]  $\lambda 5577$  line is blueshifted by  $800 \text{ km s}^{-1}$  and has a velocity dispersion of  $600 \text{ km s}^{-1}$ , while the [O I]  $\lambda\lambda 6300, 6363$  lines have the same redshift as the host galaxy, and a velocity dispersion of  $\sim 300 \text{ km s}^{-1}$ , i.e., they are instrumentally unresolved. We therefore identify the [O I]  $\lambda 5577$  high density line with the CSM interaction and the [O I]  $\lambda\lambda 6300, 6363$  lines as the host galaxy emission.

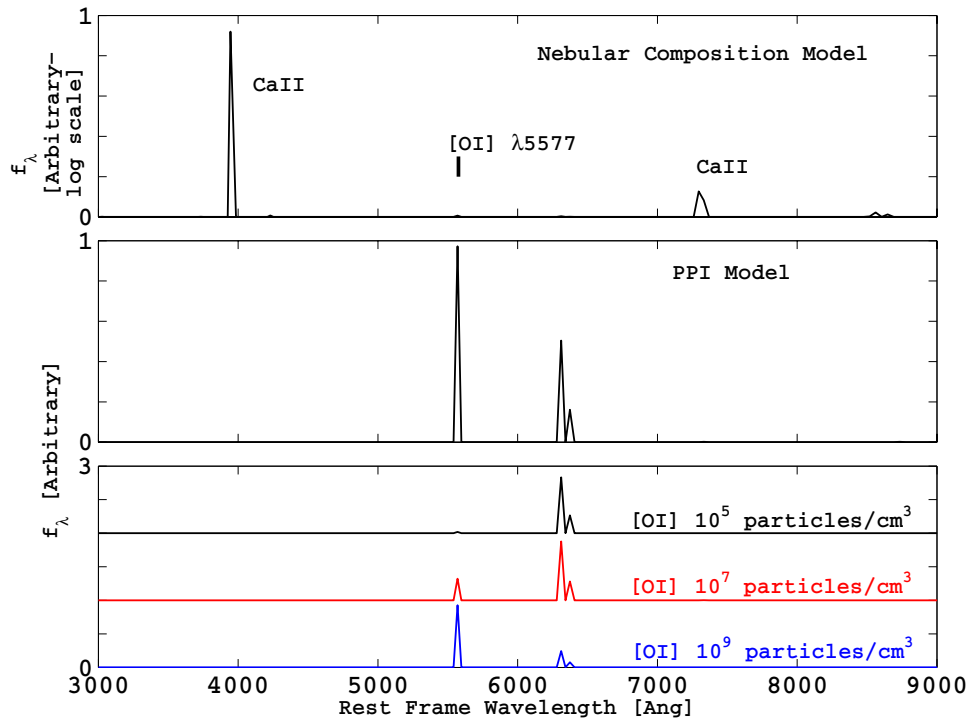
We find that the line at [O I]  $\lambda 5577$  reaches peak intensity between 2010 September 5 and November 1 UT (integrated intensity of  $7 \pm 0.4, 7.8 \pm 0.3 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$ , respectively), while the [O I]  $\lambda\lambda 6300, 6363$  lines are clearly seen in the host galaxy spectrum taken on 2012 February 20 UT. The observed evolution of the [O I]  $\lambda 5577$  line seems to be responding to the blue quasi-continuum. This behavior, as well as the resemblance of the blue quasi-continuum model to the spectral



**Figure 8.** Evolution of the [O I]  $\lambda 5577$  collisionally excited line, which requires high densities ( $\sim 10^9 \text{ cm}^{-3}$  for the [O I]  $\lambda 5577$  to be the dominant among the [O I] lines), similar to those expected in a shell ejected in a PPI event. This line is blueshifted by  $\sim 800 \text{ km s}^{-1}$  and has a velocity dispersion of  $\sim 600 \text{ km s}^{-1}$  (the dashed blue line indicates [O I]  $\lambda 5577$  blueshifted by  $800 \text{ km s}^{-1}$ ). It is getting stronger with time between 2010 July 8 and November 1 UT (integrated flux of  $1.6 \pm 0.3$  and  $7.8 \pm 0.3 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$ , respectively) and traces the flux increase seen in the interaction LC (Figure 7). In late spectra (2011 March 4 and July 5 UT), the peak integrated flux is less than  $1 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$ . While a PPI ejected shell is expected to be rich in He, C, and O (Chatzopoulos & Wheeler 2012), the [O I]  $\lambda 5577$  line is the only signature expected to fall in the optical waveband, as helium ionization requires much higher temperatures than those expected in the ejected shell, and C lines fall mainly in the near IR. The [O I]  $\lambda\lambda 6300, 6363$  lines have the same redshift as the host galaxy and a velocity dispersion of  $\sim 300 \text{ km s}^{-1}$ . They are clearly seen in the host galaxy spectrum taken on 2012 February 20 UT, Figure 11. We therefore identify them as host galaxy emission lines (dashed red).

(A color version of this figure is available in the online journal.)





**Figure 9.** SN 2010mb interaction lines. Two models are presented, both with parameters chosen so that the flux of the [O I]  $\lambda 5577$  line is similar to/greater than that of [O I]  $\lambda\lambda 6300, 6363$  as observed: (1) A composition of S, C, O, Si, Ca, and Fe, with ratios in accordance with the 2010 March 4 UT nebular model results (Table 3)—top panel. Ca II lines are dominant (more than two orders of magnitude brighter than the strongest [O I] line), but are not observed. We find this model is ruled out. (2) A composition of 90% He, 5% C, 5% O, and traces of iron—the expected shell composition from a PPI event (Chatzopoulos & Wheeler 2012)—middle panel. This model is in agreement with the observed lines. The property of the [O I]  $\lambda 5577$  line as a tracer of high densities of oxygen is shown in the bottom panel. Only at densities  $\sim 10^7 \text{ cm}^{-3}$  (red curve), do the [O I]  $\lambda 5577$  and the [O I]  $\lambda\lambda 6300, 6363$  lines have similar fluxes. For the [O I]  $\lambda 5577$  to be dominant among the [O I] lines, a density of  $\sim 10^9 \text{ cm}^{-3}$  is required (blue curve).

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continuum of SN 2010mb (e.g., the 2011 March 4 UT spectrum, Figure 6), shows that at late times the observed LC is powered by interaction of the SN ejecta with hydrogen-free CSM.

### 3.5. Circumstellar Material Mass

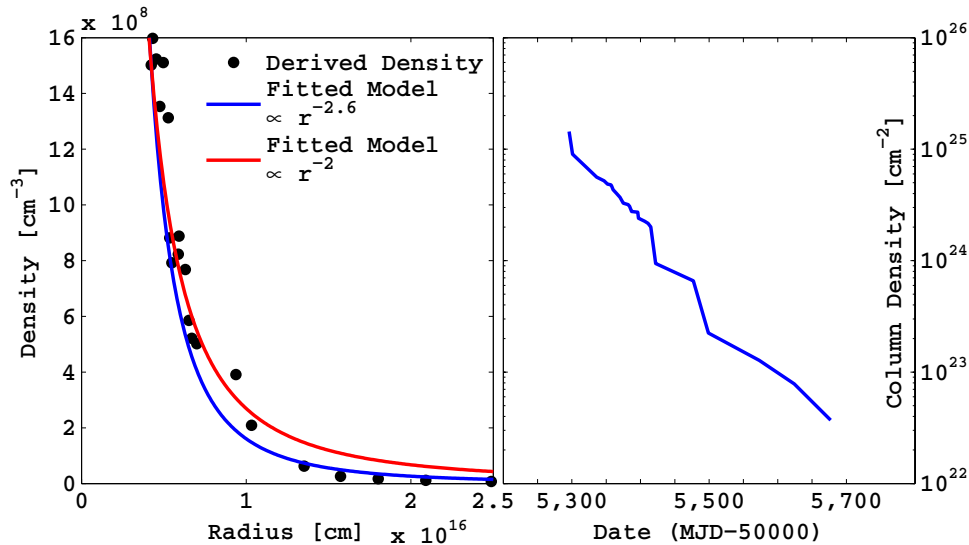
We can estimate the amount of CSM needed to convert the SN ejecta kinetic energy to radiation through interaction. From the LC, we estimate the amount of energy radiated between consecutive data points  $\Delta E$ , and from the line width in the observed spectra, we get the ejecta velocity,  $v \approx 5000 \text{ km s}^{-1}$ . Next, using the relation  $(1/2)\Delta M v^2 = \epsilon \Delta E$ , where  $\epsilon$  is the conversion efficiency and is determined in an iterative way using the relation  $\epsilon = \Delta M_{\text{CSM}}/M_{\text{ej}} + \Delta M_{\text{CSM}}$  (Murase et al. 2011), we derive the mass needed for energy conversion  $\Delta M$ . We estimate the radius at which interaction takes place at any given time by assuming an expansion of the SN ejecta at a constant velocity of  $5000 \text{ km s}^{-1}$ . We then derive a lower limit on the density by dividing the CSM mass by the shell volume swept by the ejecta between two consecutive points in the LC. We get a total CSM mass of  $\sim 3.3 M_{\odot}$  distributed around the progenitor with a density profile  $\propto r^{-2.6}$ , (Figure 10, left panel). The derived densities ( $\sim 10^9 \text{ cm}^{-3}$  peak density) are in agreement with the observed evolution of the [O I]  $\lambda 5577$  recombination line.

Based on this analysis, we conclude the following: (1)  $\sim 3 M_{\odot}$  of hydrogen-free CSM is present at the vicinity of the progenitor, with a mass loss rate higher than those associated with standard winds from Wolf-Rayet (WR) stars (i.e.,  $\sim 10^{-5} M_{\odot} \text{ year}^{-1}$  compared to  $\sim 10^{-5} M_{\odot} \text{ year}^{-1}$ ). (2) Assuming bremsstrahlung is the main cooling mechanism of this plasma, the cooling time is

orders of magnitude shorter than the time gap between adjacent points in the LC, i.e.,  $\sim 10^3 \text{ s}$  at the derived densities versus  $\sim 10^6 \text{ s}$  between adjacent measurements on the LC, and so the CSM–ejecta interaction energy is radiated promptly. (3) SN 2010mb was imaged with the *Swift X-ray telescope* (Burrows et al. 2005) on 2011 April 22 UT. Using an exposure of 5435 s, we derive a  $3\sigma$  upper limit of  $5 \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}$  for the 0.3–10 keV band assuming a photon index of 2. The column density is high enough to explain why X-ray photons have not been detected (Figure 10, right panel).

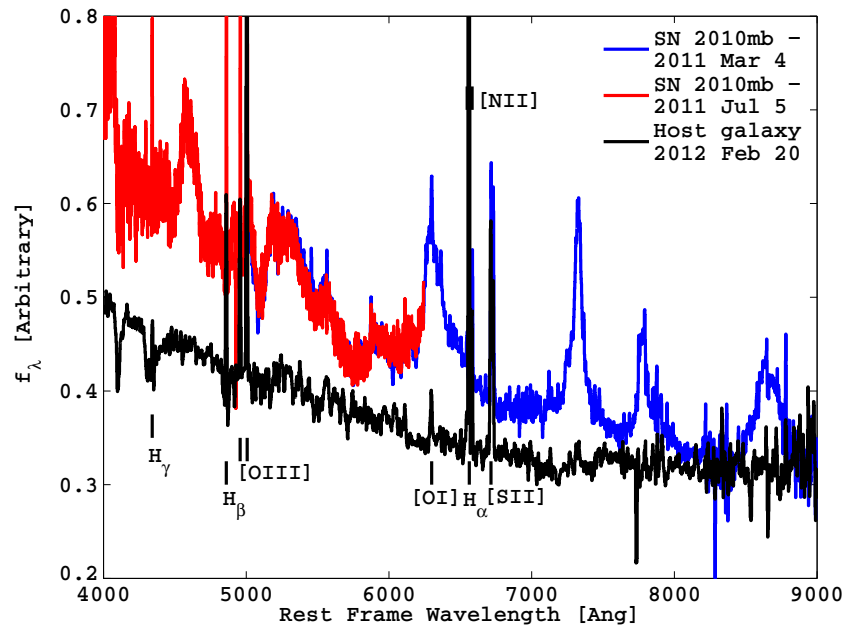
### 3.6. Host Galaxy

The host galaxy spectrum taken on 2012 February 20 UT (taken at the galaxy nucleus  $1'3 \text{ E}$  and  $2'2 \text{ S}$  from the SN location; Figure 11) shows narrow Balmer emission lines filling slightly wider Balmer absorption features. This is typical to galaxies with a dominant population of older stars ( $\sim 1 \text{ Gyr}$ ), mixed with a population of younger hotter stars (Dressler & Gunn 1983; Covino et al. 2006), which in turn indicates that star formation is still present in the galaxy. The spectrum is dominated by lines of  $H\alpha/\beta \lambda\lambda 6563, 4861$ , [N II]  $\lambda\lambda 6548, 6583$ , [O III] lines  $\lambda\lambda 4959, 5007$ , as well as [O II]  $\lambda 3727$ . The difference between the SEDs of the host galaxy (black curve) and SN 2010mb (blue and red curves), also evident from color comparison of the host ( $g - r = 0.45$ ; Sloan Digital Sky Survey, SDSS) and the SN ( $g - r = -0.93$  on 2010 July 10 UT and  $g - r = -2.73$  on 2011 January 12 UT), shows that the blue quasi-continuum is not a result of residual contamination from the host galaxy.



**Figure 10.** Characterizing the interacting CSM. Left: a CSM density profile derived from the LC of SN 2010mb, assuming a symmetric spherical CSM distribution (solid points; see text for details). The data are best fitted by a power law profile of  $n \propto r^{-2.6}$  (blue line). For comparison, an  $n \propto r^{-2}$  profile is also presented (red line). Right: CSM column density derived by integrating over the density profile shown in the left panel. The high column density explains why X-ray photons have not been detected.

(A color version of this figure is available in the online journal.)



**Figure 11.** Host Galaxy nucleus spectrum. The spectrum was taken on 2012 February 20 UT using LRIS mounted on Keck-I 10 m telescope and is calibrated to  $r$  band photometry. Two spectra of SN 2010mb (2011 March 4 and July 5 UT) are given for comparison, showing the different spectral energy distribution between the host galaxy and the SN, and demonstrating that the SN is much bluer than the host, also evident from color comparison of the host ( $g - r = 0.45$ ; SDSS) and the SN ( $g - r = -0.93$  on 2010 July 10 UT and  $g - r = -2.73$  on 2011 January 12 UT). Spectra are available in digital form from the Weizmann Interactive Supernova Data Repository (WISEREP; Yaron & Gal-Yam 2012; <http://www.weizmann.ac.il/astrophysics/wiserep/>).

(A color version of this figure is available in the online journal.)

Host galaxy analysis was performed using two different spectra: for the SN position using the spectrum taken on 2010 September 5 UT, and for the galaxy nucleus, using the spectrum taken on 2012 February 20 UT. The SN offset from the host galaxy center is  $2''.55$  or 5.9 kpc in the radial direction at the redshift of the host (see Section 2.2). We detected and measured the emission line fluxes of  $H\alpha$ ,  $H\beta$ , [OII], [OIII], and [NII] using the standard IRAF procedure *splot*. We find that the host galaxy properties derived from the nucleus and from

the SN position agree with each other within the error bars, with the SN position values having larger errors because of the lower signal-to-noise ratio at the SN position. We estimate a reddening of  $E(B - V) = 0.19 \pm 0.03$  mag for the nucleus and  $E(B - V) \leq 0.09$  mag for the SN site, having used the observed ratio of the Balmer lines, and assuming the Cardelli et al. (1989) extinction law with  $R_V = 3.1$ , as well as case B recombination with a Balmer decrement of  $H\alpha/H\beta = 2.86$ . In the following, we correct the detected emission line fluxes

for reddening. To compute the oxygen abundance from these H II region emission lines, we follow Modjaz et al. (2011 and references therein) and employ the scales of Pettini & Pagel (2004; PP04-O3N2) and of Kewley & Dopita (2002, hereafter KD02), to obtain oxygen abundance values for the host galaxy nucleus of  $12 + \log(\text{O}/\text{H})_{\text{PP04-O3N2}} = 8.39 \pm 0.01$  and  $12 + \log(\text{O}/\text{H})_{\text{KD02}} = 8.60 \pm 0.05$ , respectively, where we consider only statistical uncertainties. The oxygen abundance at the SN site is consistent with that of the host galaxy nucleus within the error bars. We conclude that the metallicity of the SN host galaxy is  $0.5^{+0.42}_{-0.19} Z_{\odot}$ , having used the scale of Pettini & Pagel (2004) and a solar oxygen abundance value of  $12 + \log(\text{O}/\text{H}) = 8.69$  (Asplund et al. 2009). The total star formation rate that we measured using the H $\alpha$  luminosity (Kennicutt 1998) is  $\sim 0.32 M_{\odot} \text{ yr}^{-1}$ .

#### 4. DISCUSSION

##### 4.1. Indications for a Pulsational Pair Instability Event

The observed data are consistently explained by a PPI event followed by a CC SN. During the first  $\sim 150$  days after discovery,  $\sim 90\%$  of the optical display is coming from the SN explosion. At later times, as the ejecta reach the hydrogen-free CSM, a shock wave is generated at the interface between the two media, and X-ray photons generated at the shock front are exciting both the CSM and the ejecta, thus introducing a new source of energy. Eventually, this new source dominates the optical radiation—as can be seen from the long-lasting optical display compared to other Type Ic SNe, the change in the SED, and the evolution of the forbidden [O I] line. If we assume a non-spherical geometry for the CSM, the high energy photons will excite mostly the ejecta in the plane of the torus, thus perhaps generating the observed internal structure in the nebular lines (Figure 5). Mazzali et al. (2007) suggest that a flat-topped profile might be the result of a disk observed edge-on (leading to a double peak line profile), where the inner component of the disk would contain more mass, so that its emission could fill the valley between the two peaks.

A recent model of a  $110 M_{\odot}$  star in a low metallicity environment produces a stripped C/O core when encountering the pair instability for the first time (Chatzopoulos & Wheeler 2012). The period of pair instability will result in the ejection of a  $\sim 3 M_{\odot}$  shell that is significantly enhanced in He, C, and O. Subsequent pulses may be even richer in carbon and oxygen, as they probe inner regions of the progenitor. Later, the progenitor will return to the normal evolution track, until it ends its life in an iron-CC SN. In the case of SN 2010mb, we estimate the time gap between the last pulse and the CC event was  $\sim 2.2$  yr, if we assume a velocity of  $800 \text{ km s}^{-1}$  for the matter ejected in the final pulse based on the observed [O I]  $\lambda 5577$  recombination line blueshift, again in accordance with models (Woosley et al. 2007). While in the past PISNe and PPI events were assumed to happen only in very low metallicity environments, recent works (Langer et al. 2007; Chatzopoulos & Wheeler 2012; Yusof et al. 2013) predict these explosions can also occur in local dwarf galaxies with metallicity values which are well within the errors of those measured for SN 2010mb, i.e., metallicities as high as  $Z = Z_{\odot}/3$ .

##### 4.2. H-free CSM Disposition Around Massive Stars

Interaction of SN ejecta with hydrogen-free material is also observed in Type Ibn SNe (Pastorello et al. 2008a; Smith et al. 2012). These events are defined by strong He I emission lines,

with a blue quasi-continuum, which are suggested to be the result of a blend of Fe II lines that are observed in at least two cases (SN 1999cq, Matheson et al. 2000a; SN 2006jc, Foley et al. 2007). Type Ibn SNe are associated with the explosion of WR progenitors embedded in a helium-rich hydrogen-diluted CSM. The interacting material composition is consistent with that of WR winds, i.e., hydrogen and helium for WN progenitors, and helium for WC/WO progenitors (Pastorello et al. 2008b). The main differences between Type Ibn SNe and SN 2010mb are as follows: (1) Type Ibn SNe have much shorter timescales than that of SN 2010mb; therefore,  $^{56}\text{Ni}/^{56}\text{Co}$  mixing in the outer shells is not a possible explanation for the blue quasi-continuum in our case as it decays too fast (a half-life of 8/77 days, respectively). Scattering from interstellar dust Matheson et al. (2000a) as echoes also cannot explain the plateau because the integrated flux during the plateau is higher than the flux from the peak in the case of SN 2010mb. (2) In Type Ibn SNe, *B* band and *g* band magnitudes decrease with time. (3) While in Type Ibn the interaction has a modest influence on the observed LC and a fairly modest amount of CSM is involved (less than  $\sim 0.3 M_{\odot}$ ; Pastorello et al. 2008a), in the case of SN 2010mb, CSM interaction is the dominant power source at late times, and it probably involves a much higher mass of CSM ( $> 3 M_{\odot}$ ). We conclude that a similar scenario to the one generating Type Ibn SNe is not adequate for SN 2010mb.

The presence of a large amount of CSM can also be attributed to wave excitation by vigorous convection in the late stages of stellar evolution (Si nuclear burning phase; Quataert & Shiode 2012). In this scenario, convective motions excite internal gravity waves that in some cases are converted to sound waves as they tunnel toward the stellar surface. As the sound waves dissipate while crossing through the star envelope, they release a large amount of energy that can unbind up to several solar masses of the stellar envelope. In case of SN 2010mb, the interaction goes for long time, indicating the CSM must extend over large distances, and it was ejected at an earlier phase than the Si nuclear burning phase of the progenitor. It remains to be demonstrated that high mass loss rates can be produced during earlier nuclear burning stages in order to explain the optical display of SN 2010mb.

Luminous blue variable stars (Humphreys & Davidson 1994, Meynet et al. 2011) are also associated with strong eruptions. Maeder (1992) has shown that in the outer shells of these stars, the free fall time becomes longer than the thermal diffusion time, leading to super-Eddington luminosity, and the ejection of a large amount of material. In cases where this process sets in after the progenitor lost its hydrogen and helium outer shells, the ejected material will be comprised from intermediate mass elements, mainly O and C, as seen in our case. This option is speculative at this time and needs to be further explored.

Finally, another event where interaction of SN ejecta with H-free CSM was suggested as a way to explain an increased, short lived luminosity is SN 2009dc, an SN interpreted as a 'Super Chandrasekhar' Type Ia SN. Hachinger et al. (2012) suggest that a non-violent white dwarf merger can culminate in a Type Ia SN interacting with H-/He-free CSM, causing an increase in the observed flux with respect to the flux of a typical Type Ia SN. In such a case, if the emitting CSM is concentrated in clumps or in a torus, the observed emission will be a sum of two components: one from the SN explosion, and the other from the ejecta-CSM interaction, as is the case for SN 2010mb. However, the total amount of radiation that can be added by interaction with white dwarf debris is much smaller than seen for SN 2010mb.

### 4.3. Other Scenarios

#### 4.3.1. Photon Diffusion

Type II-P SNe also exhibit a plateau in their LCs. The mechanism driving this plateau is photon diffusion through the expanding ejecta. We analyze the SN 2010mb LC in the context of the recombination front model Popov (1993). In this model, the opacity is approximated by a step function and is constant above the ionization temperature,  $T_{\text{ion}}$ , and equals zero below that. We assume that above the ionization temperature the opacity is dominated by Thomson scattering, and since the envelope contains mainly metals, we take  $k = 0.2 \text{ cm}^2 \text{ g}^{-1}$ . This approximation to the opacity is an upper bound on the real opacity, which is already smaller than  $0.2 \text{ cm}^2 \text{ g}^{-1}$  at temperatures of  $T_{\text{ion}} \approx 3 \times 10^4 \text{ K}$  Rabinak & Waxman (2011). We assume the effective temperature of the radiation is  $\approx 5,000 \text{ K}$  from the models described in Section 3.2, and we get an ion temperature of  $T_{\text{ion}} \approx (T_{\text{eff}}/2^{1/4}) \approx 4170 \text{ K}$ . Based on this analytic model, we get the following estimation for the ejecta mass, explosion energy, and progenitor radius:

$$M_{\text{ej}} \sim 1160 t_{185}^4 T_{4170}^4 v_5^3 L_{41}^{-1} \kappa_{0.2}^{-1} M_{\odot}$$

$$E \sim 8.4 \times 10^{52} t_{185}^4 T_{4170}^4 v_5^5 L_{41}^{-1} \kappa_{0.2}^{-1} \text{ erg}$$

$$R_s \sim 2.3 \times 10^{11} t_{185}^{-2} T_{4170}^{-4} v_5^{-4} L_{41}^2 \kappa_{0.2} \text{ cm},$$

where  $t_{185} = t_p/185$  days,  $T_{4170} = T_{\text{ion}}/4170 \text{ K}$ ,  $L_{41} = L/10^{41} \text{ erg s}^{-1}$ ,  $v_5 = v_{ph}/5 \times 10^8 \text{ cm s}^{-1}$ , and  $\kappa_{0.2} = \kappa/0.2 \text{ cm}^2 \text{ g}^{-1}$ . Even when considering the uncertainties in the measurements, the model suggests an unrealistically large ejecta mass (several hundred solar masses of He, C, and O), and therefore, we pursue it no more.

#### 4.3.2. Magnetar

Magnetars, very highly magnetic ( $B > 10^{14} \text{ G}$ ) and rapidly rotating ( $P_s \approx 2\text{--}5 \text{ ms}$ ) NSs, generated at the time of SN explosion, can have a large impact on SN LCs (Mazzali et al. 2006; Maeda et al. 2007; Kasen & Bildsten 2010; Woosley 2010; Inserra et al. 2013). These magnetars spin down in a few months (comparable to the radiative diffusion time in SN ejecta,  $t_d$ ), depositing enough energy to power the rare superluminous ( $L > 10^{44} \text{ erg s}^{-1}$ ) hydrogen poor SLSN-I Quimby et al. (2011); Gal-Yam (2012); Inserra et al. (2013). The much lower (and nearly constant) luminosity,  $L_{\text{SN}} \approx 10^{42} \text{ erg s}^{-1}$ , of SN 2010mb over a prolonged ( $> 500$  days) period requires a different interpretation. Specifically, a young NS with a magnetic dipole spin-down luminosity  $L_p \approx 10^{42} \text{ erg s}^{-1}$ , and a spin-down timescale,  $t_p$ , longer than 10 yr, so that  $L_p$  is approximately constant over the 500 days, is needed. Such a power source (Kasen & Bildsten 2010) would overwhelm the adiabatic losses from the initial explosion, and reset the SN luminosity to  $L_p$ . Setting  $L_p = 10^{42} \text{ erg s}^{-1}$  gives the relation  $B \approx 3 \times 10^{13} \text{ G} (P_s/8 \text{ ms})^2$ , and requiring a spin-down timescale longer than 10 yr implies that  $P_s < 8 \text{ ms}$ . Though more rapidly rotating than the fastest known young pulsar, PSR J0537-6910 (Marshall et al. 1998) (at  $P_s = 16 \text{ ms}$ ), such an initial spin and magnetic field are certainly allowed (e.g., the implied  $B \approx 4 \times 10^{12} \text{ G}$  for an initial spin of 3 ms is entirely reasonable). In this scenario, the young NS spin-down luminosity is reprocessed by the SN ejecta, allowing for a long, visible event. Though we remain unsure as to the form

(i.e., Poynting flux, particles, or radiation) of the spin-down power,  $L_p$ , we will assume that the thinning of the ejecta due to expansion will eventually lead to an inefficient coupling and, hence, a reduction in the optical brightness of the event. For an ejecta mass of  $M_{\text{ej}} \approx 5 M_{\odot}$  moving at an average velocity fixed by  $E = M_{\text{ej}} V^2/2$ , the column depth in the ejecta will evolve with time as

$$\frac{M_{\text{ej}}}{4\pi(Vt)^2} = 2 \frac{\text{g}}{\text{cm}^2} \left( \frac{500 \text{ days}}{t} \right)^2 \left( \frac{10^{51} \text{ erg}}{E} \right) \left( \frac{M_{\text{ej}}}{5 M_{\odot}} \right)^2, \quad (1)$$

implying thinning to fast energetic ions and hard radiation at about the time where the SN faded.

The magnetar model can explain the LC timescale but fails to explain the color evolution of SN 2010mb, the high density narrow O lines, and the presence of strong nebular lines that require an optically thin medium, while the reprocessing of the radiation emitted by the magnetar requires a thick medium.

## 5. CONCLUSIONS

Observations of SN 2010mb give direct evidence for interaction of SN ejecta with a large amount of hydrogen-free CSM. The extended LC, the rise in magnitude in the  $g$  band and  $B$  band, and the blue quasi-continuum that becomes more significant with time, show that a mechanism external to the SN explosion is injecting energy at late times. The presence of the [O I]  $\lambda 5577$  line and its evolution shows the external source is interaction of the SN ejecta with a large mass of hydrogen-free CSM at high densities at the vicinity of the SN. Finally, the late transition from photospheric to nebular emission and the spectral decomposition based on the similar SN 1997ef shows that SN 2010mb ejected a large mass  $\geq 12 M_{\odot}$ , compared to  $1.7 M_{\odot}$  in “normal” Type Ic SNe (Drout et al. 2011), as is expected in an explosion of an extremely massive progenitor.

Combining all these signatures, a PPI scenario naturally comes to mind. In such a scenario, the period of PPI will result in the ejection of a few solar masses of matter that is composed mostly of He, C, and O. After the PPI period, the progenitor will return to the normal evolution track associated with massive stars, until it ends its life due to photo-disintegration and neutronization at its core (a CC). In case of a successful SN following the CC, the expected optical signature is a long-lasting event, where the SN ejecta will interact with the large amount of CSM ejected in the recent past. Other scenarios, such as wave-driven mass loss during late stages of nuclear burning, may be possible as well. A similar, if more intense, hydrogen-free CSM interaction may explain the energy source behind SLSN-I events (Quimby et al. 2011; Chomiuk et al. 2011; Pastorello et al. 2010; Leloudas et al. 2012; Inserra et al. 2013; see Gal-Yam 2012 for a review), which are found in growing numbers by various sky surveys in recent years.

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